

# Interpretations of recent MINOS results

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Los Alamos National Laboratory, October 21st, 2010



partly based on work done collaboration with  
Pedro A. N Machado and Stephen J. Parke (arXiv:1009.0014)

# Outline

- 1 The MINOS experiment and its results—a hint for CPT violation?
- 2 Explanation attempts
  - Low statistics?
  - A systematic error?
  - “Real” *CPT* violation?
  - Effective *CPT* violation: Non-standard neutrino matter effects
  - A new long range force?
  - A *CP*-violating charged current interaction?
- 3 Non-standard neutrino interactions in renormalizable models
- 4 Conclusions

# Hinchcliffe's theorem

“When a title is in the form of a question,  
the answer is always NO.”

see, however:

**IS HINCHLIFFE'S RULE TRUE? ·**

Boris Peon

Abstract

Hinchcliffe has asserted that whenever the title of a paper is a question with a yes/no answer, the answer is always no. This paper demonstrates that Hinchcliffe's assertion is false, but only if it is true.

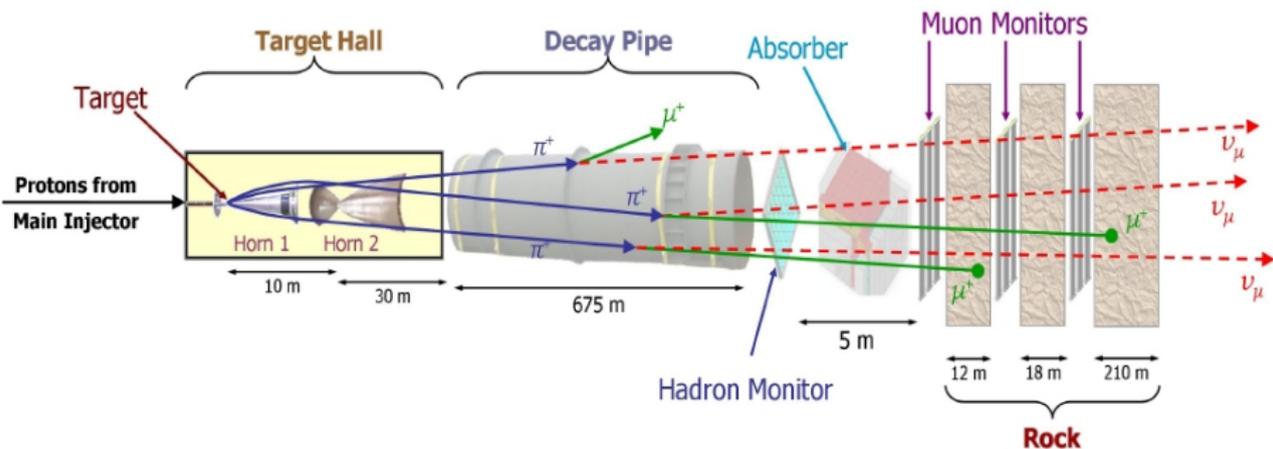
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# Disclaimer

I'm **not a member of the MINOS collaboration**,  
and what I'm going to say in this talk  
is entirely **my own responsibility**.

# The NuMI beam



- $\nu_\mu$  from decay in flight of focused  $\pi^+ + K^+$  beam
- Intrinsic backgrounds:  $\bar{\nu}_\mu, \nu_e$  from subdominant  $\pi^+, K^+$  decays (e.g.  $K^+ \rightarrow \pi^0 e^+ \nu_e$ ) and from muon decays
- Polarity of focussing system can be inverted to obtain  $\bar{\nu}_\mu$  beam.

Image credit: MINOS collaboration, <http://www.numi.fnal.gov/>

# The MINOS experiment



Far detector:

- 5.4 kt magnetized iron, interleaved with solid scintillator plates to record tracks

Near detector:

- Similar to the far detector but smaller
- Measures unoscillated neutrino flux  
⇒ reduction of systematic uncertainties

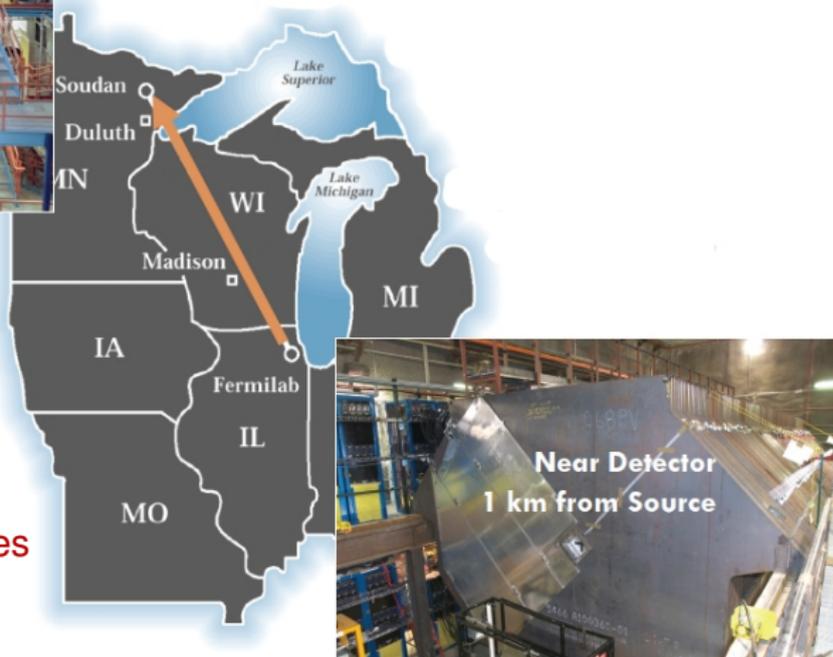
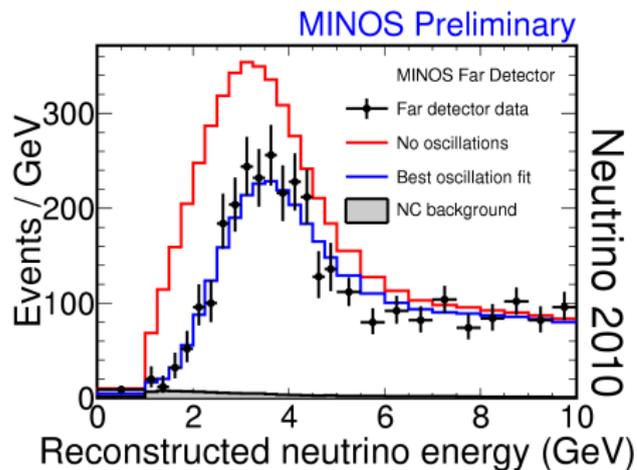
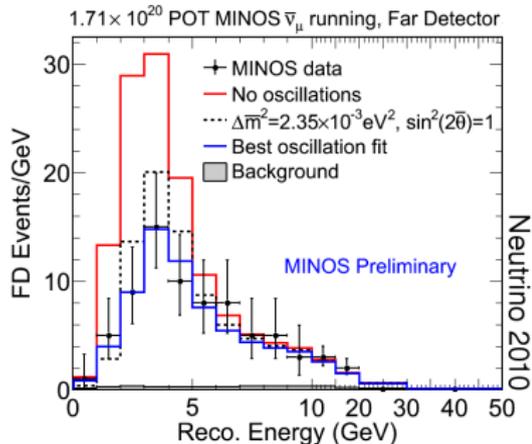


Image credit: MINOS collaboration, <http://www.numi.fnal.gov/>

# MINOS $\nu_\mu, \bar{\nu}_\mu$ disappearance data



$\nu_\mu$

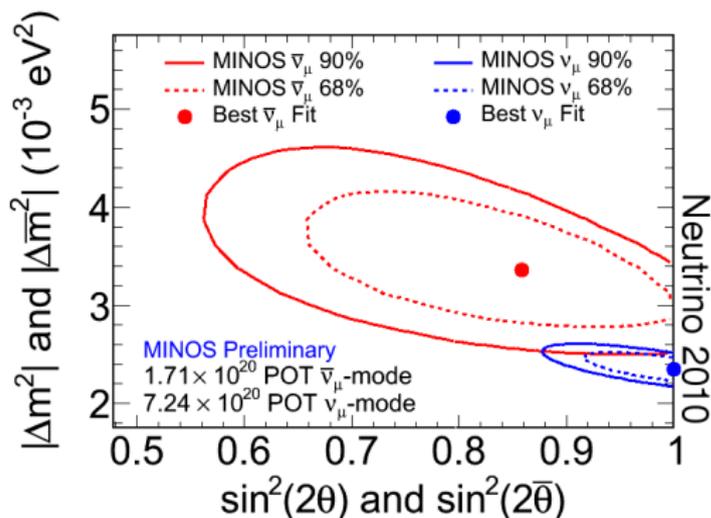


$\bar{\nu}_\mu$

Image credit: MINOS collaboration, <http://www.numi.fnal.gov/>  
This result first presented by P. Vahle at Neutrino 2010

# MINOS oscillation fit

## MINOS fit



## Comparison to our fit

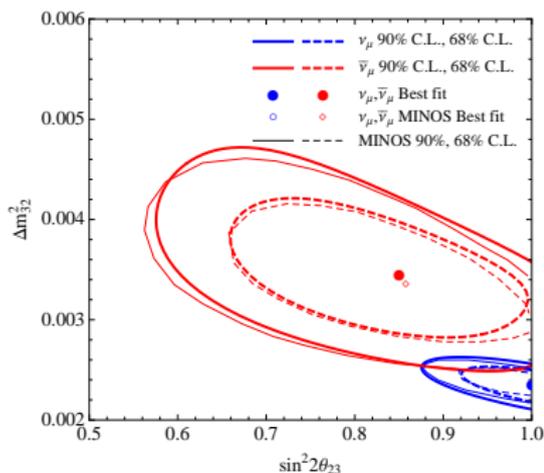


Image credit: MINOS collaboration, <http://www.numi.fnal.gov/>  
 This result first presented by P. Vahle at Neutrino 2010

- Two-flavor fits:  $P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta \sin^2 \frac{\Delta m^2 L}{4E}$
- Separate fits for neutrinos and anti-neutrinos differ at  $\sim 2\sigma$  confidence level.

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# Low statistics?

$\bar{\nu}_\mu$  sample is about 20 times smaller than  $\nu_\mu$  sample.

⇒ Effect might go away with more statistics

# A systematic error?

I can only speculate . . .

# CPT violation?

## Why not just CP violation?

- $\nu_\mu \rightarrow \nu_\mu$  is a  $T$ -invariant process
- By virtue of  $CPT$ , it must conserve  $CP$ .
- Note:  $CP$  violation in interactions is a possibility—see later

## Phenomenological parameterization

Barenboim Lykken arXiv:0908.2993

- Assume mixing matrices for  $\nu$  and  $\bar{\nu}$  to be completely independent.
- Fit to MINOS (older dataset), Super-K, KamLAND, CHOOZ yields

	$\nu$	$\bar{\nu}_\mu$
$\Delta m_{31}^2$ [eV <sup>2</sup> ]	0.0025	0.02
$\sin^2 2\theta_{23}$	1.0	0.55



$\Delta \bar{m}_{\text{atm}}^2$

$\Delta m_{\text{atm}}^2$



$\Delta m_{\text{solar}}^2$



$\Delta \bar{m}_{\text{solar}}^2$

# CPT violation?

Another parameterization:

- Introduce Lorentz- and *CPT*-violating operators like  $A_\mu \bar{\psi} \gamma^\mu \psi$   
(with  $A_\mu$  a constant 4-vector)
- Studied in detail in Dighe Ray [arXiv:0802.0121](https://arxiv.org/abs/0802.0121)  
(but not in the context of MINOS)

# A model of spontaneous $CPT$ violation

Mukohyama Park arXiv:1009.1251

- Ghost condensation ( $\langle \partial_0 \phi \rangle \neq 0$ ) on a distant brane in 5D.  
⇒ preferred frame
- Right-handed neutrinos propagating in the bulk couple to  $\partial_\mu \phi$  and to  $\nu_L$ .
- After ghost-condensation, Lorentz-violating neutrino mass terms are generated.

# Neutrino matter effects

In the Standard Model:

$$\begin{aligned}\mathcal{L}_{\text{eff}} &\sim -2\sqrt{2}G_F [\bar{e}\gamma^\mu P_L \nu_e] [\bar{\nu}_e \gamma_\mu P_L e] \\ &\sim -2\sqrt{2}G_F [\bar{e}\gamma^\mu P_L e] [\bar{\nu}_e \gamma_\mu P_L \nu_e]\end{aligned}$$

In ordinary matter

$$\begin{aligned}\langle \bar{e}\gamma^0 e \rangle &= n_e & \langle \bar{e}\vec{\gamma} e \rangle &= \langle \vec{v}_e \rangle = 0 \\ \langle \bar{e}\gamma^0 \gamma^5 e \rangle &= \langle \vec{\sigma}_e \vec{p}_e / E_e \rangle = 0 & \langle \bar{e}\vec{\gamma} \gamma^5 e \rangle &= \langle \vec{\sigma}_e \rangle = 0\end{aligned}$$

Potential felt by electron neutrinos in ordinary matter:

$$V = \sqrt{2}G_F n_e$$

Sign changes for  $\nu_\mu \leftrightarrow \bar{\nu}_\mu$

⇒ **Effective CPT violation** due to **CPT**-asymmetric background matter

In the SM, these effects are **far too small** to explain MINOS  $\nu_\mu$  disappearance data since they are **suppressed** by  $\theta_{13}$ ,  $\Delta m_{21}^2 / \Delta m_{31}^2$

# Non-standard matter effects

Many previous works on NSI in MINOS, but mostly focussing on the  $\nu_\mu \rightarrow \nu_e$  appearance channel.

Friedland Lunardini

Two modes of searching for new neutrino interactions at MINOS

hep-ph/0606101

Kitazawa Sugiyama Yasuda

Will MINOS see new physics?

hep-ph/0606013

Blennow Ohlsson Skrotzki

Effects on non-standard interactions in the MINOS experiment

hep-ph/0702059

# Non-standard matter effects

Consider a neutral current (NC) **non-standard interaction** (NSI) of the form

$$\mathcal{L}_{\text{NSI}} \sim -2\sqrt{2}G_F\epsilon_{\alpha\beta}^f [\bar{f}\gamma^\mu f] [\bar{\nu}_\alpha\gamma_\mu P_L\nu_\beta] \quad f = e, \mu, \tau,$$

leading to **off-diagonal** (flavor-violating) and/or **non-universal** matter potential. In the flavor basis,

$$V = \sqrt{2}G_F n_e \begin{pmatrix} 1 + \epsilon_{ee} & \epsilon_{e\mu} & \epsilon_{e\tau} \\ \epsilon_{e\mu}^* & \epsilon_{\mu\mu} & \epsilon_{\mu\tau} \\ \epsilon_{e\tau}^* & \epsilon_{\mu\tau}^* & \epsilon_{\tau\tau} \end{pmatrix}.$$

The oscillation probability is

$$P(\nu_\alpha \rightarrow \nu_\beta) = |\langle \nu_\beta | e^{-iHt} | \nu_\alpha \rangle|^2, \quad H = \frac{1}{2E} U \begin{pmatrix} 0 & & \\ & \Delta m_{21}^2 & \\ & & \Delta m_{31}^2 \end{pmatrix} U^\dagger + V.$$

For  $\bar{\nu}$ :  $U \rightarrow U^*$ ,  $V \rightarrow -V$   
 $\Rightarrow$  **Effective CPT violation**

# Non-standard matter effects in the $\mu$ - $\tau$ sector

Two-flavor calculation leads to

$$P(\nu_\mu \rightarrow \nu_\mu) = 1 - \sin^2 2\theta_N \sin^2 \left( \frac{\Delta m_N^2 L}{4E} \right)$$

with

$$\Delta m_N^2 = [(\Delta m_{32}^2 \cos 2\theta_{23} + \epsilon_{\tau\tau} A)^2 + |\Delta m_{32}^2 \sin 2\theta_{23} + 2\epsilon_{\mu\tau} A|^2]$$

$$\sin^2 2\theta_N = |\Delta m_{32}^2 \sin 2\theta_{23} + 2\epsilon_{\mu\tau} A|^2 / \Delta m_N^4,$$

and  $A = A = 2\sqrt{2}G_F n_e E$ . (we set  $\epsilon_{\mu\mu} = 0$  since flavor-universal terms can be subtracted from  $V$ )

Note the following symmetries:

$$\arg(\epsilon_{\mu\tau}) \rightarrow 2\pi n - \arg(\epsilon_{\mu\tau})$$

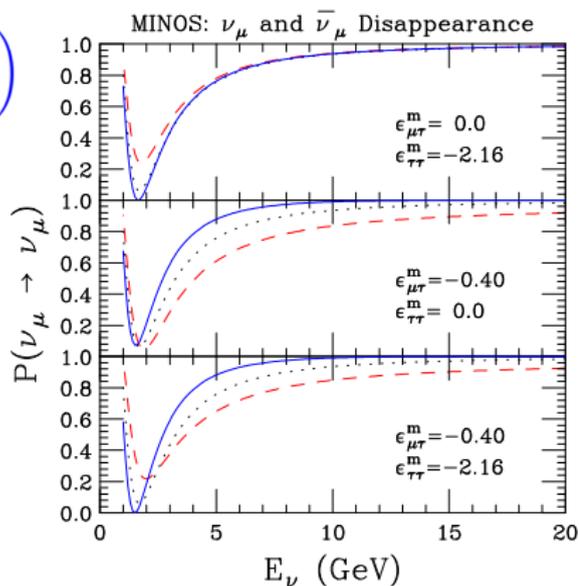
$$\epsilon_{\mu\tau} \rightarrow -\epsilon_{\mu\tau},$$

$$\epsilon_{\tau\tau} \rightarrow -\epsilon_{\tau\tau},$$

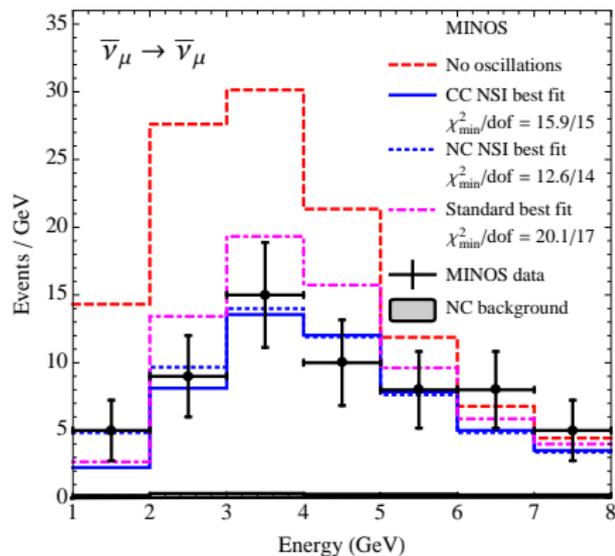
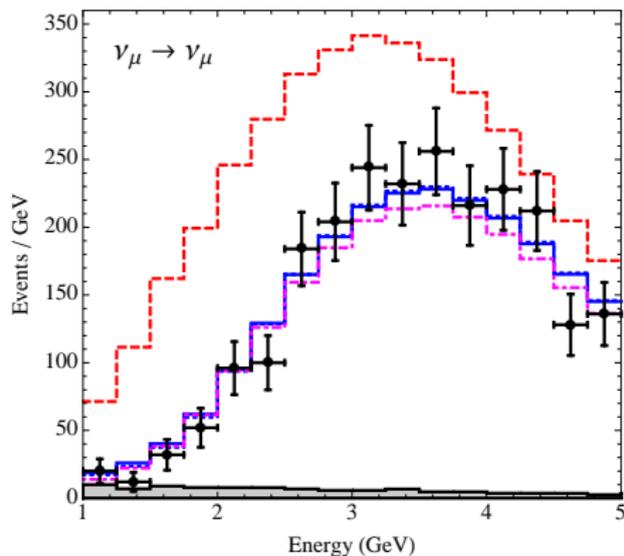
$$\Delta m_{32}^2 \rightarrow -\Delta m_{32}^2,$$

$$\epsilon_{\tau\tau} \rightarrow -\epsilon_{\tau\tau},$$

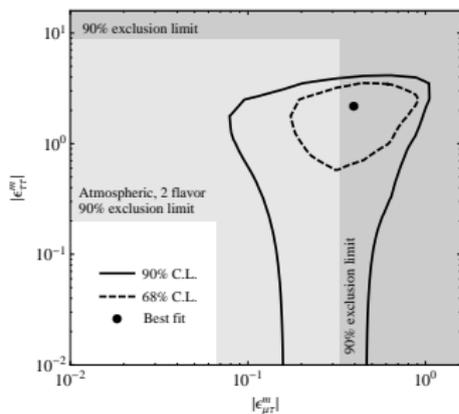
$$\theta_{23} \rightarrow \frac{\pi}{2} - \theta_{23}.$$



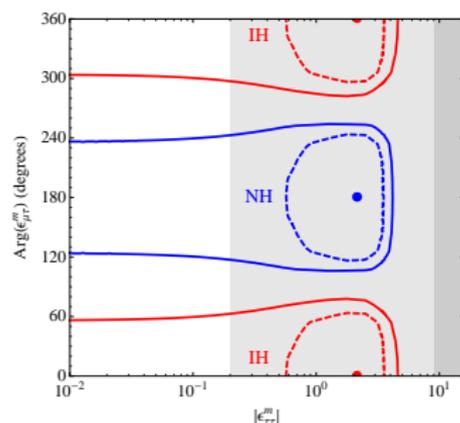
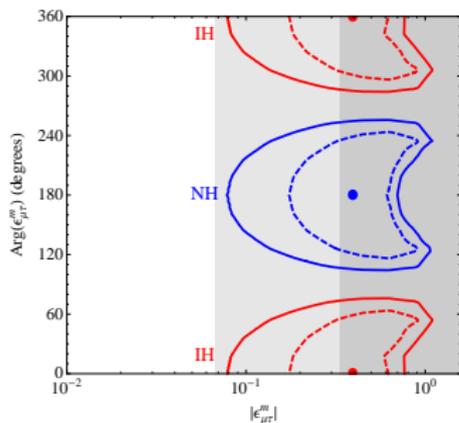
# Non-standard matter effects in MINOS?



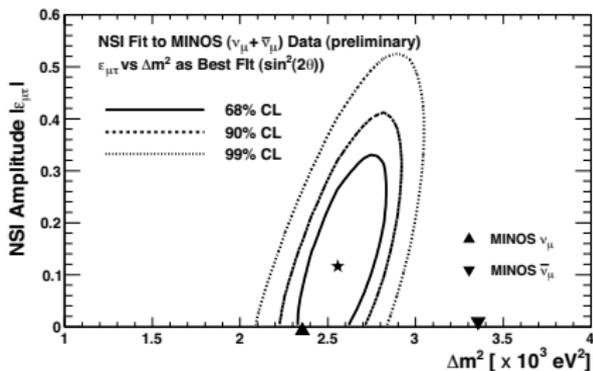
# Non-standard matter effects in MINOS? (2)



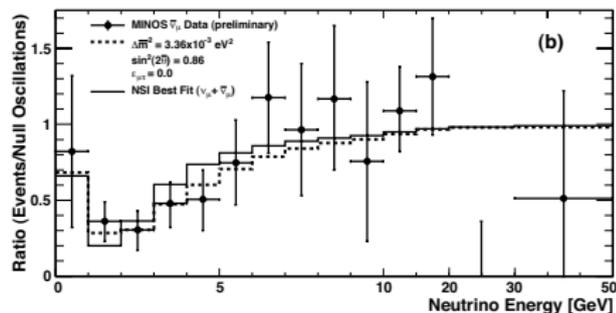
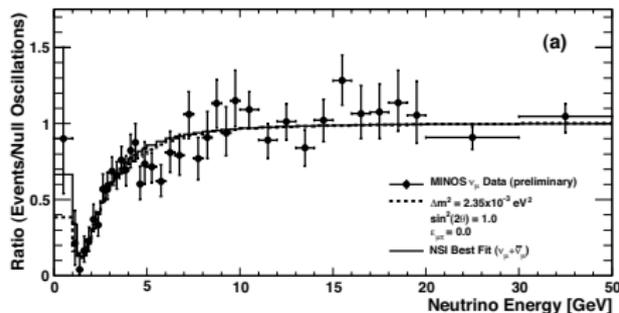
- $|\epsilon| \gtrsim 0.1$  required (almost as strong as SM weak interactions)
- **Consistent** with constraints on  $\epsilon_{\mu\tau}$  from CHARM ( $\nu_\mu e \rightarrow \nu e$ ) and NuTeV ( $\nu_\mu q \rightarrow \nu q$ )
- **Consistent** with constraints on  $\epsilon_{\tau\tau}$  from  $\Gamma_{inv}^{Z^0}$
- **Disfavored** by atmospheric neutrinos (These are 2-flavor limits, may not be robust)
- **Model-dependent constraints:** See later



# A similar analysis



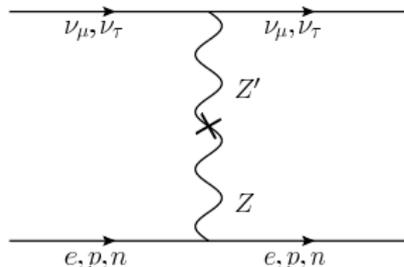
- Assume only  $\epsilon_{\mu\tau} \neq 0$
- Fit to extracted **oscillation probability** rather than **spectrum**.
- Results agree with ours **qualitatively**, but not **quantitatively**.
- Possible reason: Fit to **probability** cannot fully include effect of **experimental energy resolution**



# A new long-range force?

Heeck Rodejohann arXiv:1007.2655

- A very light  $L_\mu - L_\tau$  gauge boson  $Z'$  ( $m_{Z'} \lesssim 10^{-18}$  eV  $\sim 1$  a.u. $^{-1}$ )
- Very weak couplings ( $\alpha \lesssim 10^{-50}$ )
- Mixing with the SM  $Z$



- $\nu_\mu, \nu_\tau$  feel potential generated by the Sun (contribution from the Earth is  $\sim 3$  times smaller)
- Since the Sun contains no anti-matter, and since  $\nu$  and  $\bar{\nu}$  have opposite  $L_\mu - L_\tau$  charges), this leads to effective  $CPT$  violation.
- Phenomenologically equivalent to  $\epsilon_{\mu\mu}, \epsilon_{\tau\tau}$ .

# Non-standard charged current interactions

- Remember:  $\nu_\mu \rightarrow \nu_\mu$  is  $CP$ -invariant
- **But:**  $\pi(\text{ source}) \rightarrow ??? \rightarrow \mu(\text{ detector})$  does not have to be.
- Two possibilities
  - ▶ Modified  $\nu_\mu$  flux at far detector, **but not** at near detector.  
 $\nu_\tau$  contamination in the NuMI beam?  
 $\Rightarrow$  **Ruled out by NOMAD.**
  - ▶ A new interaction of the form

$$\nu_\tau + N \rightarrow X + \mu,$$

e.g.

$$\mathcal{L}_{\text{NSI}} \supset -2\sqrt{2}G_F\epsilon_{\tau\mu}^d V_{ud} [\bar{u}\gamma^\rho d] [\bar{\mu}\gamma_\rho P_L\nu_\tau] + h.c.$$

- If the new interaction is **vector-like**, it will not contribute to  $\pi \rightarrow \mu\nu_\tau$ , which is constrained by NOMAD.

# Non-standard charged current interactions (2)

“Apparent” oscillation probability:

$$\begin{aligned} \tilde{P}(\nu_\mu \rightarrow \nu_\mu) = & \\ & 1 - \left[ 1 + 2 |\epsilon_{\tau\mu}^d| \cot 2\theta_{23} \cos [\arg(\epsilon_{\tau\mu}^d)] - |\epsilon_{\tau\mu}^d|^2 \right] \sin^2 2\theta_{23} \sin^2 \left( \frac{\Delta m_{32}^2 L}{4E} \right) \\ & + 2 |\epsilon_{\tau\mu}^d| \sin 2\theta_{23} \sin [\arg(\epsilon_{\tau\mu}^d)] \sin \left( \frac{\Delta m_{32}^2 L}{4E} \right) \cos \left( \frac{\Delta m_{32}^2 L}{4E} \right) \end{aligned}$$

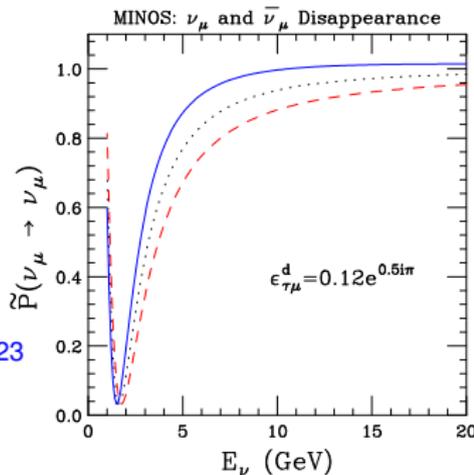
For anti-neutrinos:

$$\arg(\epsilon_{\tau\mu}^d) \rightarrow -\arg(\epsilon_{\tau\mu}^d)$$

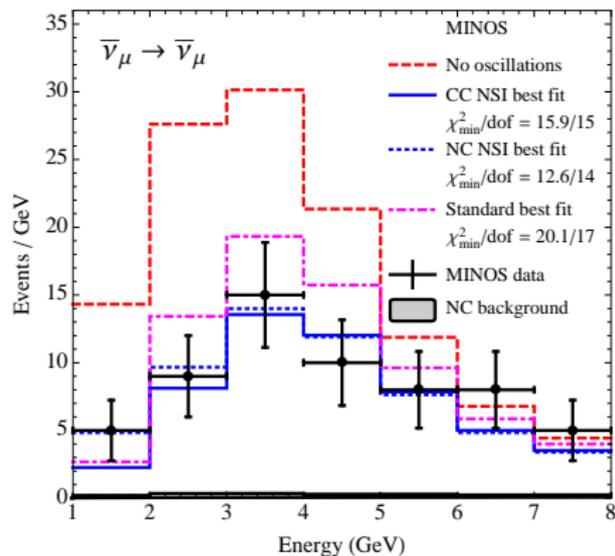
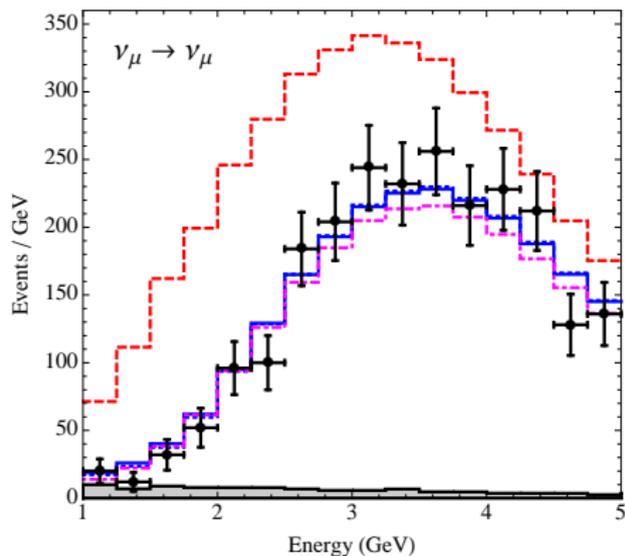
Symmetries:

$$\begin{aligned} \arg(\epsilon_{\tau\mu}^d) &\rightarrow 2\pi n - \arg(\epsilon_{\tau\mu}^d), & \Delta m_{32}^2 &\rightarrow -\Delta m_{32}^2 \\ \arg(\epsilon_{\tau\mu}^d) &\rightarrow (2n+1)\pi - \arg(\epsilon_{\tau\mu}^d), & \theta_{23} &\rightarrow \frac{\pi}{2} - \theta_{23} \end{aligned}$$

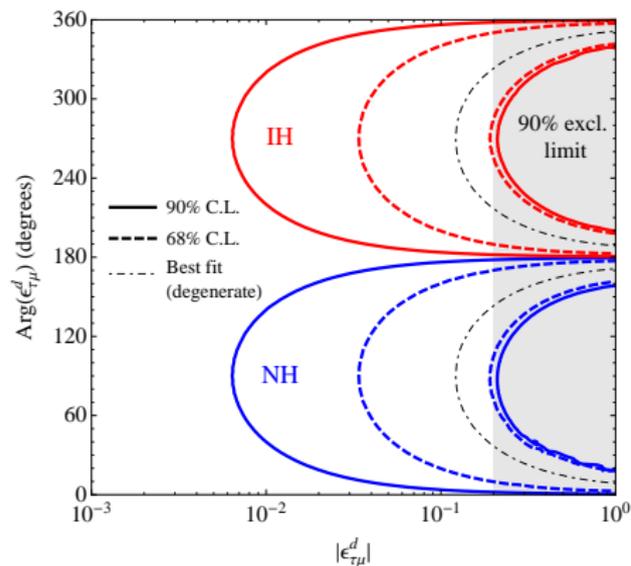
(The second of these can be generalized to a continuous symmetry.)



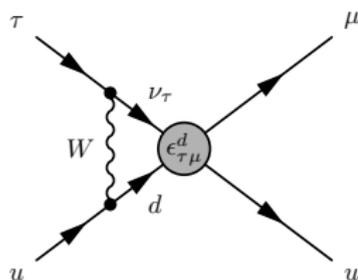
# CC NSI in MINOS?



## CC NSI in MINOS? (2)



- $|\epsilon| \gtrsim 0.1$  required (almost as strong as SM weak interactions)
- Consistent with model-independent constraint from  $\tau \rightarrow \mu + \text{hadrons}$

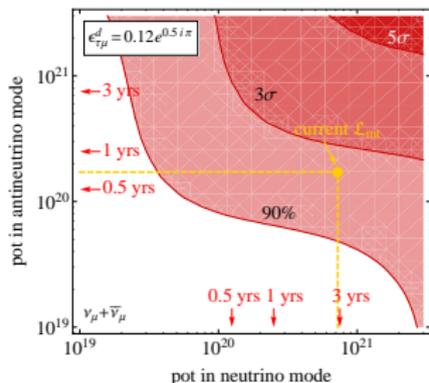


(Model-independent = consider only log-divergent part)

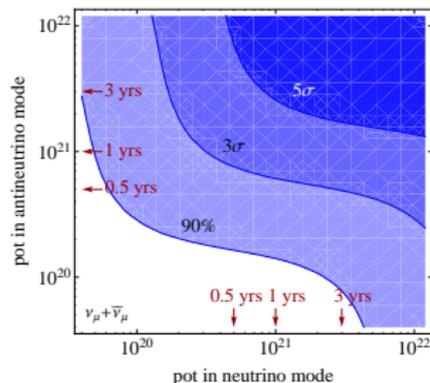
- Hard to embed in a renormalizable model

# Future tests

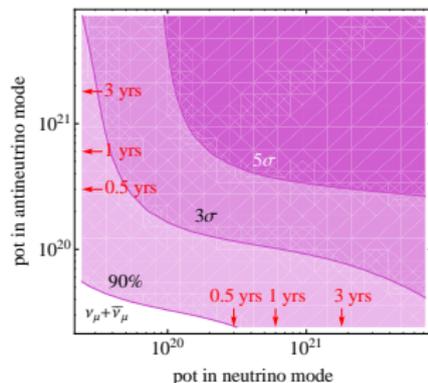
CC NSI discovery reach in MINOS



CC NSI discovery reach in T2K



CC NSI discovery reach in NOvA



- Ideally, slightly more time spent on  $\bar{\nu}$  running  
 $\Rightarrow$  Similar statistics in  $\nu$  and  $\bar{\nu}$  in spite of lower  $\bar{\nu}$  cross section
- At this time: More  $\bar{\nu}$  running in MINOS desirable

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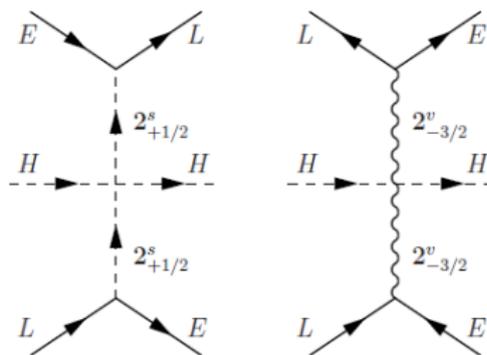
# Neutrino NSI from new physics at high scales

Aim: Relate NSI operators to **renormalizable** model

- $SU(2)$  invariant operators for neutrino NSI are usually **accompanied by charged lepton NSI**, which are heavily constrained.  
(Exception: NC  $[\bar{\nu}_\tau \nu_\tau][\bar{f}f]$  couplings)

see e.g. Antusch Baumann Fernández-Martínez arXiv:0807.1003  
Gavela Hernandez Ota Winter arXiv:0809.3451

- One way out: **Dimension 8 operators**, e.g.  $[\bar{E}^c_\gamma \gamma^\rho L_\alpha][\bar{L}^\beta \gamma_\rho E^c_\delta]$



- ▶ Requires **new mediators**
- ▶ Requires **cancellation** between couplings to avoid large **dim-6** effects.

# Neutrino NSI from light new physics

- Many constraints on NSI come from **high-energy** ( $\mathcal{O}(\text{GeV})$ ) processes.
- On the other hand, assume new mediator(s) with **very small masses**  $m$  and with **extremely** weak coupling  $g$

Nelson Walsh arXiv:0711.1363; Engelhardt Nelson Walsh arXiv:1002.4452

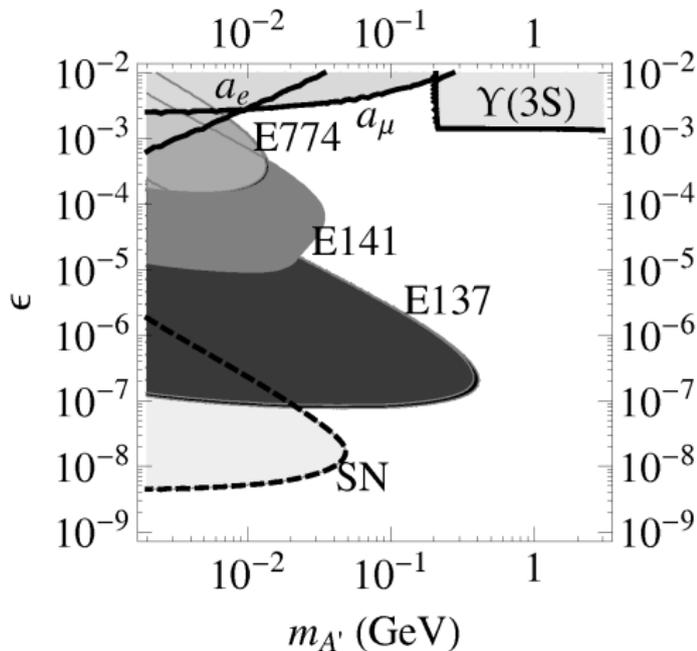
- ▶ high-energy cross sections/rates suppressed by  $g^4$
  - ▶ **Coherent forward scattering** ( $q^2 = 0$ ) only suppressed by  $(g^2 \sin^2 \theta_w / e^2)(M_W^2 / m^2)$
  - ▶ ... can be **relatively large**
- Light new physics also motivated by **Dark Matter** (Sommerfeld enhancement)

# Constraints on light new physics

Constraints on new kinetically mixed  $U(1)$  gauge boson from

- Fixed target (beam dump) experiments
- Supernova cooling
- Electron/Muon  $g-2$
- BaBar search for  $\Upsilon(3S) \rightarrow \gamma Z' \rightarrow \gamma \mu^+ \mu^-$

Note: For flavor-violating couplings, some constraints may become weaker



Bjorken Essig Schuster Toro arXiv:0906.0580

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# Conclusions

- MINOS observes interesting  $\sim 2\sigma$  discrepancy between  $\nu_\mu$  and  $\bar{\nu}_\mu$  oscillations.
- Possible explanations:
  - ▶ Low  $\bar{\nu}_\mu$  statistics
  - ▶ Systematic effect
  - ▶  $CPT$  violation
  - ▶ Non-standard neutrino matter effects (NC NSI)
  - ▶ A long-range force from the Sun
  - ▶ A  $CP$ -violating CC interaction (CC NSI)
- All these effects are not generic in extensions of the Standard model, but can be accommodated.
- In the future
  - ▶ MINOS will collect more  $\bar{\nu}$  statistics
  - ▶ T2K and NO $\nu$ A can confirm or refute the effect, provided they are operated also in  $\bar{\nu}$  mode.

Thank you!