

# Making NLO predictions for Method 2

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# Introduction

- We must understand QCD production of:

$W + \text{jets}$

$W + b\bar{b}$

$W + b\bar{b} + \text{jets}$

as large  
backgrounds to:

$t\bar{t}$

single top

Higgs production

other new physics

- There are many LO tools (ALPGEN, COMPHEP, Madevent, etc.) but:
  - LO lacks a predictive normalization;
  - often does not include all partonic processes (for example,  $gg \rightarrow W + 1 \text{ jet}$  enters at NLO only).
- We would like to use NLO predictions throughout, but the current state of the art in this area is limited to  $W + 2 \text{ jets}$  and  $W + b\bar{b}$  (MCFM, <http://mcfm.fnal.gov/>).

# CDF's 'Method 2'

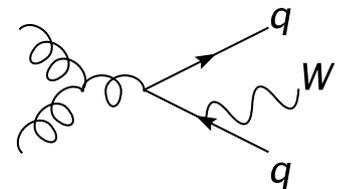
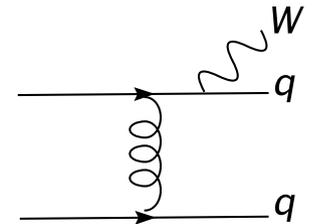
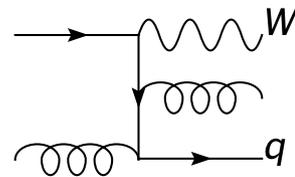
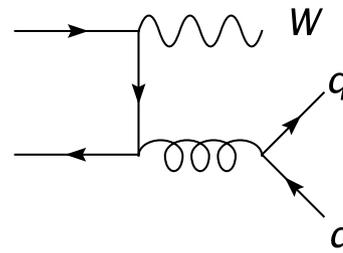
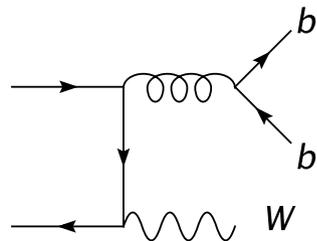
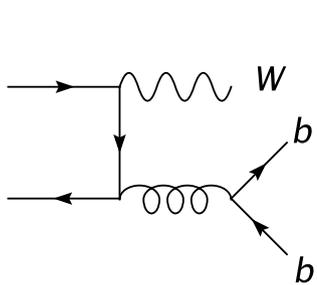
- To predict the number of  $W + b\bar{b} +$  jet events, CDF uses a mix of theory and data.
- They use ALPGEN (leading order) + Herwig to estimate the fraction of  $W + n$  jet events that contain two  $b$ 's.
- The prediction for the  $W + b\bar{b} +$  jets cross-section is then obtained from:

$$\sigma(Wb\bar{b} + (n - 2) \text{ jets}) = \left[ \frac{\sigma(Wb\bar{b} + (n - 2) \text{ jets})}{\sigma(W + n \text{ jets})} \right]_{MC} \times [\sigma(W + n \text{ jets})]_{\text{data}}$$

- One would like to know how this ratio depends on:
  - the order in perturbation theory;
  - the choice of renormalization and factorization scales;
  - the number of jets,  $n$ .
- Can investigate some of these issues for  $n = 2$  using MCFM.

# $Wb\bar{b}$ vs. $W + 2 \text{ jets}$

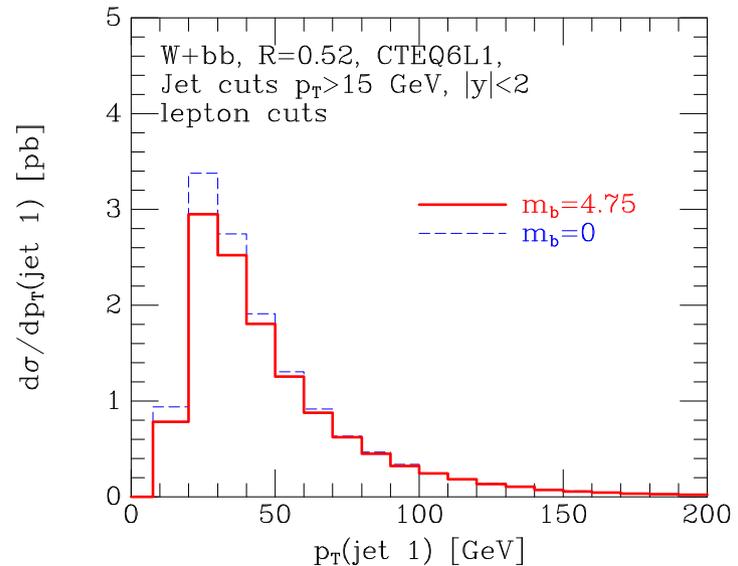
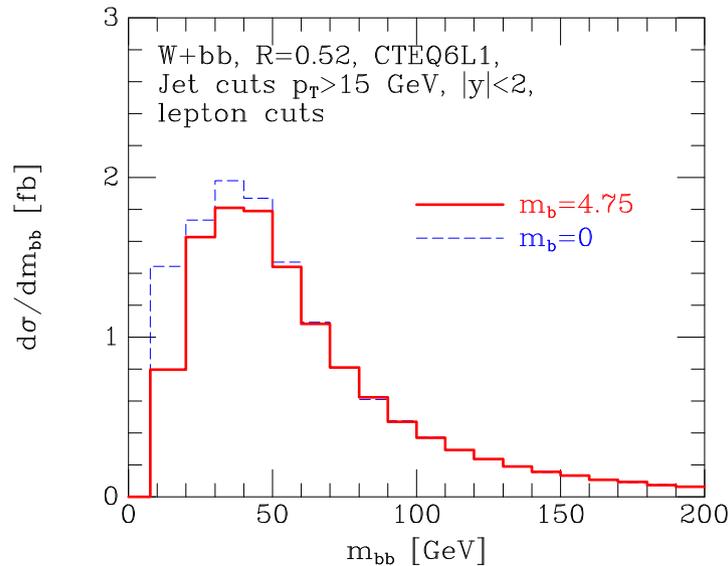
- Many more diagrams for  $W + 2 \text{ jets}$ .



- Notably,  $Wb\bar{b}$  has no gluon contribution at LO and  $b$ 's are produced by gluon splitting only.
- $b$  is treated as a massless particle in MCFM and the singularity protected by an invariant-mass cutoff.

# Mass effects

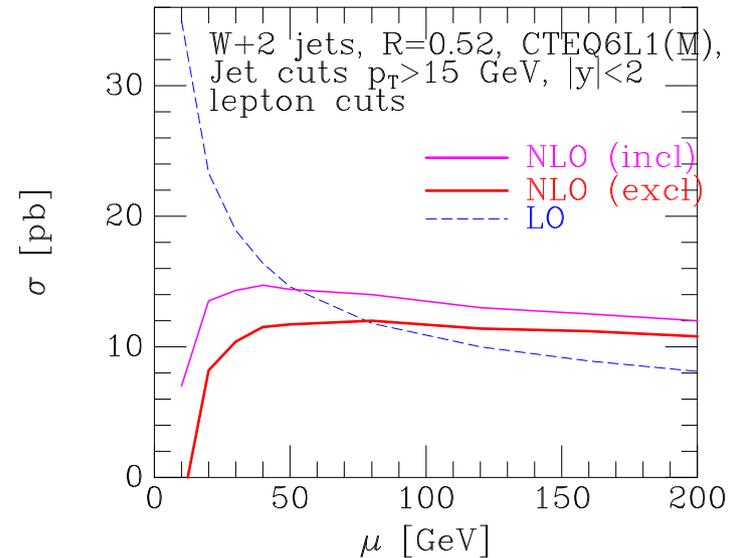
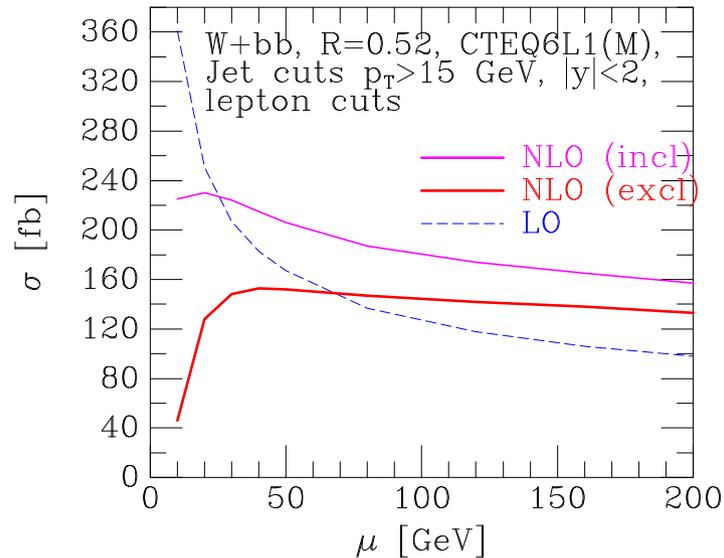
- Examine the effects of introducing the  $b$ -mass at lowest order, which is easily calculable.



- Overall the cross section decreases by approximately 10%. Kinematic distributions are not much affected away from regions of low  $p_T(b)$ .

# Scale dependence

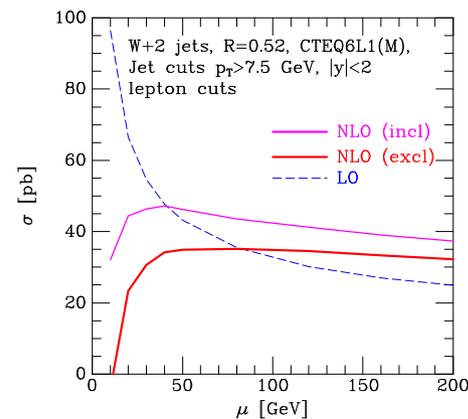
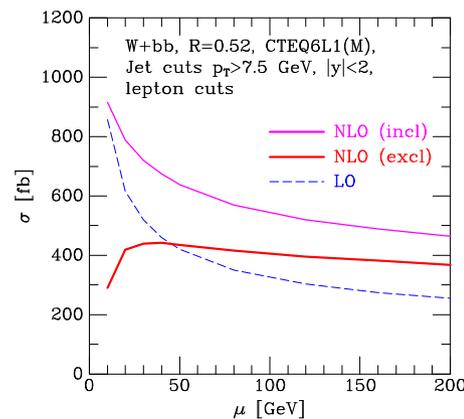
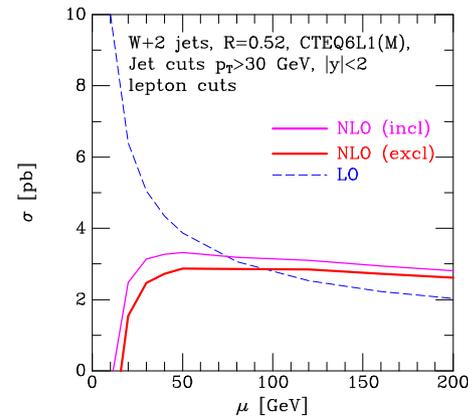
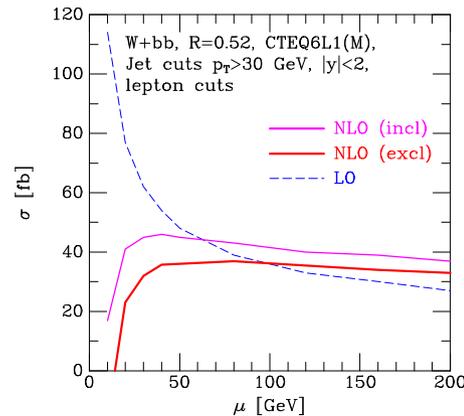
- Usual scale dependence, much reduced at NLO.



- Inclusive result (allows  $Wb\bar{b}j$ ,  $W + 3$  jet configurations) shows more scale dependence, as expected.
- Exclusive cross-sections stable over a large range of scales.

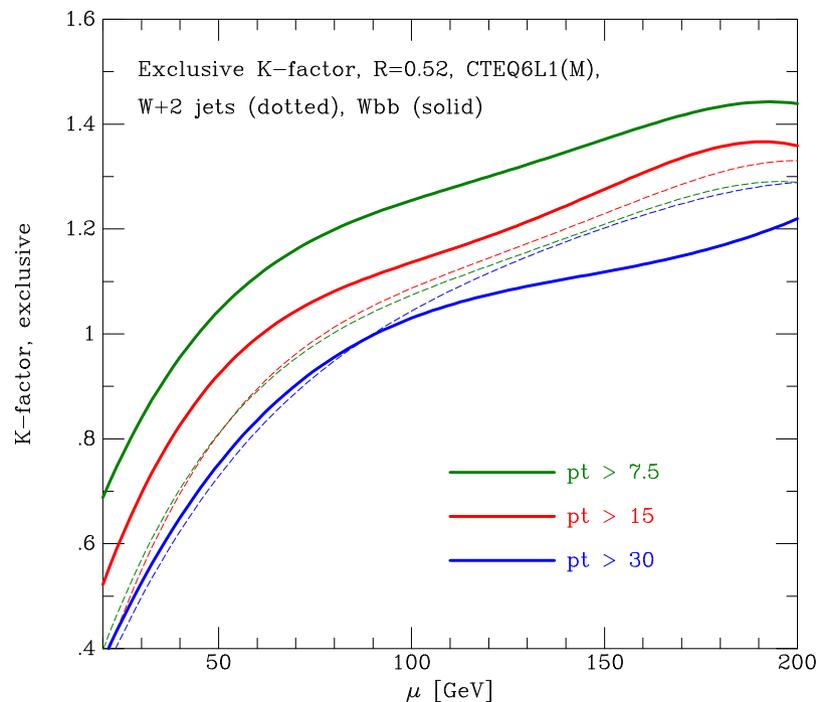
# Jet $p_T$ dependence

- Increasing the minimum jet  $p_T$  reduces the 3 jet contribution compared to the 2 jet one, so the behaviour of the inclusive cross-section improves.



# Scale dependence of $K$ -factors

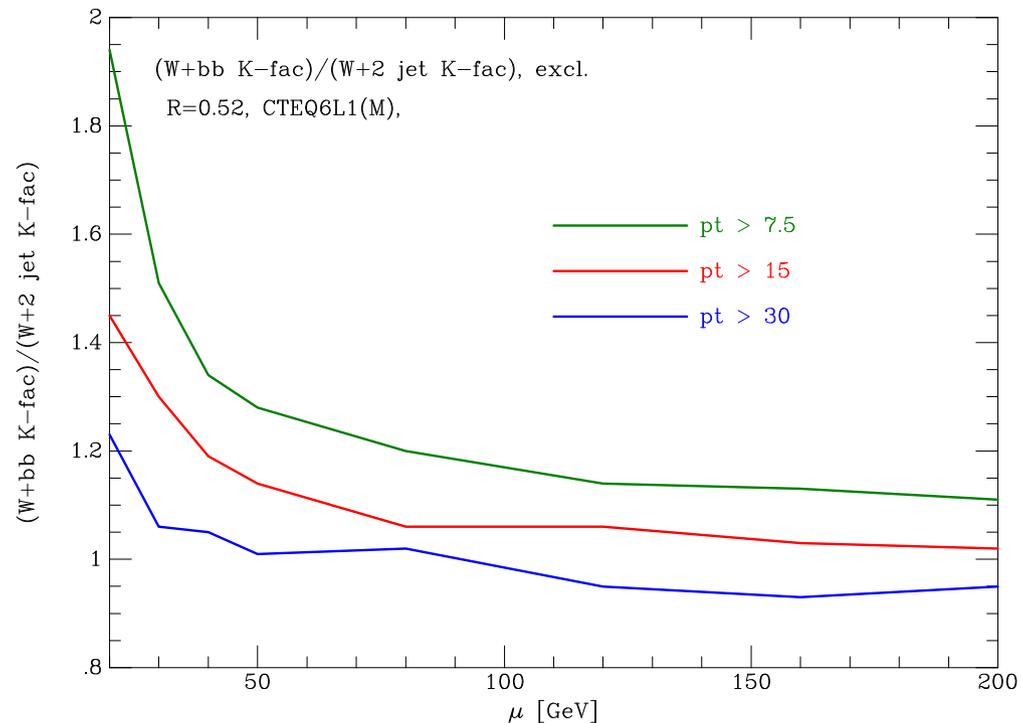
- Strong scale dependence.
- The  $Wb\bar{b}$   $K$ -factor varies greatly with the minimum jet  $p_T$ , whereas the  $W + 2$  jets one does not.



- Scale dependence has a similar shape for both processes.

# Reliability of Method 2 at NLO

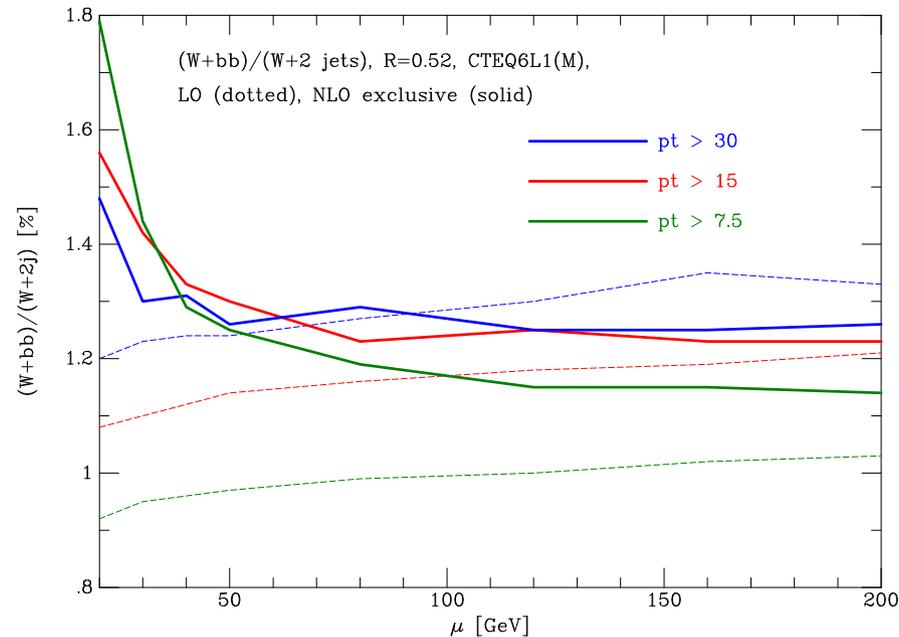
- If we are to trust Method 2, the ratio of  $K$ -factors should be  $\sim 1$ .



- This seems to be true for scale choices around 50 GeV or greater and  $p_T$  cuts of about 15 GeV or greater.
- Markedly worse for lower jet  $p_T$  cuts.

# *b*-jet fraction

- At NLO, ratio is stable across a wide range of scales.



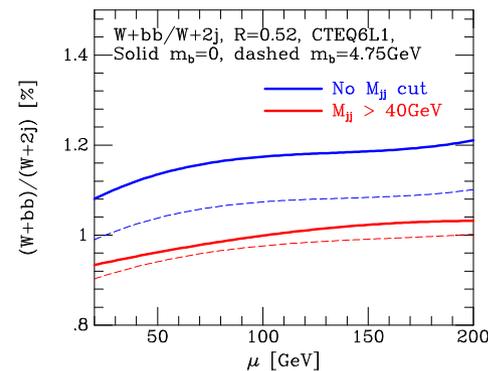
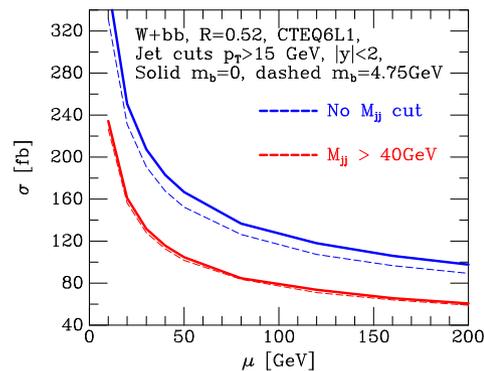
- For a  $p_T$  cut of 15 GeV and  $\mu \sim M_W$ , we have:

$$\left[ \frac{\sigma(Wb\bar{b})}{\sigma(W + 2 \text{ jets})} \right]_{LO} = 1.16\%,$$

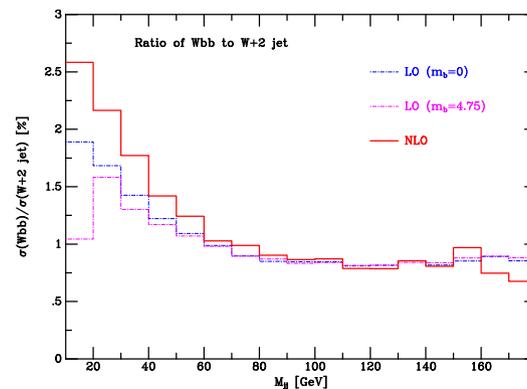
$$\left[ \frac{\sigma(Wb\bar{b})}{\sigma(W + 2 \text{ jets})} \right]_{NLO} = 1.23\%$$

# $b\bar{b}$ mass cut

- Such a cut would be helpful, if it could be experimentally enforced:
  - It improves the massless approximation

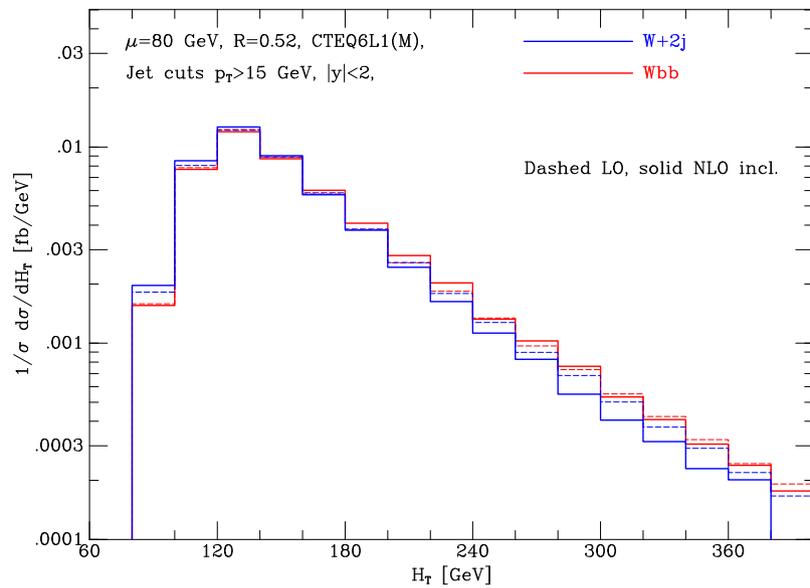
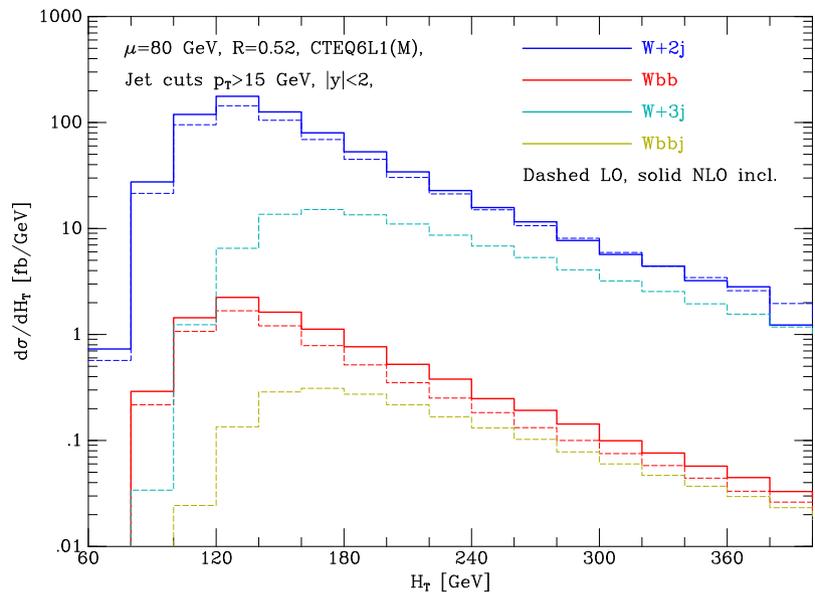


- It reduces this background compared to, for example,  $t\bar{t}$  production, since here the  $b$ 's like to lie at low invariant mass.



# Kinematic distributions

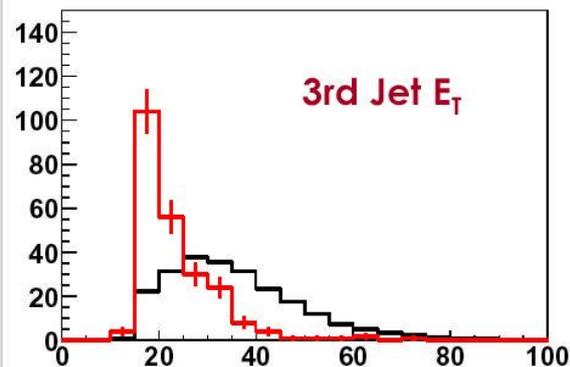
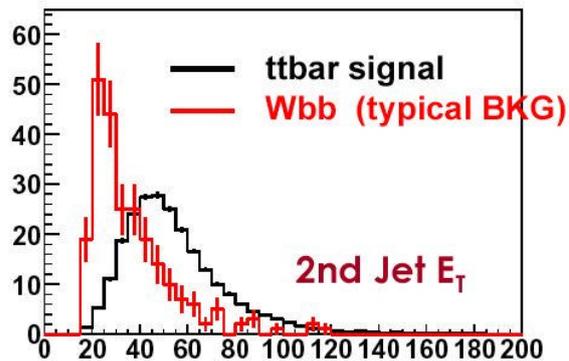
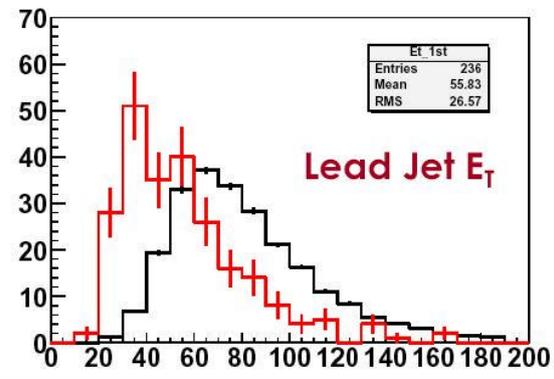
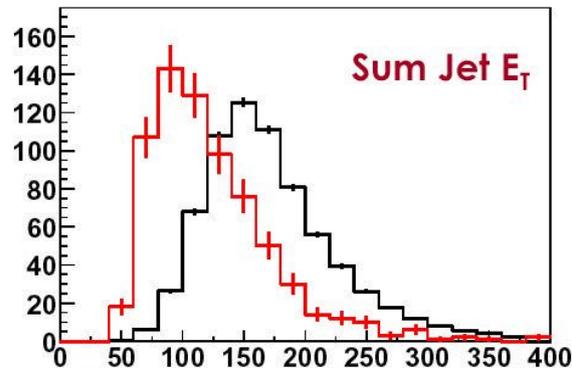
- NLO behaviour may provide clues to processes with more jets, especially for very inclusive variables such as  $\sum E_T(\text{jet})$  and  $H_T = \sum_{\text{event}} E_T$ .



- $Wb\bar{b}$  shape is relatively unchanged at NLO, compared to  $W + 2$  jets.

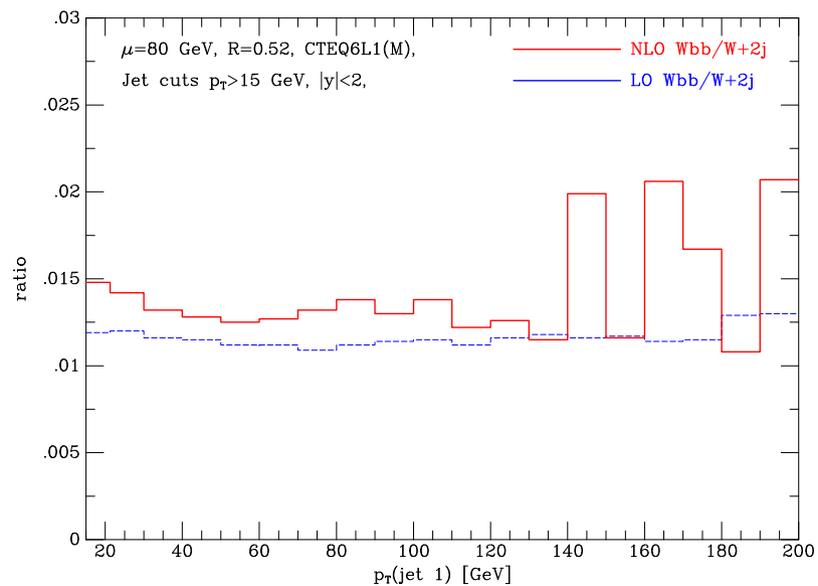
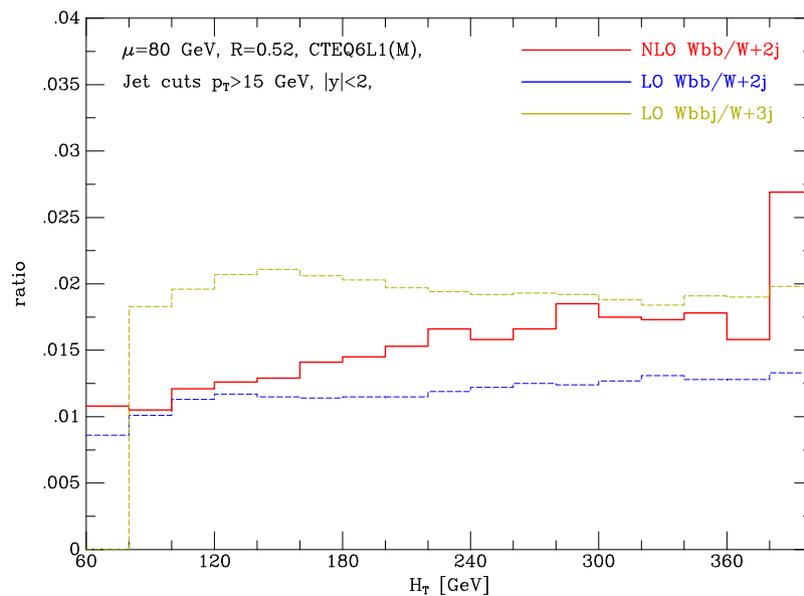
# Shape comparisons

- Top analysis, which would like to make kinematic cuts to reduce the  $W + \text{jet}$  backgrounds, relies on similar shapes of kinematic distributions.



# NLO predictions

- Statistics limited at the moment, but evidence for a change of shape in the  $H_T$  distribution.

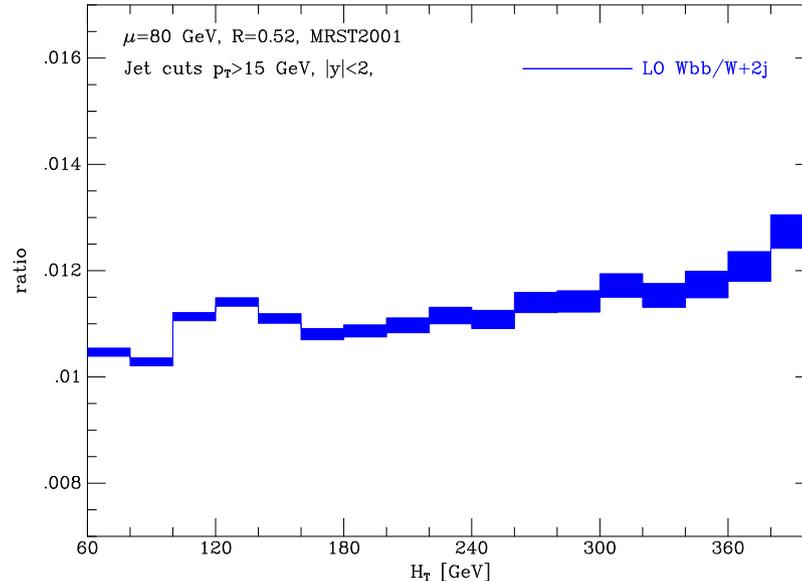


- Does not seem to be due to extra parton flux.
- Effect on analyses still being studied.

# PDF uncertainties

- Implemented in MCFM using LHAPDF (see <http://pdf.fnal.gov/>).
- Beta version of the LHAPDF code uses grid versions of the PDF's so that cross-sections using all PDF error sets can be calculated in one Monte Carlo run.
- At present, not all PDF sets are implemented in this form.

The uncertainty on the total  $Wb\bar{b}$  cross-section is  $\pm 1.5\%$ .



# Outlook

- The  $W + \text{jets}$  channel is very important for many physics searches in Run II and should be understood to the best of our ability.
- There should be lots to learn from the NLO corrections that we know about, i.e.  $Wb\bar{b}$  and  $W + 2 \text{ jets}$ .
- Too early to predict the effects of these results on Run II analyses, but preliminary results suggest that we can proceed with more confidence.
- Inclusion of PDF errors in the Monte Carlo (and LHAPDF) is a step in the right direction.
- Comparisons with parton shower approaches and data should be coming soon.