

Mu2e: Search for Muon to Electron Conversion at Fermilab

Eric Prebys
Fermilab

For the Mu2e Collaboration



Mu2e Collaboration



Boston University

Brookhaven National Laboratory

University of California, Berkeley

University of California, Irvine

California Institute of Technology

City University of New York

Duke University

Fermilab

University of Houston

University of Illinois, Urbana-Champaign

University of Massachusetts, Amherst

Lawrence Berkeley National Laboratory

Lewis University

Muons, Inc.

Northern Illinois University

Northwestern University

Pacific Northwest National Laboratory

Purdue University

Rice University

University of Virginia

University of Washington, Seattle



Istituto G. Marconi Roma

Laboratori Nazionale di Frascati

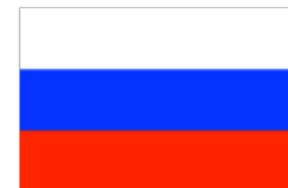
INFN Genoa

Università di Pisa, Pisa

INFN Lecce and Università del Salento

Gruppo Collegato di Udine

currently 155 collaborators
28 institutions

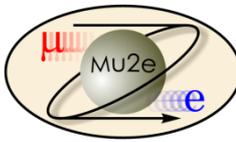


*Institute for Nuclear
Research, Moscow, Russia*

JINR, Dubna, Russia



Outline

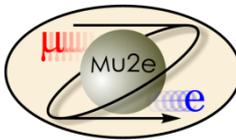


- Theoretical Motivation
- Experimental Technique
- Making Mu2e work at Fermilab
- Sensitivities
- Future Upgrades
- Conclusion

Will spend quite a bit of time on this



Provocative Comments



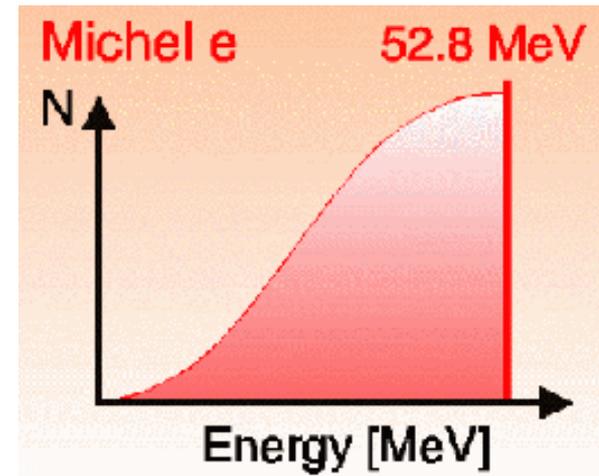
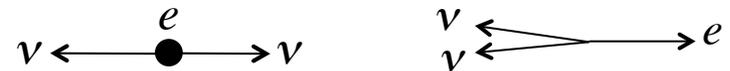
- Once upon a time, high energy physics moved forward by going to higher energies and “seeing what came out”.
 - The last time this happened was the discovery of the tau lepton and b quark in the 70s!
- For the last 40 years, all other discoveries have been preceded by strong indirect evidence
 - $K \rightarrow \mu^+ \mu^-$ suppression \rightarrow charm quark
 - CP Violation \rightarrow third generation
 - Weak decays \rightarrow W and Z particles and their masses
 - Precision tests at LEP and elsewhere \rightarrow top and Higgs masses
- With the discovery of the Higgs, we now find ourselves without guidance for the first time in half a century
 - The LHC was “guaranteed” to discover the Higgs (or it would have been even more interesting)
 - No one knows the next “sure bet” energy!
- If the past is any indicator, such guidance will likely come from indirect evidence.



History of the Muon



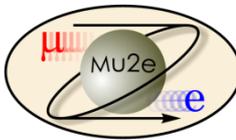
- The muon was originally discovered in 1936 by Anderson and Neddermeyer while studying cosmic ray data
- Hypothesized to be Yukawa's proposed mediator of the nuclear binding force, but did not interact strongly
 - Yukawa's particle was the pion
- Excited electron?
 - If so, expect $\mu \rightarrow e + \gamma$
 - **Not seen!**
- The muon was observed to decay to electron+"something invisible" with a spectrum consistent with a three body decay



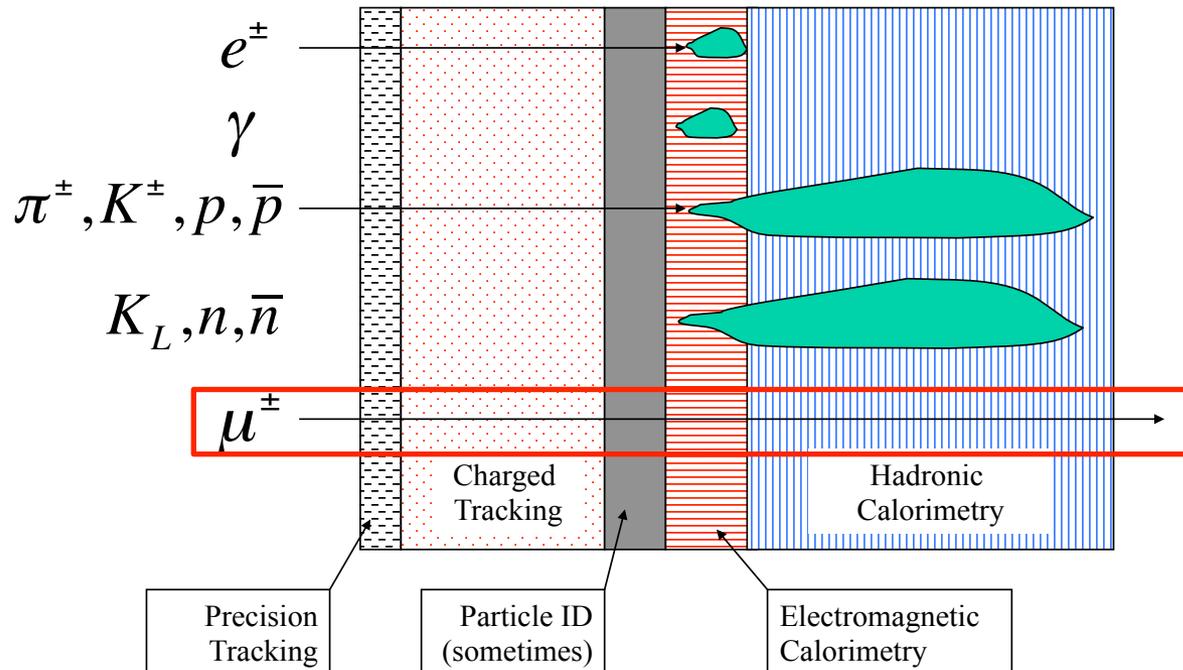
Fast forwarding (and skipping a whole bunch of stuff)...



Today's Muon



- Mass: $105.66 \text{ MeV}/c^2$ ($\sim 200m_e \sim 0.1m_p$)
- Charge: $\pm e$
- Spin: $\frac{1}{2}\hbar$ (fermion)
- Lifetime: $2.2 \mu\text{sec}$ ($c\tau=660\text{m}$)
- Interactions: Electromagnetic and Weak, but NOT strong
- Because muons are so much heavier than electrons, they are very penetrating





The Standard Model

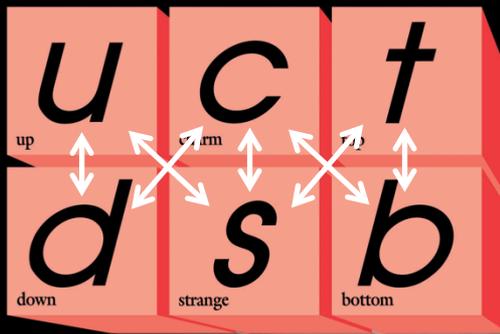


Spin 1/2 "Fermions"

Spin 1 "Bosons"

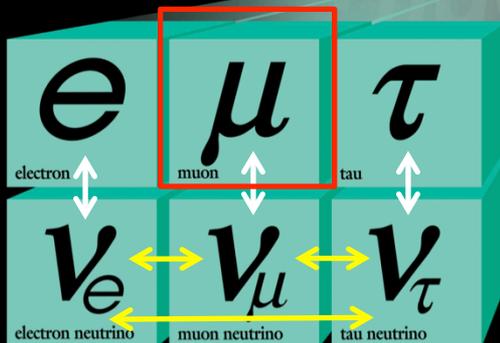
Combine to form hadrons

Quarks



Quarks can transition across generations

Leptons transition within generations...

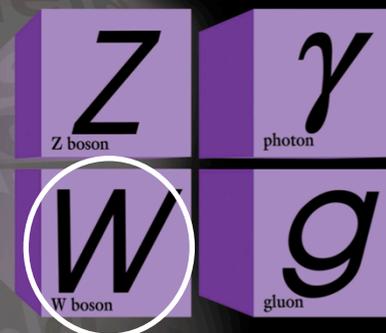


Free

Leptons

...except for neutrino mixing

Forces

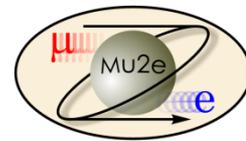


Mediate interactions

Weak charged current (W[±]) interactions "flip" fundamental fermions in weak isospin space

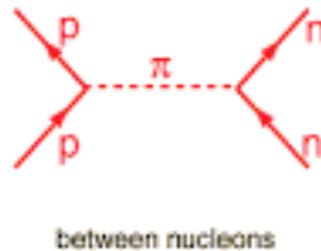
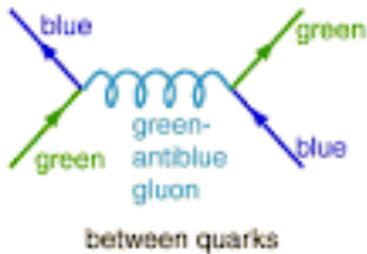
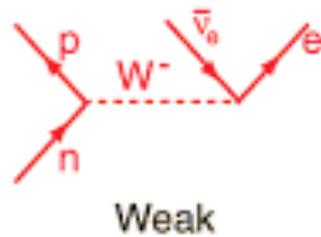
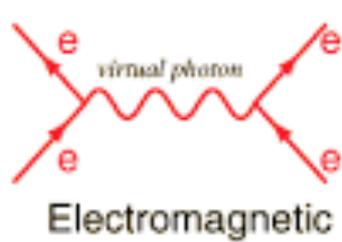


Interactions in the Standard Model

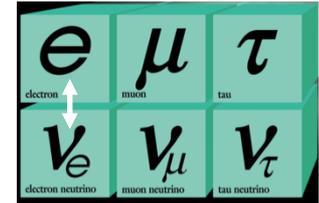
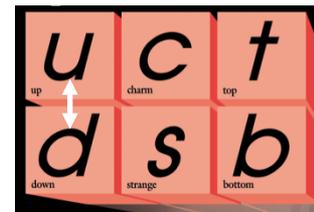
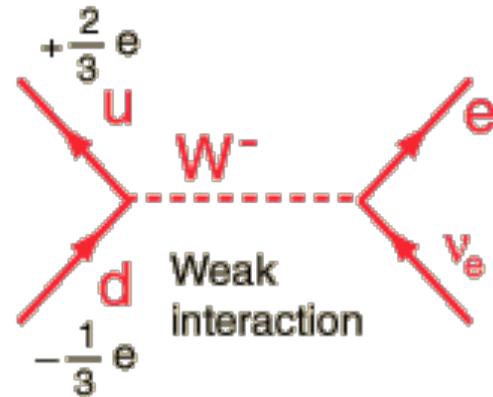


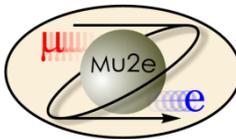
The bosons mediate interactions between the fundamental fermions

W particle causes a weak isospin transition within one *weak* quark or lepton generation



Strong Interaction





Generation (Flavor) Transitions

- In both the quark and lepton sector, the weak eigenstates are related to the mass eigenstates by a unitary matrix

$$\begin{bmatrix} d' & s' & b' \end{bmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{bmatrix} d \\ s \\ b \end{bmatrix}$$

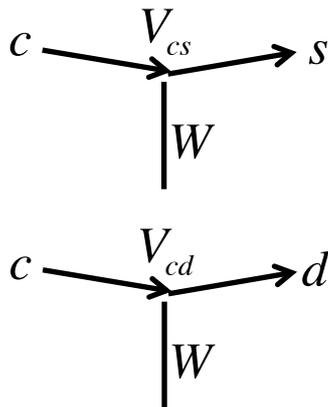
“almost” diagonal

$$\begin{bmatrix} \nu_e & \nu_\mu & \nu_\tau \end{bmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

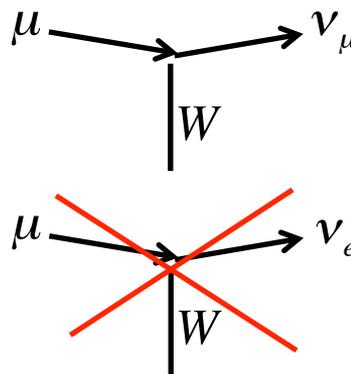
~maximum mixing

- However, because the neutrino masses and their differences are so small, the phenomenology is *very* different

Quarks: generational transitions observed

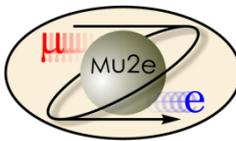


Leptons: weak transitions and mixing proceed separately



← Pure weak state. Propagates as a superposition of mass eigenstates → “neutrino mixing”

NOT observed!

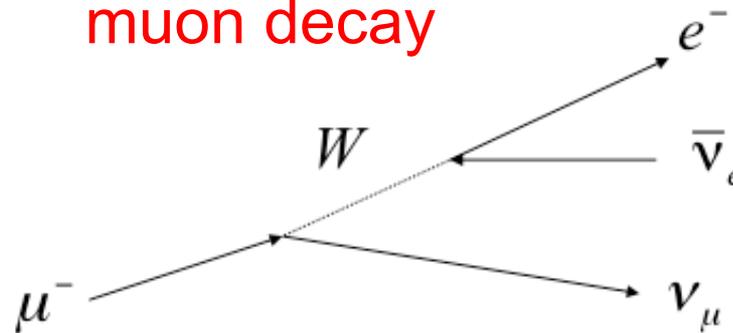


Lepton Number and Lepton Flavor Number

As a consequence, both lepton number and lepton "flavor" (generation) number are individually conserved*

	l	l_e	l_μ
μ^-	1	0	1
total	1	0	1

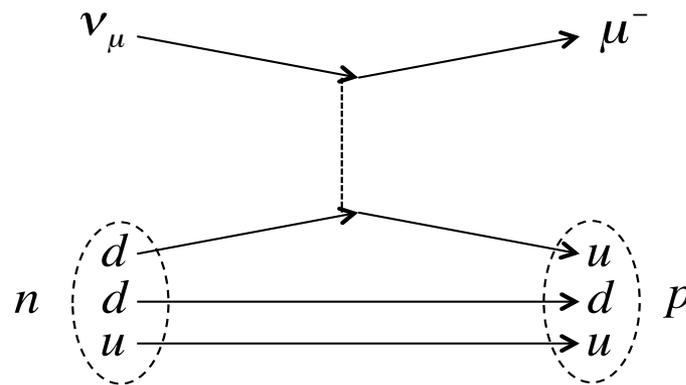
muon decay



	l	l_e	l_μ
e^-	1	1	0
$\bar{\nu}_e$	-1	-1	0
ν_μ	1	0	1
total	1	0	1

	l	l_e	l_μ
ν_μ	1	0	1
n	0	0	0
total	1	0	1

$\nu + n \rightarrow \mu^- + p$

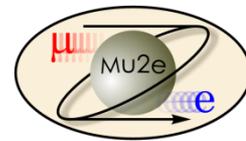


	l	l_e	l_μ
μ^-	1	0	1
p	0	0	0
total	1	0	1

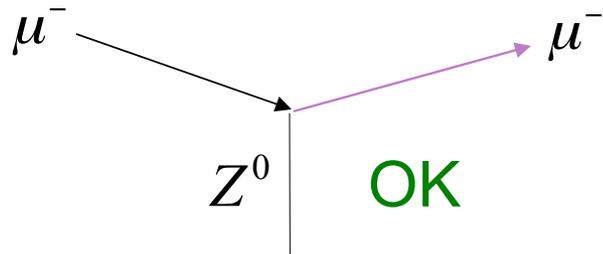
*except in neutrino mixing



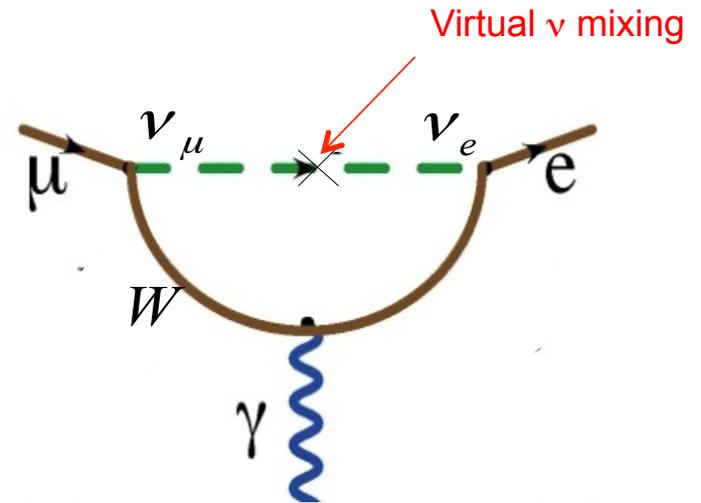
Charged Lepton Flavor Violation (CLFV)



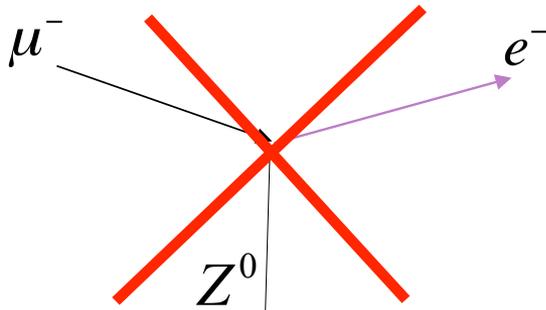
The Z^0 mediates neutral current scattering



Note: Observation of neutrino mixing shows CLFV *can* occur



However, “Flavor Changing Neutral Currents” (FCNC):

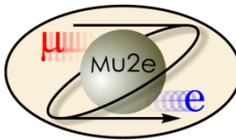


However, the Standard Model branching ratio is $\sim \mathcal{O}(10^{-52})$
(35 orders of magnitude below our goal)

are forbidden in Standard Model



Beyond the Standard Model



- Because extensions to the Standard Model couple the lepton and quark sectors, Charged Lepton Flavor Violation (CLFV) is a nearly universal feature of such models.
- The fact that it has not yet been observed already places strong constraints on these models.
- CLFV is a powerful probe of multi-TeV scale dynamics
 - complementary to direct collider searches
- Among various possible CLFV modes, rare muon processes offer the best combination of new physics reach and experimental sensitivity

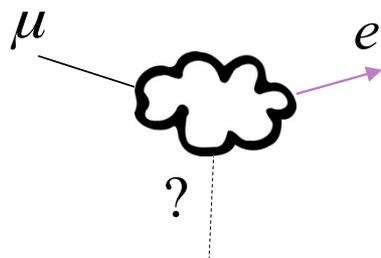


Generic Beyond Standard Model CLFV



There are two broad classes of CLFV reactions...

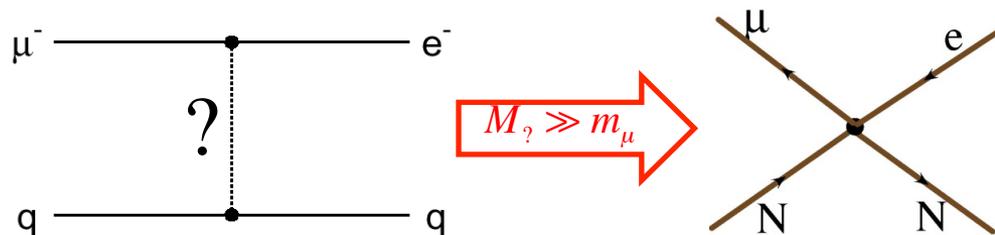
Flavor Changing Neutral Current



➤ Mediated by *virtual* massive neutral Boson, e.g.

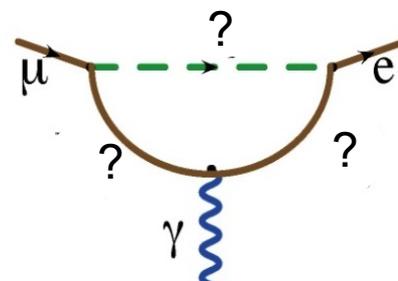
- Leptoquark
- Z'
- Composite

➤ Approximated by “four fermi interaction”

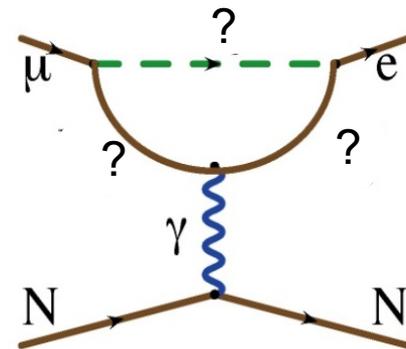


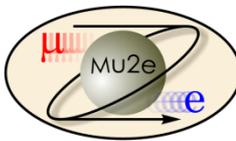
Dipole (penguin)

➤ Can involve a real photon



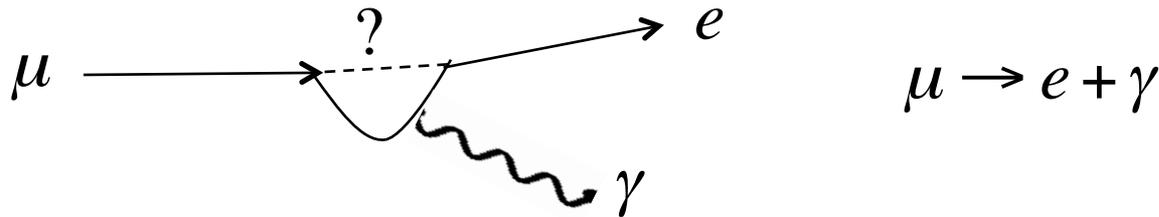
➤ Or a virtual photon



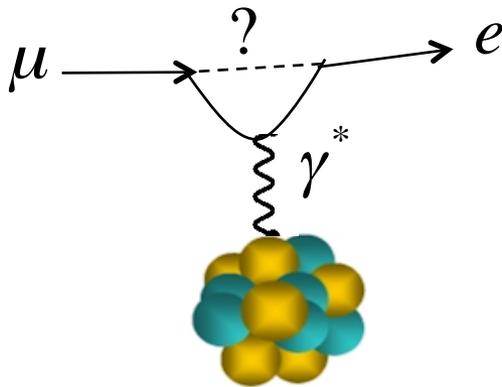


Decay vs. Conversion

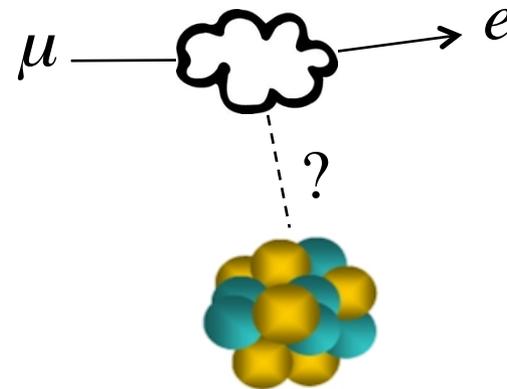
- Only the “dipole”-like reactions can lead to a decay



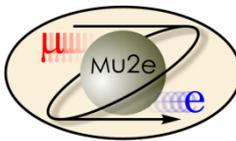
- However, if we capture a μ^- on a nucleus, it could “convert” to an e^- via exchange of a virtual particle in both scenarios



photon



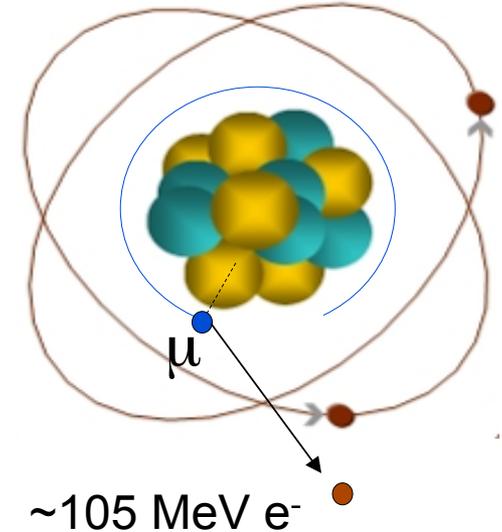
heavy neutral boson



Experimental Signature of $\mu+N \rightarrow e+N$

- When captured by a nucleus, a muon will have an enhanced probability of exchanging a virtual particle with the nucleus.
- This reaction recoils against the entire nucleus, producing a *mono-energetic* electron carrying most of the muon rest energy

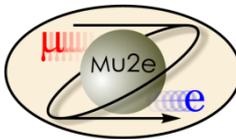
$$E_e = m_\mu c^2 - \frac{(m_e c^2)^2}{2m_N c^2}$$



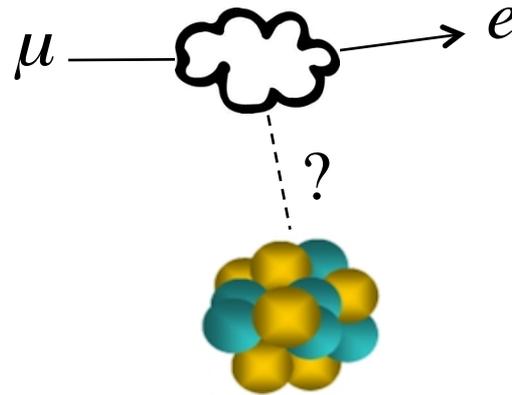
- Similar to $\mu \rightarrow e\gamma$, with important advantages:
 - No combinatorial background.
 - Because the virtual particle can be a photon or heavy neutral boson, this reaction is sensitive to a broader range of new physics.
- Relative rate of $\mu \rightarrow e\gamma$ and $\mu N \rightarrow eN$ would be the most important clue regarding the details of the physics



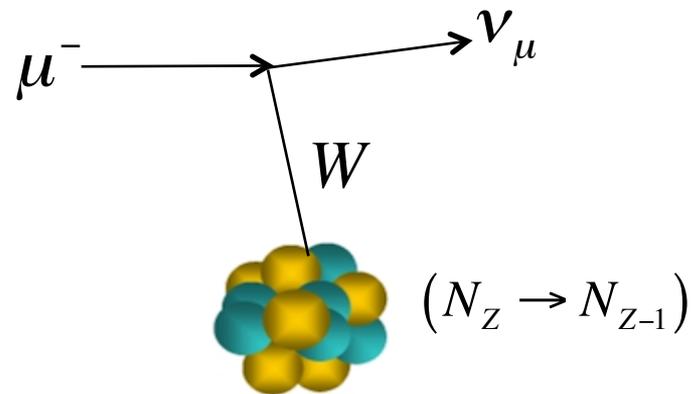
What We (Plan to) Measure



- We will measure the rate of μ to e conversion...



...relative to ordinary μ capture

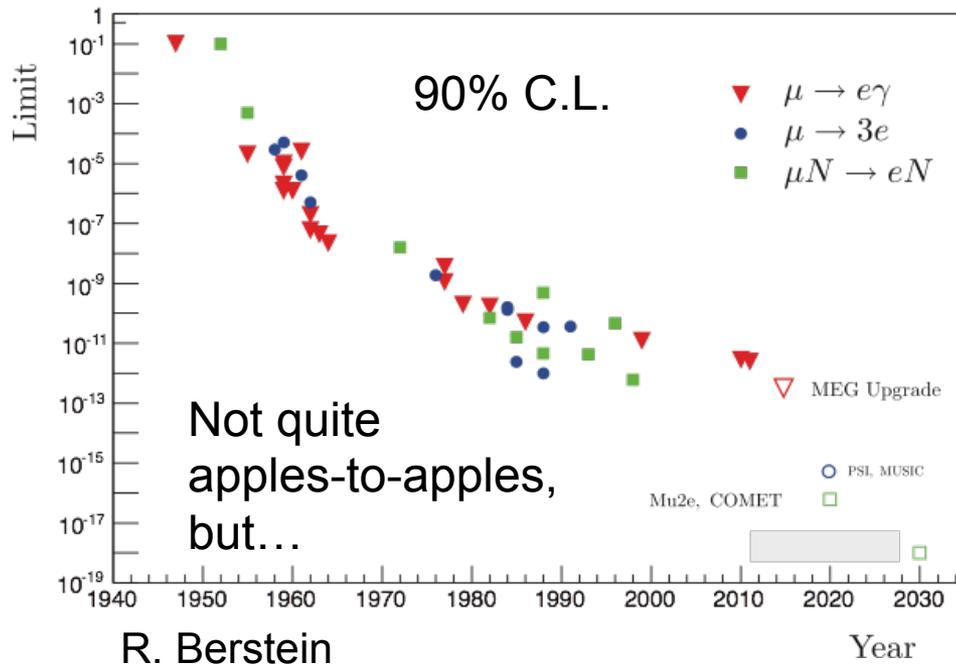
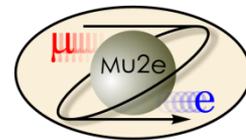


- This is defined as

$$R_{\mu e} \equiv \frac{\Gamma(\mu^- N(A, Z) \rightarrow e^- + N(A, Z))}{\Gamma(\mu^- N(A, Z) \rightarrow \nu_\mu + N'(A, Z-1))}$$



History of Lepton Flavor Violation Searches



➤ Best Limits (all from PSI)

- $\text{Br}(\mu \rightarrow e\gamma) < 6 \times 10^{-13}$ (MEG 2013)
- $\text{Br}(\mu \rightarrow 3e) < 1 \times 10^{-12}$ (Sindrum-I 1988)
- $R_{\mu e} < 7 \times 10^{-13}$ (Sindrum-II 2006)

Four orders of magnitude improvement!

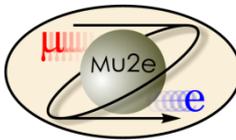
Mu2e will measure:

$$R_{\mu e} \equiv \frac{\Gamma(\mu^- N(A, Z) \rightarrow e^- + N(A, Z))}{\Gamma(\mu^- N(A, Z) \rightarrow \nu_\mu + N'(A, Z-1))}$$

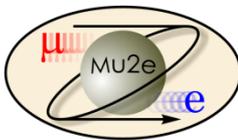
Goal: single event sensitivity of $R_{\mu e} = 3 \times 10^{-17}$



Just to be clear...



- We are *not* planning to make a measurement and compare it to a calculation.
- We are looking for something that (effectively) doesn't exist in the Standard Model.
- Our goal is to build an experiment with negligible backgrounds, such that any observed signal will be *unambiguous evidence of new physics*.
- We are planning for an improvement of roughly four orders of magnitude in sensitivity over the best previous measurement.

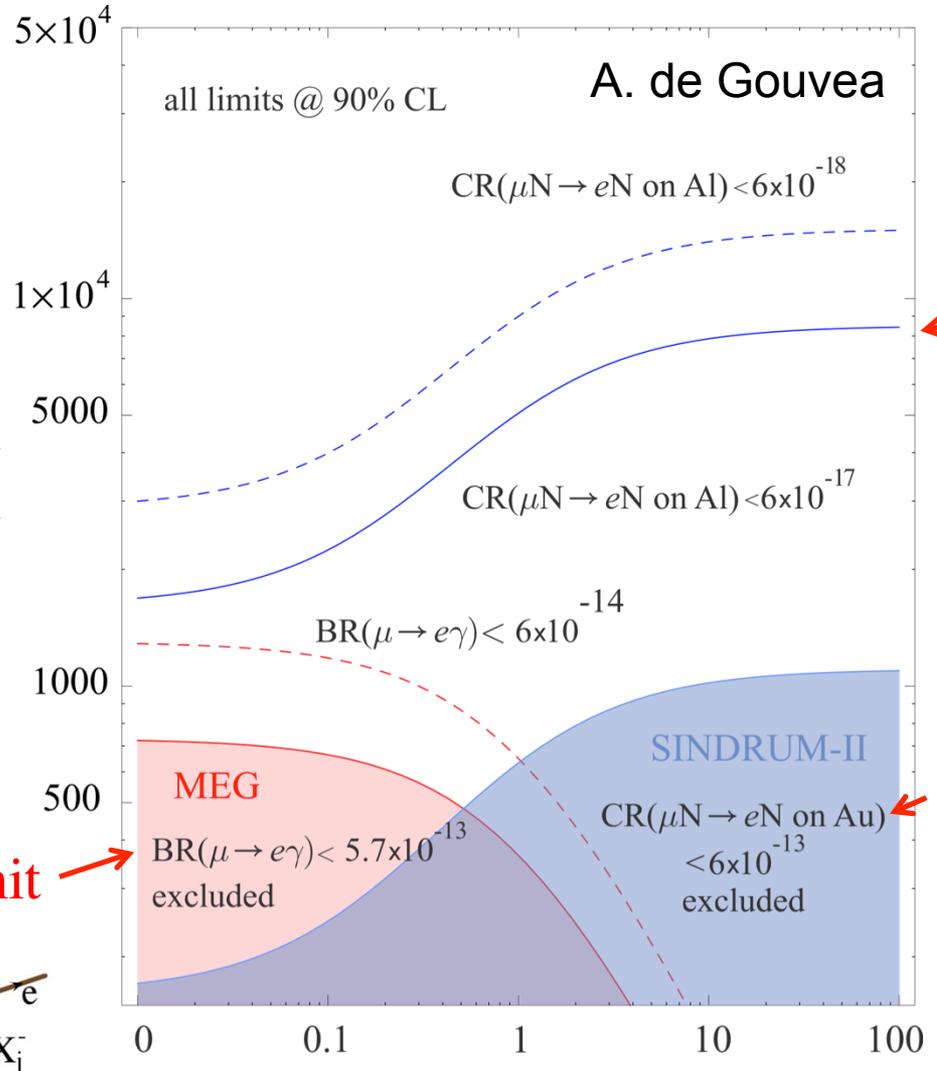


Dipole vs. Contact Reaction

Mass Scale

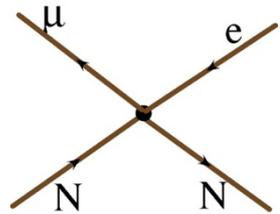
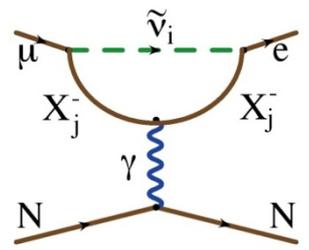
$$\left(\text{Rate} \propto \frac{1}{\Lambda^4} \right)$$

Λ (TeV)



Best $\mu \rightarrow e\gamma$ limit

Best $\mu N \rightarrow e N'$ limit



K
(different for different models)

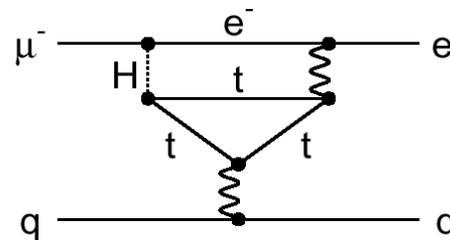
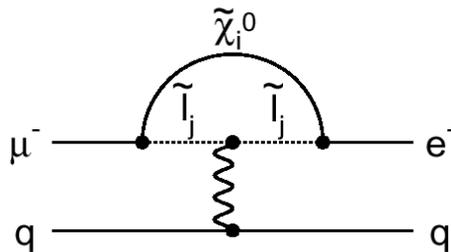


Example Sensitivities*



Supersymmetry

Predictions at 10^{-15}

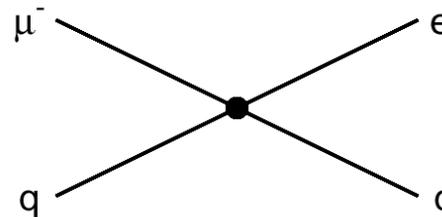
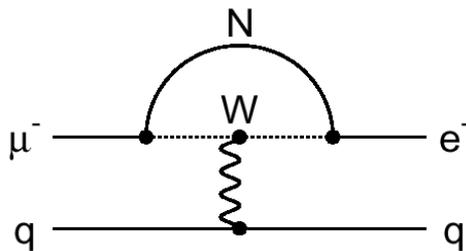


Second Higgs doublet

$$g_{H_{\mu e}} = 10^{-4} \times g_{H_{\mu\mu}}$$

Heavy Neutrinos

$$|U_{\mu N}^* U_{eN}|^2 = 8 \times 10^{-13}$$

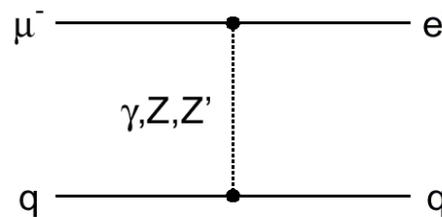
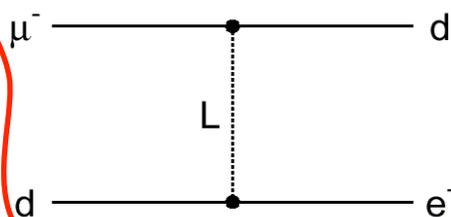


Compositeness

$$\Lambda_C = 3000 \text{ TeV}$$

Leptoquarks

$$M_L = 3000 \sqrt{\lambda_{\mu d} \lambda_{e d}} \text{ TeV}/c^2$$



Heavy Z' ,
Anomalous Z
coupling

$$M_{Z'} = 3000 \text{ TeV}/c^2$$
$$B(Z \rightarrow \mu e) < 10^{-17}$$

*After W. Marciano

No $\mu \rightarrow e\gamma$ signal



Example: $\mu \rightarrow e$ in Supersymmetry*



← SUSY Models

	AC	RVV2	AKM	δ LL	FBMSSM	LHT	RS
$D^0 - \bar{D}^0$	★★★	★	★	★	★	★★★	?
ϵ_K	★	★★★	★★★	★	★	★★	★★★
$S_{\psi\psi}$	★★★	★★★	★★★	★	★	★★★	★★★
$S_{\phi K_S}$	★★★	★★	★	★★★	★★★	★	?
$A_{CP}(B \rightarrow X_s \gamma)$	★	★	★	★★★	★★★	★	?
$A_{7\mu}(B \rightarrow K^* \mu^+ \mu^-)$	★	★	★	★★★	★★★	★★	?
$A_9(B \rightarrow K^* \mu^+ \mu^-)$	★	★	★	★	★	★	?
$B \rightarrow K^{(*)} \nu \bar{\nu}$	★	★	★	★	★	★	★
$B_s \rightarrow \mu^+ \mu^-$	★★★	★★★	★★★	★★★	★★★	★	★
$K^+ \rightarrow \pi^+ \nu \bar{\nu}$	★	★	★	★	★	★★★	★★★
$K_L \rightarrow \pi^0 \nu \bar{\nu}$	★	★	★	★	★	★★★	★★★
$\mu \rightarrow e \gamma$	★★★	★★★	★★★	★★★	★★★	★★★	★★★
$\tau \rightarrow \mu \gamma$	★★★	★★★	★	★★★	★★★	★★★	★★★
$\mu + N \rightarrow e + N$	★★★	★★★	★★★	★★★	★★★	★★★	★★★
d_n	★★★	★★★	★★★	★★	★★★	★	★★★
d_c	★★★	★★★	★★	★	★★★	★	★★★
$(g-2)_\mu$	★★★	★★★	★★	★★★	★★★	★	?

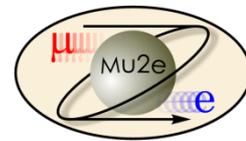
← All SUSY models predict both $\mu \rightarrow e \gamma$ and $\mu N \rightarrow e N$

Table 8: “DNA” of flavour physics effects for the most interesting observables in a selection of SUSY and non-SUSY models ★★★ signals large effects, ★★ visible but small effects and ★ implies that the given model does not predict sizable effects in that observable.

*from Altmannshofer, Buras, *et al*, Nucl.Phys.B830:17-94, 2010



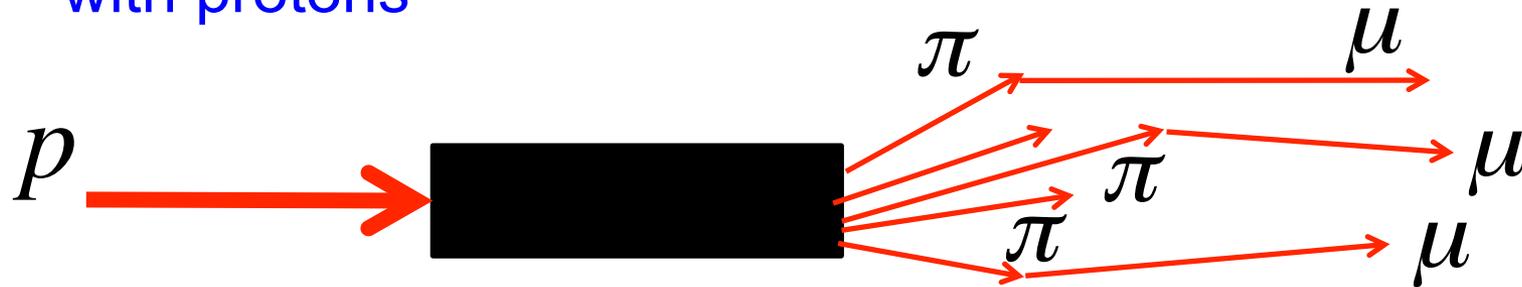
How do we make muons?



Hit a target with protons

This produces mostly pions

These quickly decay to muons

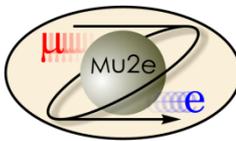


$$\tau_{\pi^\pm} = 26 \text{ ns}$$



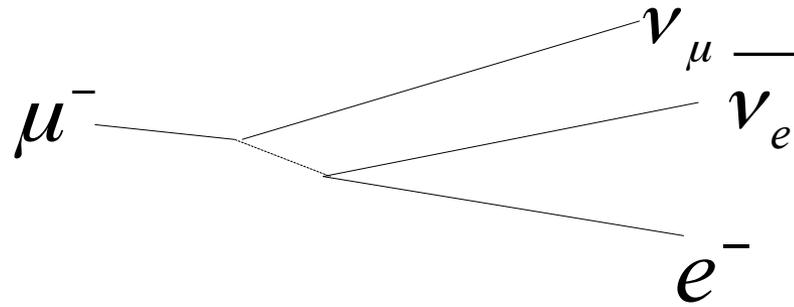
$$\tau_{\mu^\pm} = 2200 \text{ ns}$$

Muons go much further

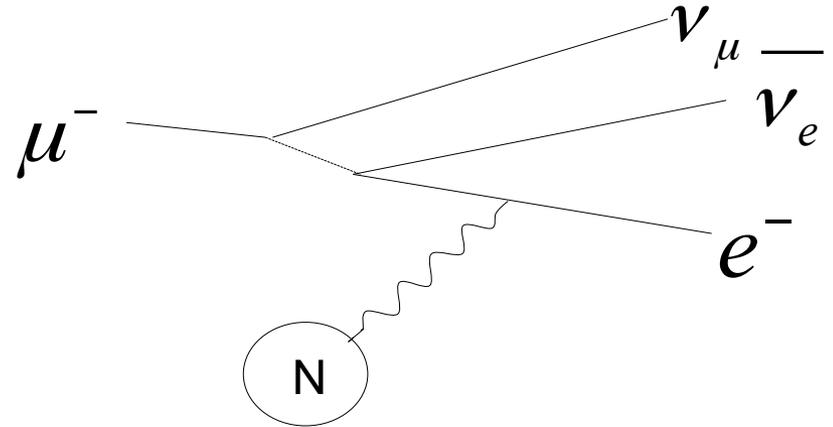


Biggest Issue: Decay in Orbit (DIO)

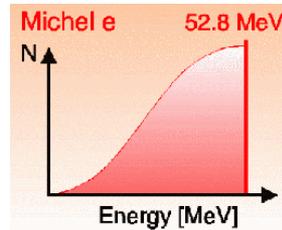
Free μ^- Decay:



Coherent DIO:



- Very high rate
- “Michel Spectrum”
 - Peak energy ~ 53 MeV



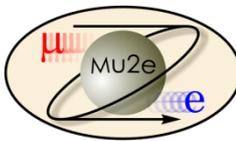
- Must design detector to be very *insensitive* to these.

- Nucleus coherently balances momentum and smears out Michel Spectrum.
- Rate approaches conversion (endpoint) energy as $\sim (E_{\text{conversion}} - E)^5$

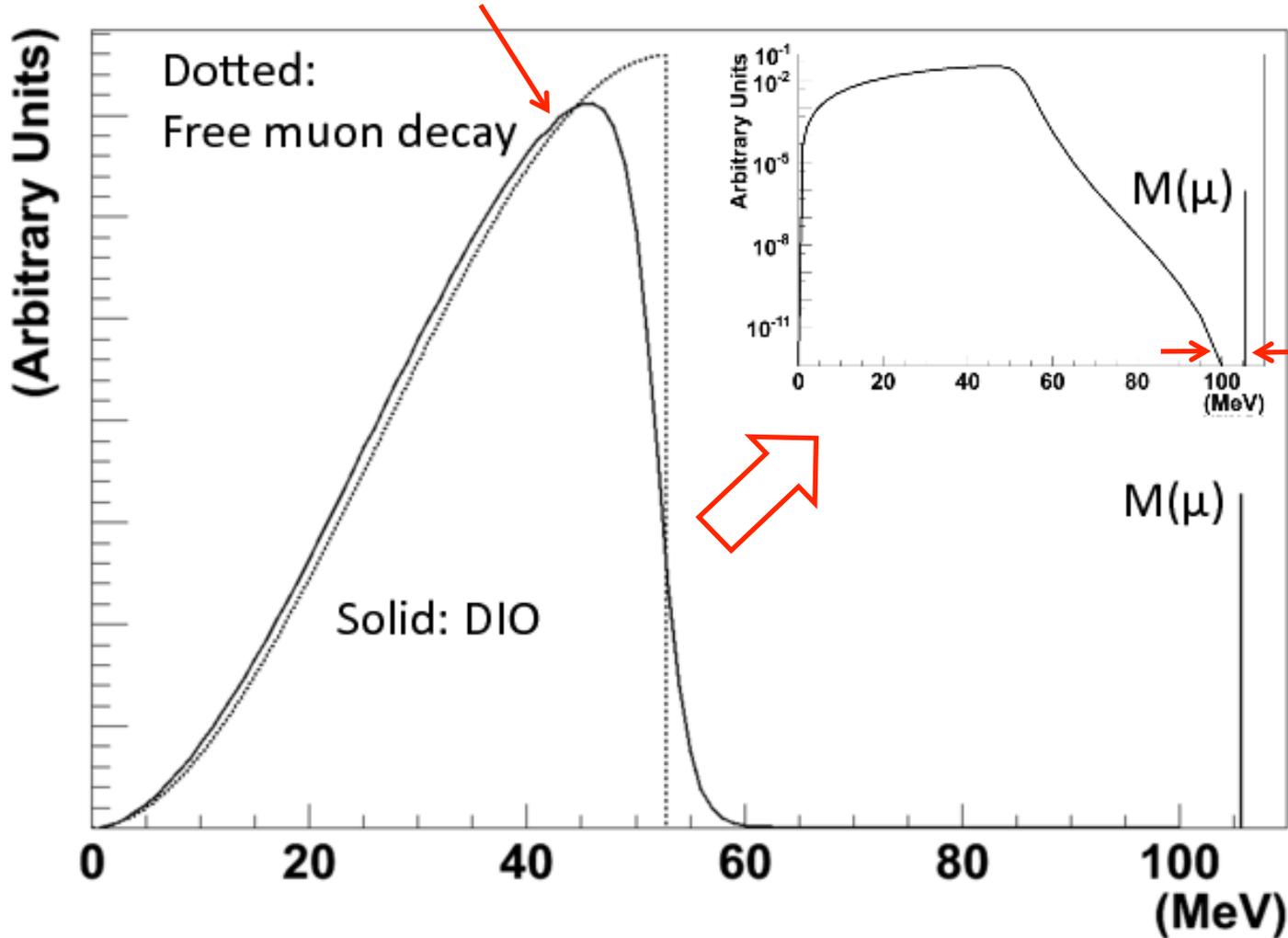
- Drives resolution requirement.



DIO Spectrum



We want to be blind to this
(acceptance)



We must
resolve this



Prompt Backgrounds



- There are significant backgrounds related to the production and transport of the muons.

- Radiative π^- capture
 $\pi^- N \rightarrow N^* \gamma, \gamma Z \rightarrow e^+ e^-$

Biggest worry

- Muon decay in flight

$$\mu^- \rightarrow e^- \nu \nu$$

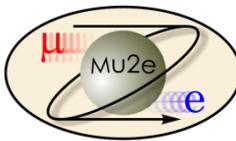
- Pion decay in flight

$$\pi^- \rightarrow e^- \nu_e$$

- Prompt electrons

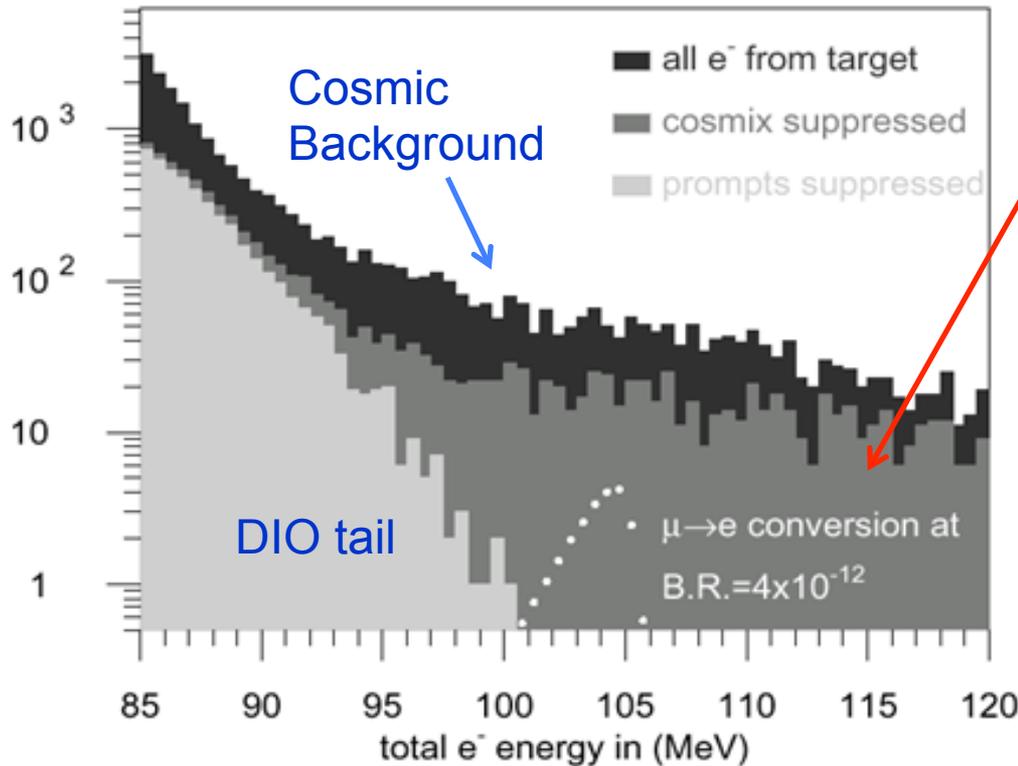
- General approach

- Produce muons
- Transport muons to target where some are captured.
- Wait(!) for prompt backgrounds to go away
- Open detection window to look for conversion of captured muons.



Experimental Challenge of “Waiting”

$\mu \rightarrow e$ Conversion: Sindrum II



- Most backgrounds are ~prompt with respect to the proton beam
 - Mostly radiative pion capture
- Previous experiments suppressed these backgrounds *by vetoing all observed electrons* for a period of time after the arrival of *each proton*.
 - This leads to a fundamental to a rate limitation.

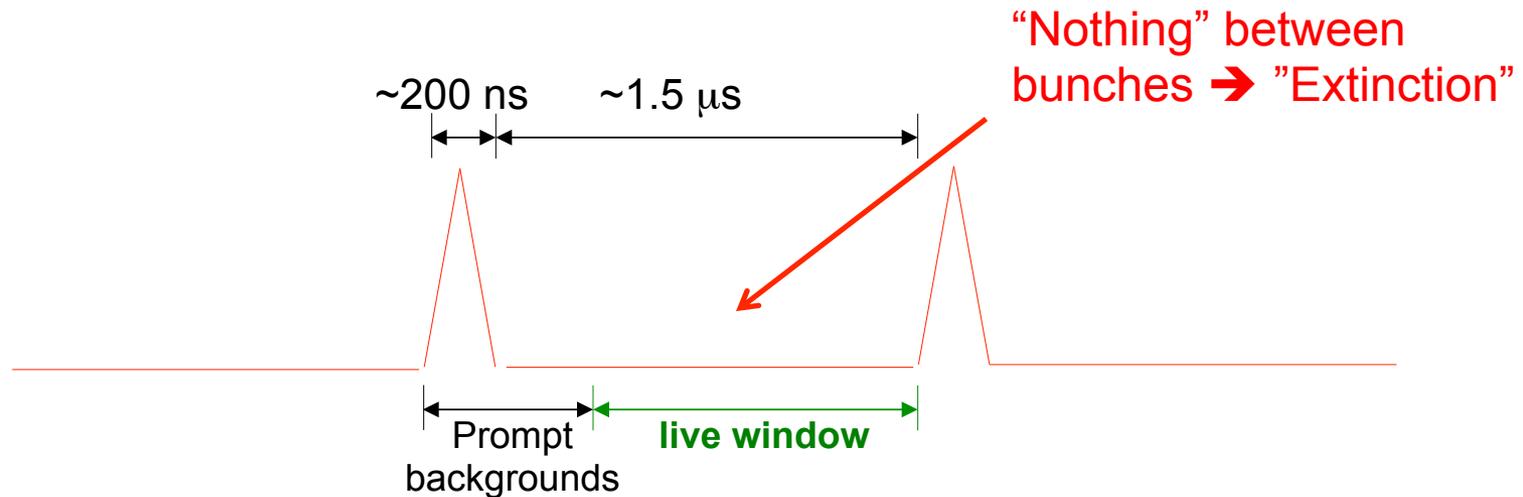
$$R_{\mu e} \equiv \frac{\Gamma(\mu^- Au \rightarrow e^- Au)}{\Gamma(\mu^- Au \rightarrow \text{capture})} < 7 \times 10^{-13}$$



Pulsed Beams (first proposed for MELC*)



- Replace individual protons with short proton *pulses*, separated by a time on the order of a muon life time.
- Veto the time after the pulse to eliminate prompt backgrounds.

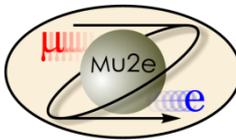


- Design a transport channel to optimize the transport of right-sign, low momentum muons from the production target to the muon capture target.
- Design a detector which is very insensitive to electrons from ordinary muon decays, and has excellent tracking resolution.

*1992, Moscow Meson Factory



Summary: Experimental Needs

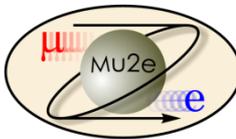


- Proton beam:
 - Bunches, separated by \sim muon lifetime with “nothing” in between them.
- Muon transport:
 - Optimize for low momentum, *negative* muons
- Detector:
 - Completely blind to any particle with $p \lesssim 60 \text{ MeV}/c$
 - Excellent energy resolution for $105 \text{ MeV } e^-$
 - \rightarrow Very low mass for both target and tracker!

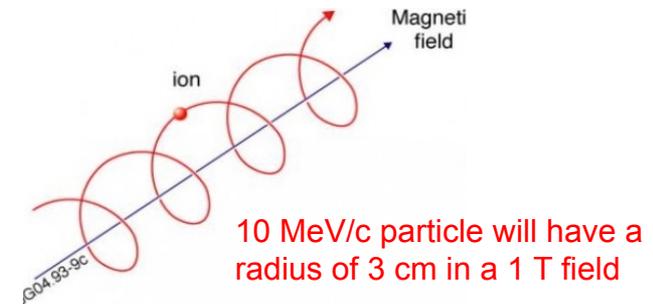
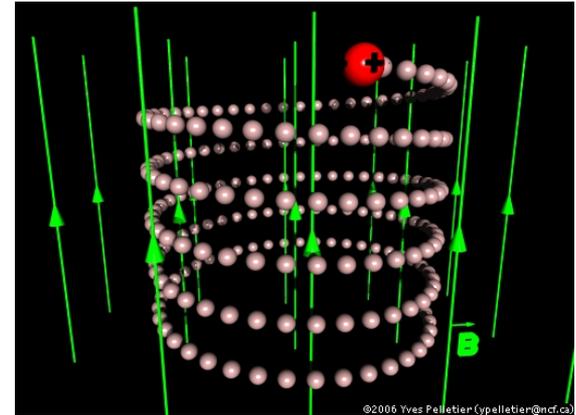
} Solenoids!



Refresher: Fun with Solenoids

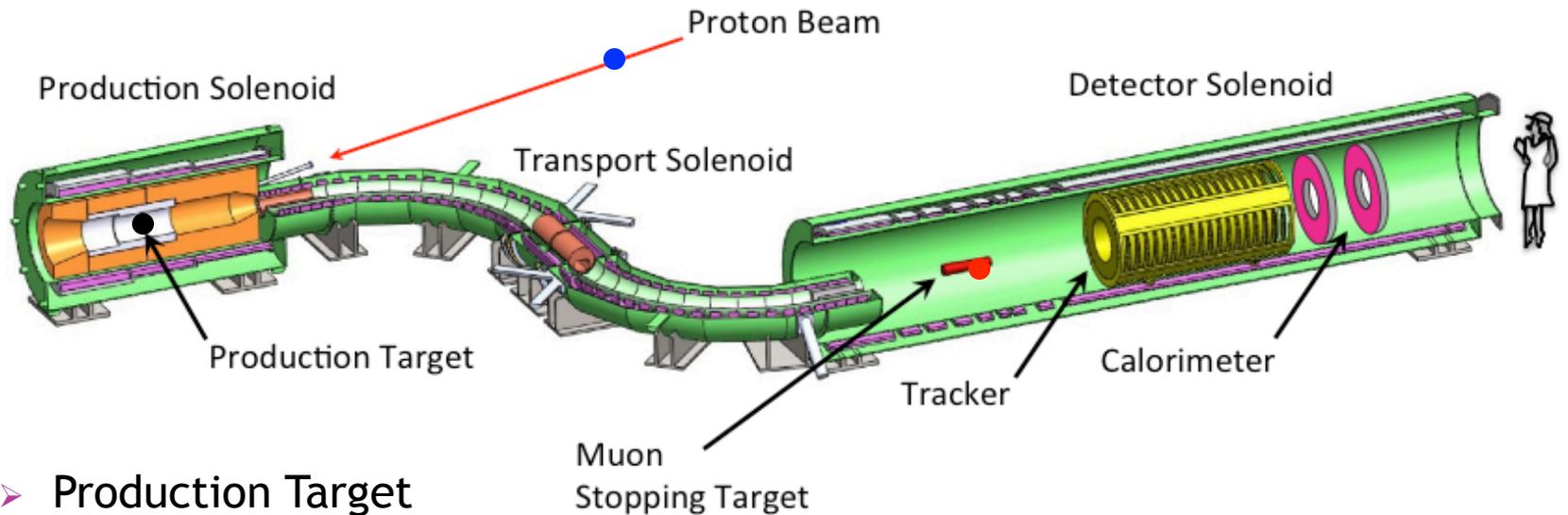
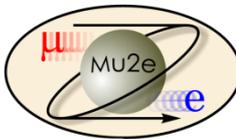


- Particles in a solenoidal field will generally move in a helical path
- Low momentum particles are effectively “trapped” along the field lines
 - We use this to transport muons
- A particle trapped along a *curved* solenoidal field will drift *out of the plane of curvature*
 - This is how we will resolve muon charge and momentum in the transport line
- For higher momentum particles, the curvature can be used to measure momentum
 - This is how we will measure the momentum of electrons from the capture target





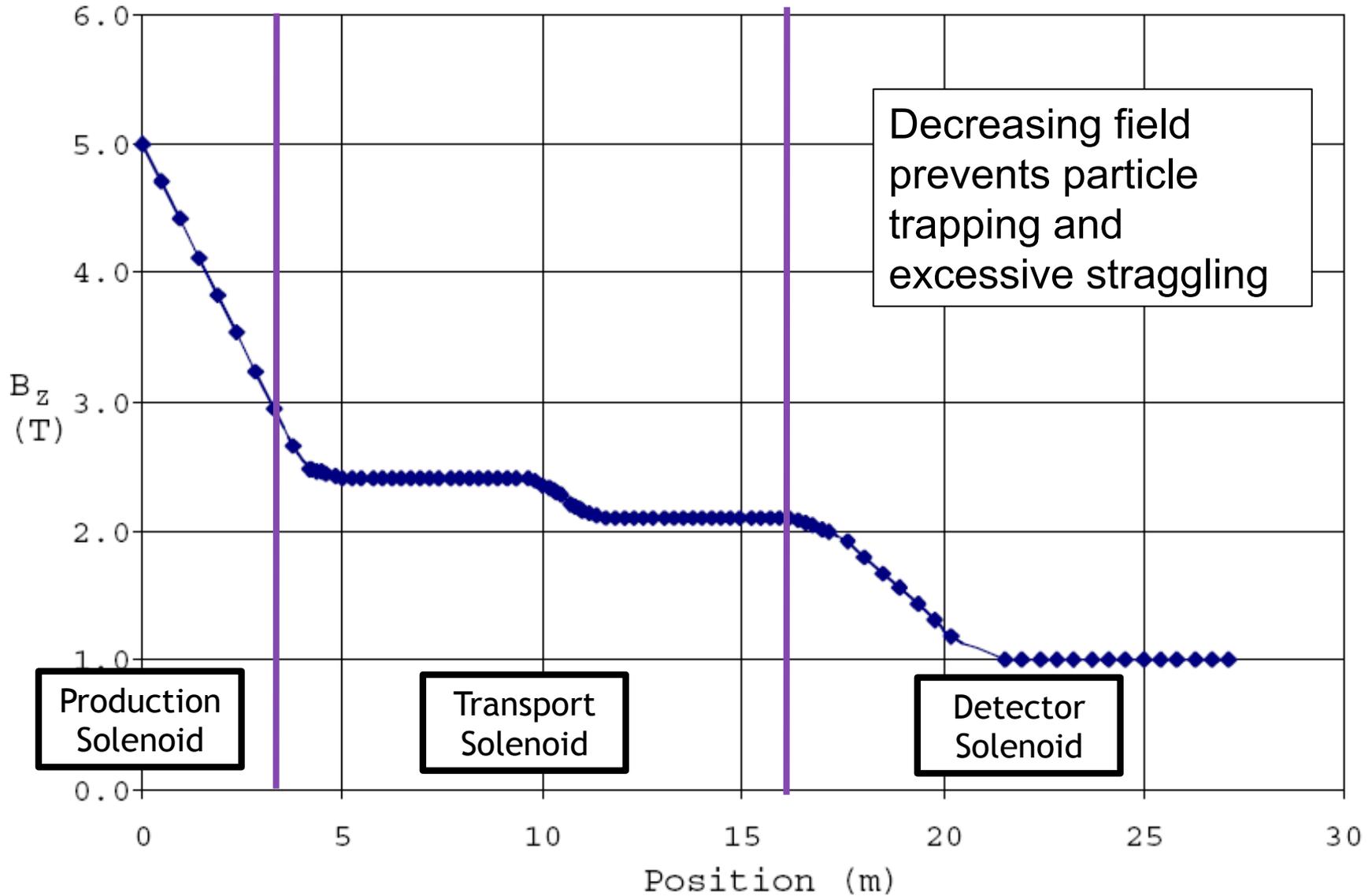
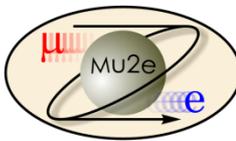
Mu2e: The Big Picture



- **Production Target**
 - Proton beam strikes target, producing mostly pions
- **Production Solenoid**
 - Contains backwards pions/muons and reflects slow forward pions/muons
- **Transport Solenoid**
 - Selects low momentum, negative muons
- **Capture Target, Detector, and Detector Solenoid**
 - Capture muons on target and wait for them to decay
 - Detector blind to ordinary (Michel) decays, with $E \leq \frac{1}{2}m_{\mu}c^2$
 - Optimized for $E \sim m_{\mu}c^2$

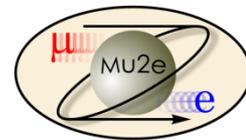


Magnetic Field Gradient





Target and Heat Shield



- Produces pions which decay into muons

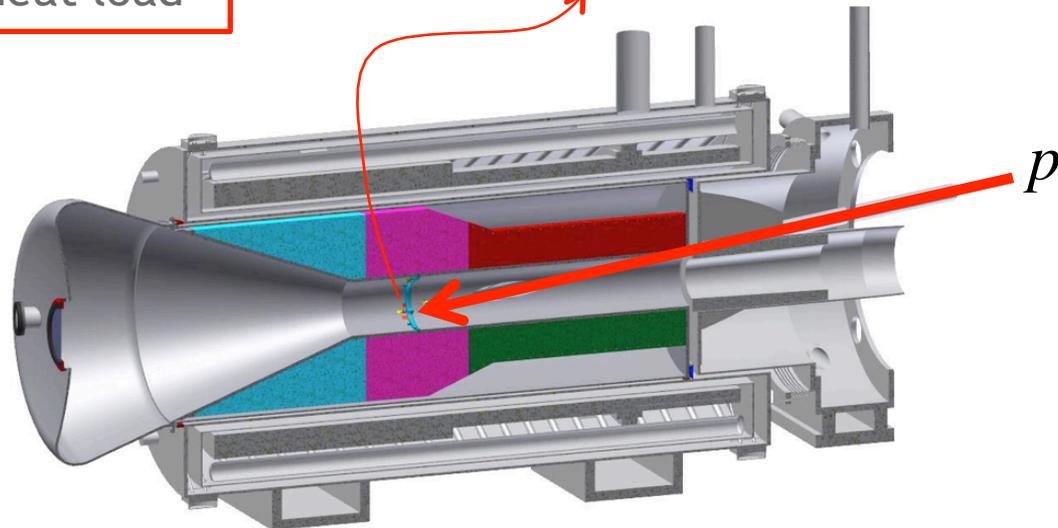
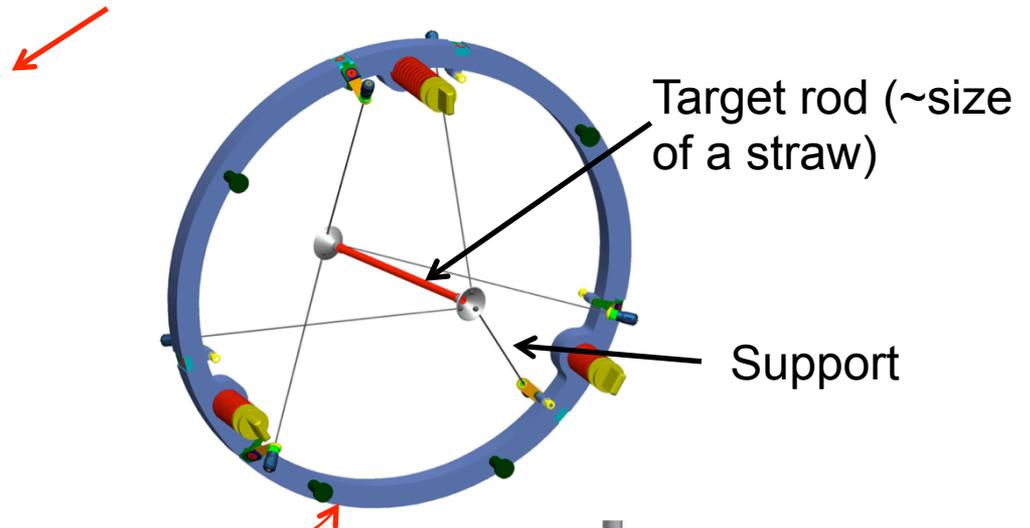
- Tungsten Target

- 8 kW beam
- 700 W in target
- Radiatively cooled

- Heat Shield

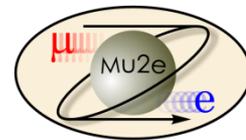
- Bronze insert
- 3.3 kW average heat load

Remember, this is inside a superconducting magnet

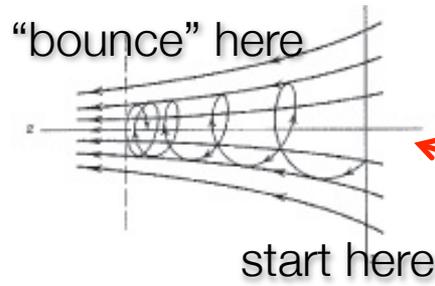
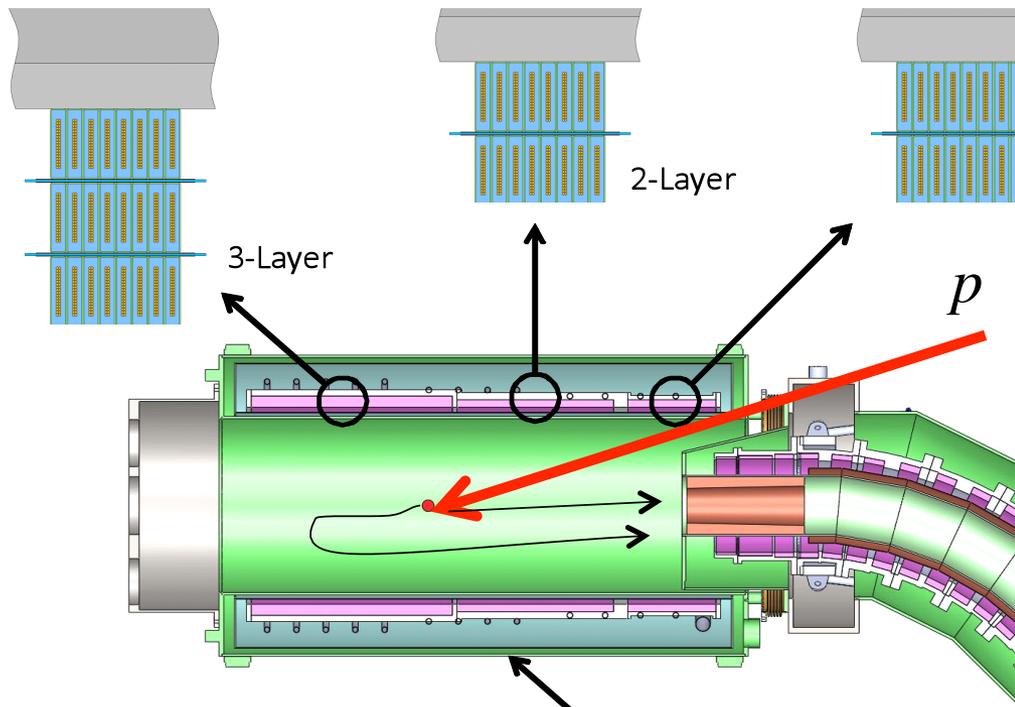




Production Solenoid

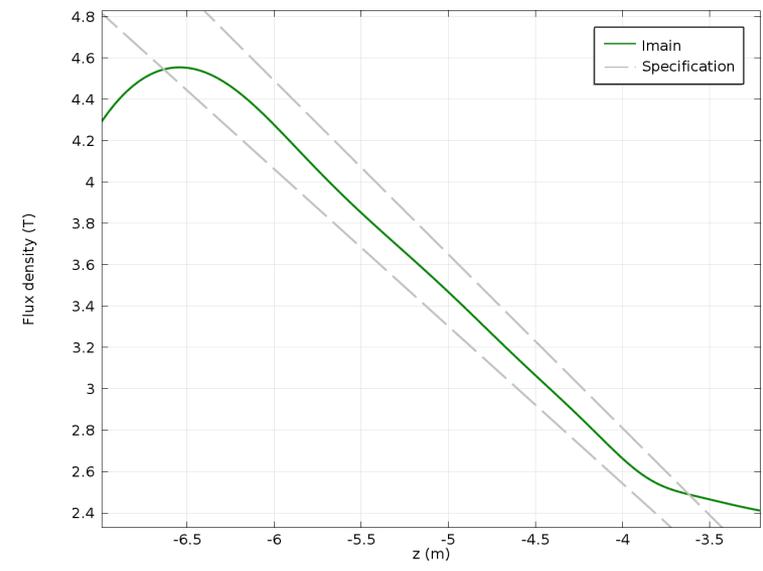


- Axially graded ($\sim 5\text{T} \rightarrow 2.5\text{T}$) solenoid captures low energy backward and reflected pions, directing to the Transport Solenoid



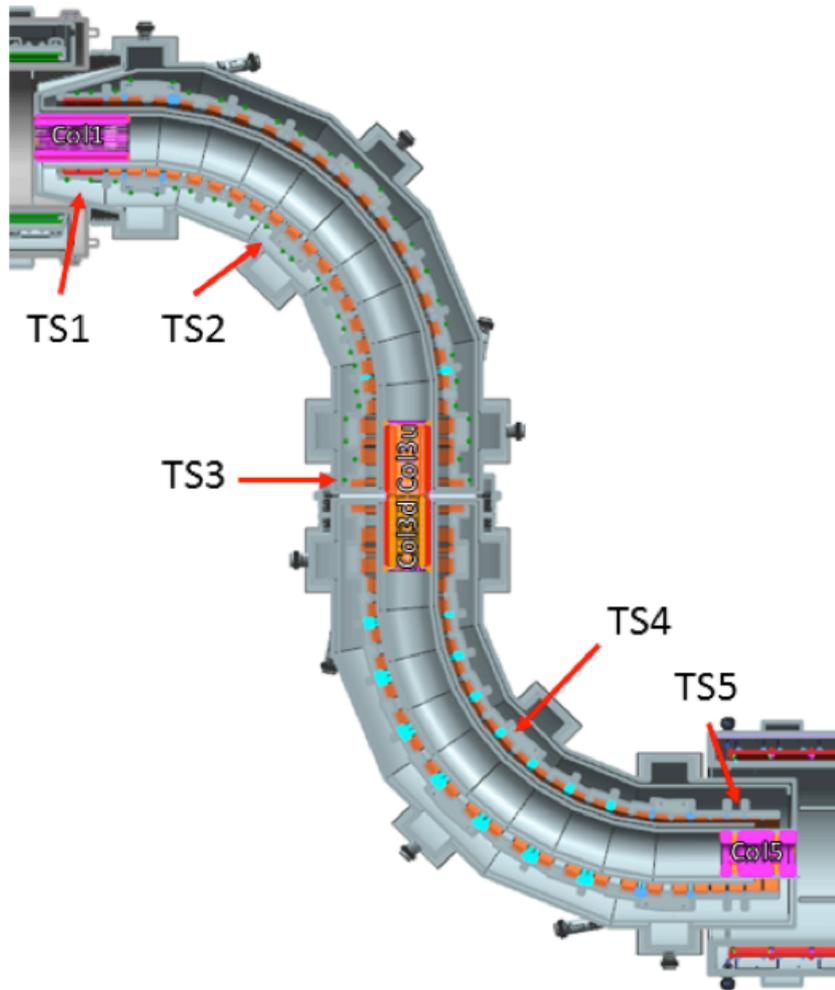
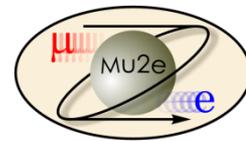
Magnetic reflection (pinch confinement)

Magnetic Gradient

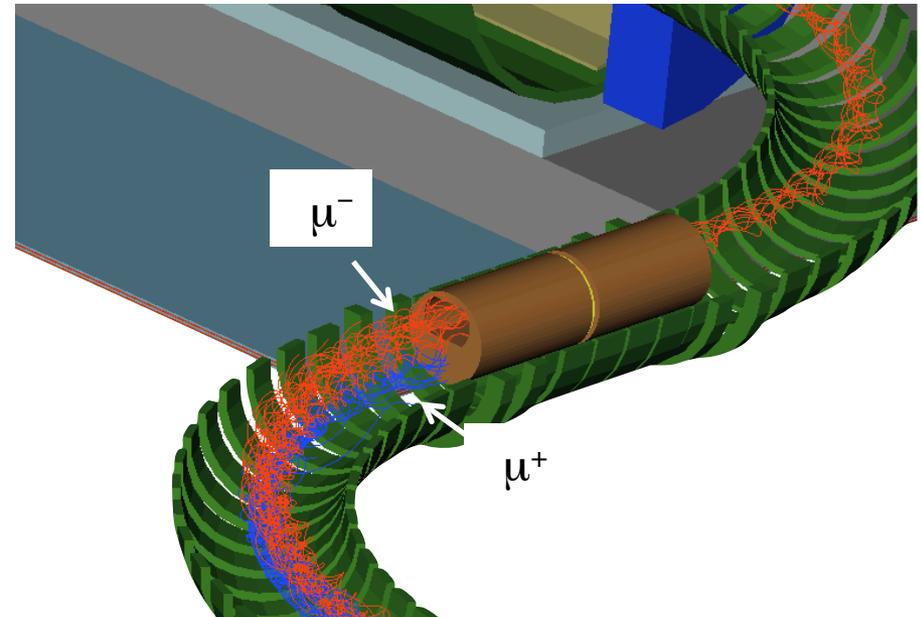


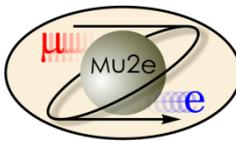


Transport Solenoid



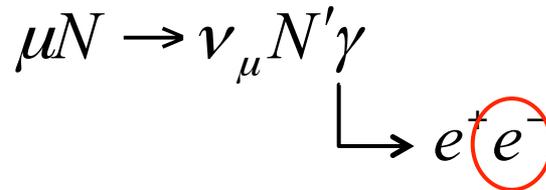
- Transports muons from production target to capture target
- Curved solenoid eliminates line-of-sight backgrounds
- Collimator in center selects low momentum negative muons
 - RxB drift causes sign/momentum dependent *vertical* displacement





Choosing the Capture Target

- The probability of exchanging a virtual particle with the nucleus goes up with Z , however
- Lifetime is *shorter* for high- Z
 - Decreases useful live window
- Also, need to avoid background from radiative muon capture limits choices



⇒ Want $M(Z) - M(Z-1)$
< signal energy

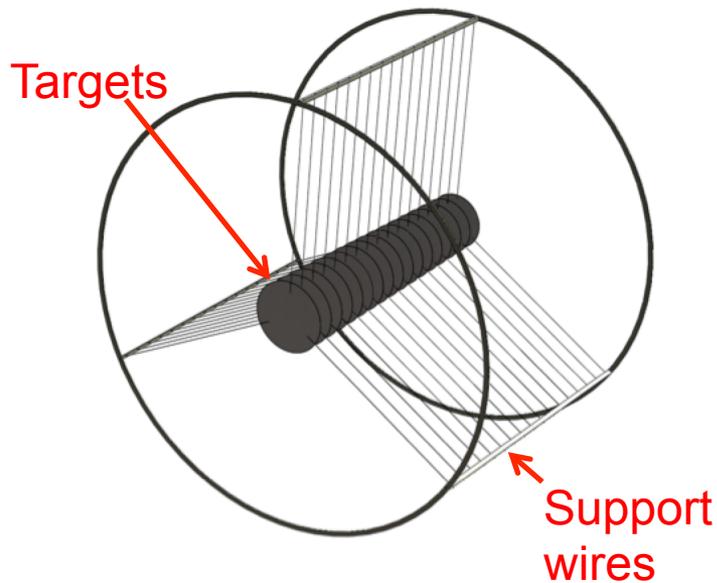
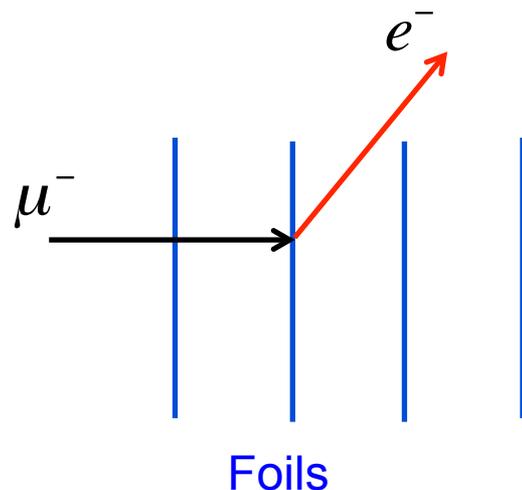
⇒ Aluminum is initial choice for Mu2e

Nucleus	$R_{\mu e}(Z) / R_{\mu e}(Al)$	Bound lifetime	Atomic Bind. Energy(1s)	Conversion Electron Energy	Prob decay >700 ns
Al(13,27)	1.0	.88 μs	0.47 MeV	104.97 MeV	0.45
Ti(22,~48)	1.7	.328 μs	1.36 MeV	104.18 MeV	0.16
Au(79,~197)	~0.8-1.5	.0726 μs	10.08 MeV	95.56 MeV	negligible

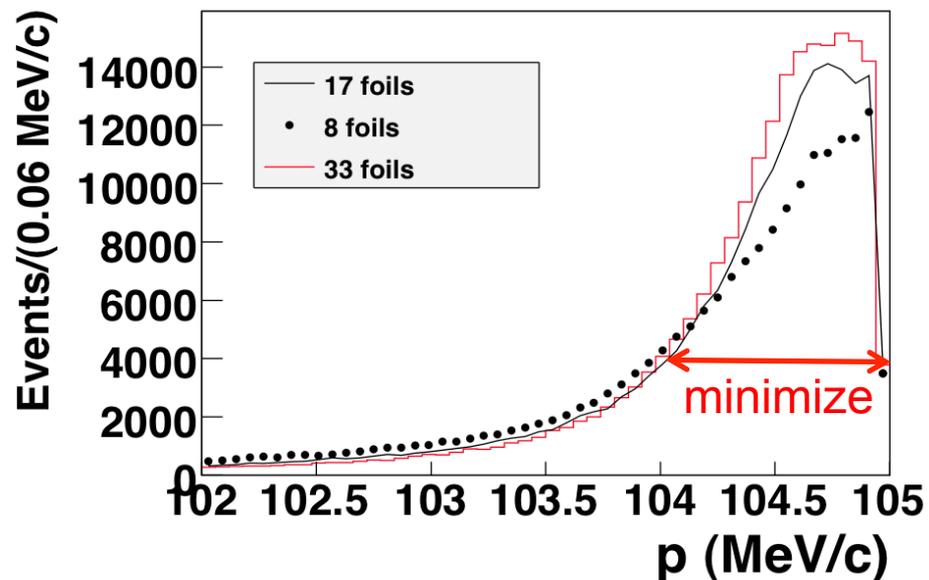


Stopping (capture) Target

- Multiple thin layers to allow decay or conversion electrons to exit with minimal scattering
 - 17 Aluminum foils
 - 200 μm thick
- Stops 49% of arriving muons

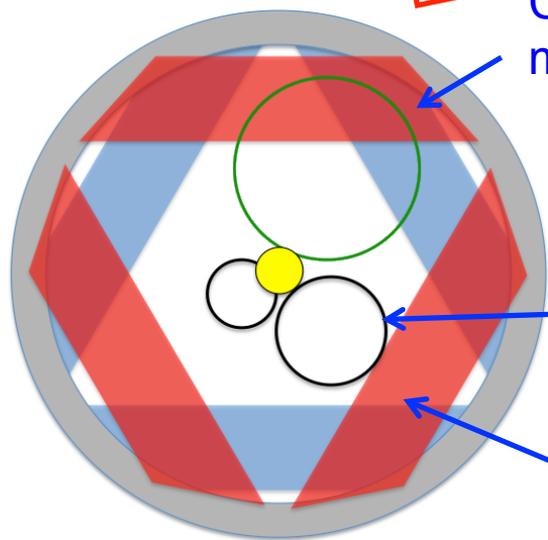
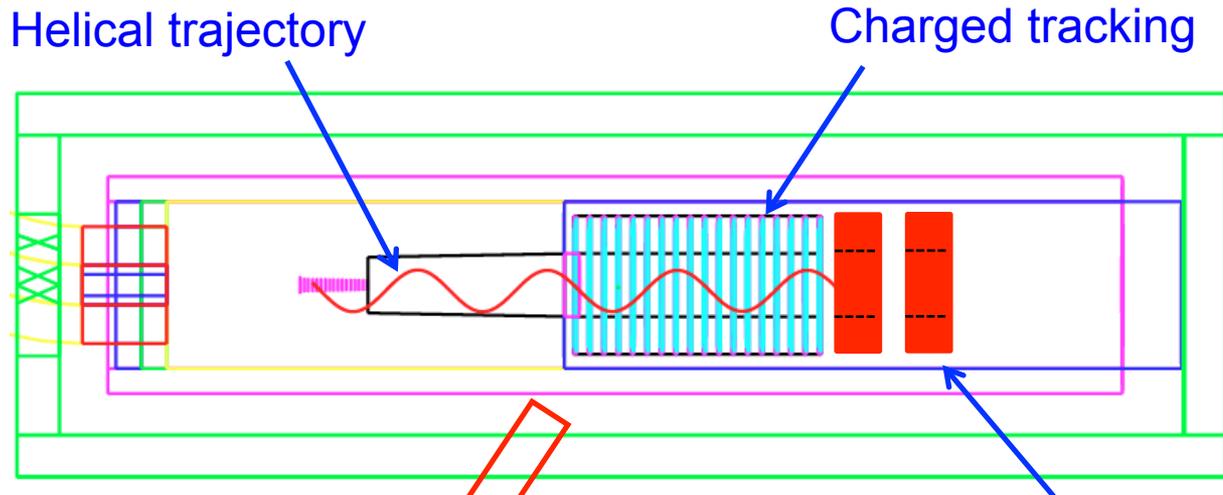


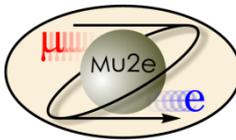
Conversion electron spectrum:





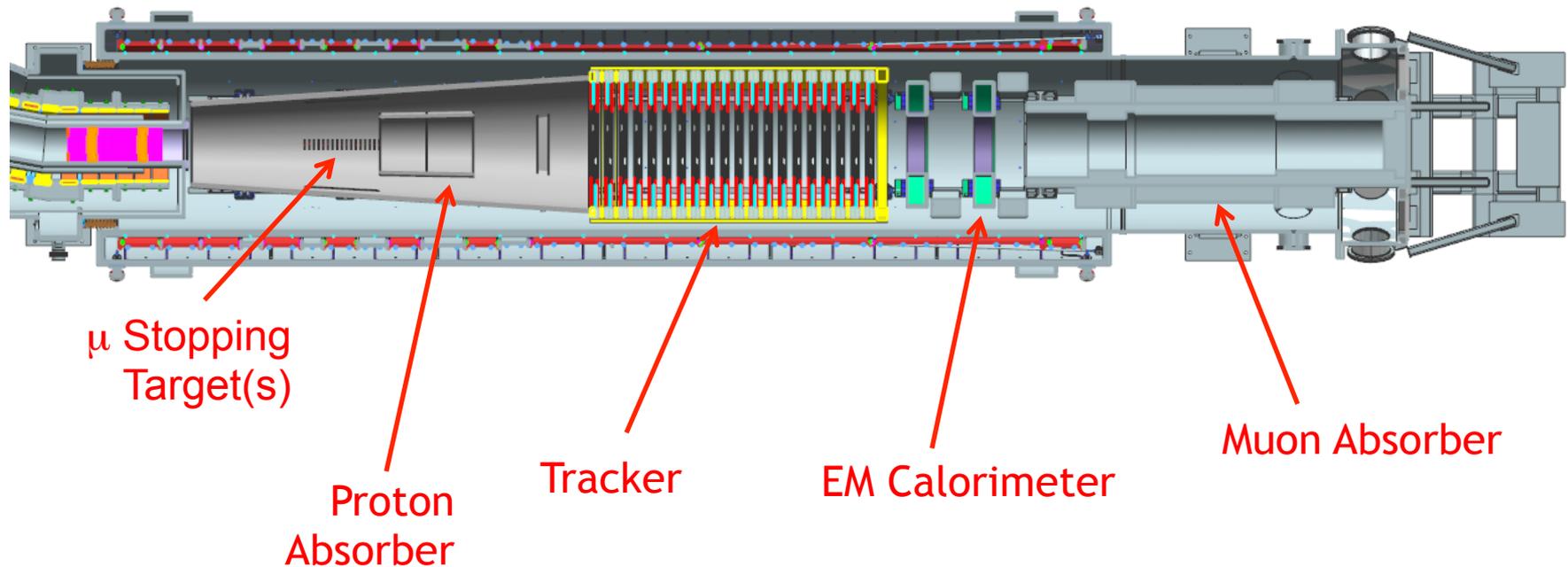
Particle Detector

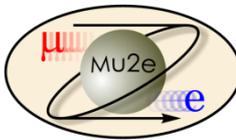




Detector and Detector Solenoid

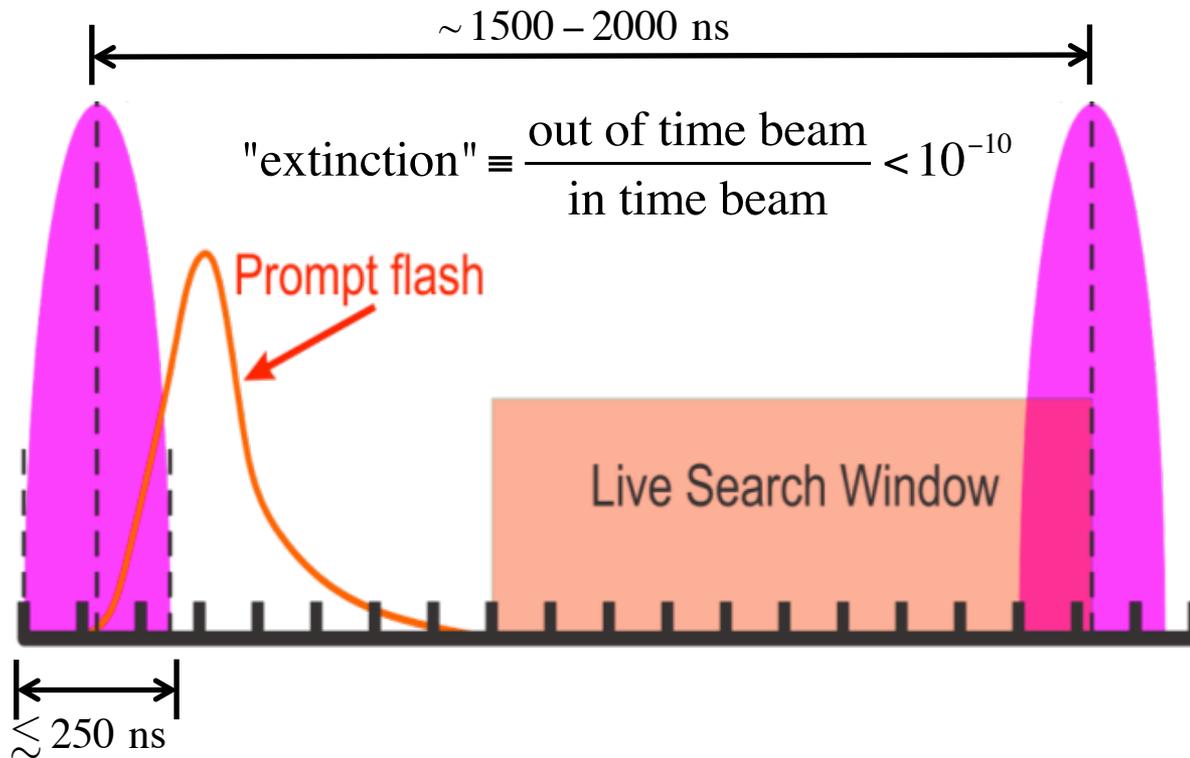
- Graded field around stopping target to increase acceptance
 - Magnetic reflection again
- Uniform field in tracking volume
- Electromagnetic calorimeter to tag electrons.





Beam Needs

- We've talked about the experiment. Now where do we put it?
- Remember, we need a beam that looks kind of like this



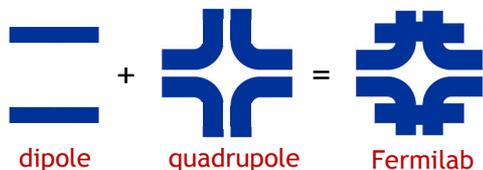
- This is where Fermilab comes in...



A Brief History of Fermilab



Trivia: original Main Ring was the first “separated function” synchrotron



- 1968: construction begins
- 1972: first beams from Main Ring
 - 200→400 GeV proton beams to fixed targets
 - Highest energy lab for next 36 years!
- ~1985:
 - “Tevatron”: first superconducting synchrotron shares tunnel with Main Ring
 - 900GeV x 900 GeV p-pBar collisions
 - Highest energy collider for 23 years.
- 1997: Major upgrade
 - Main Injector replaces Main Ring -> more intensity
 - 980 GeV x 980 GeV p-pBar collisions
 - Intense neutrino program
- 2011: Tevatron permanently turned off after the LHC came full online.
- **So what is the lab doing now?**



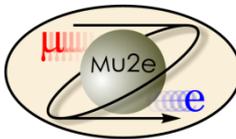
Guidance: The P5 Report



- The Particle Physics Project Prioritization Panel (P5) advises the DOE Office of High Energy Physics.
- In 2013, the P5 was charged to determine priorities in US particle physics (primarily priorities for Fermilab) under various funding scenarios
- In 2014, the panel report recommended proceeding with Mu2e under all funding scenarios.

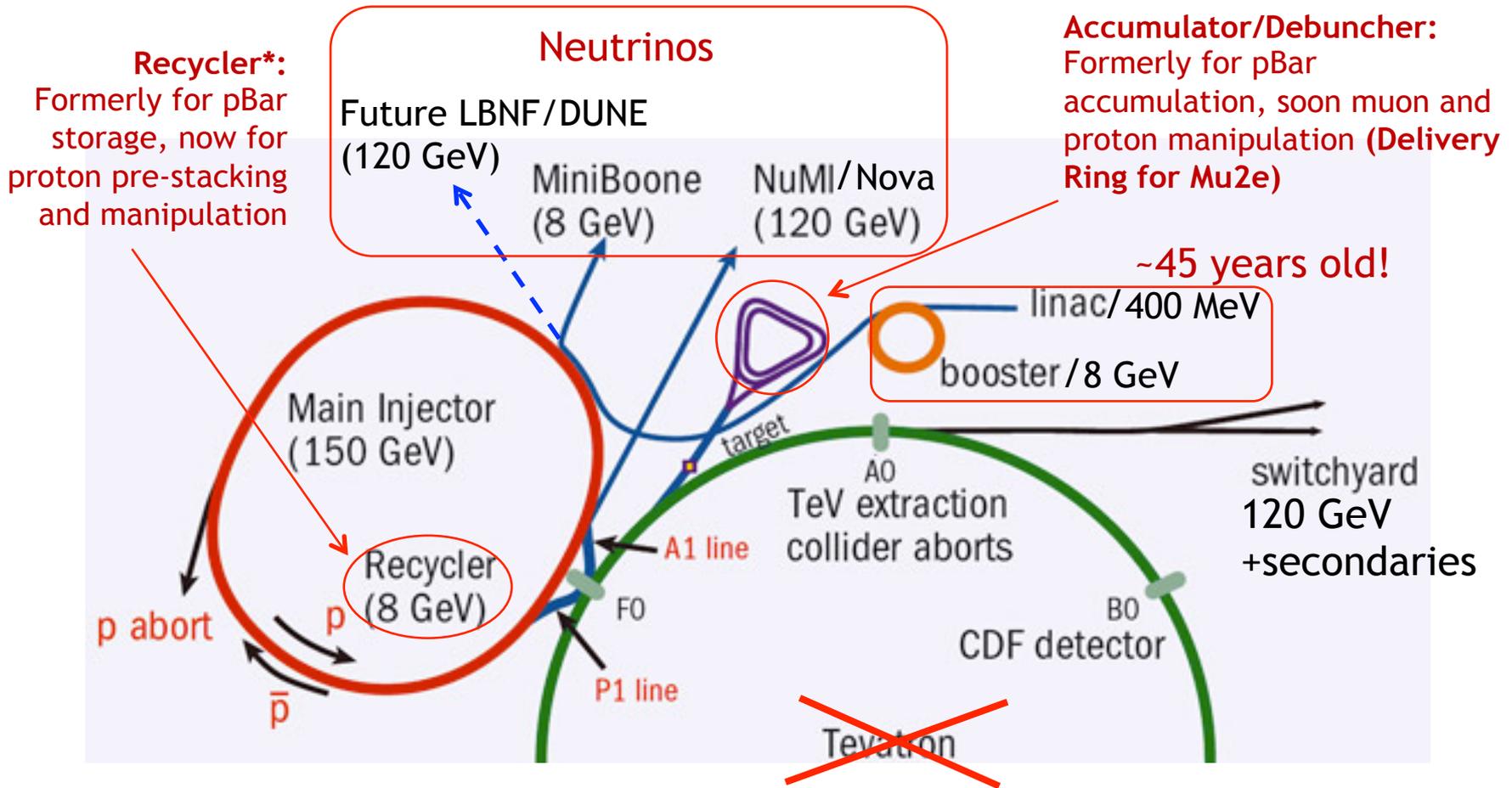
Project/Activity	bleak Scenario A	flat Scenario B	fantasy Scenario C
Large Projects			
Muon program: Mu2e, Muon g-2	Y, <small>Mu2e small reprofile needed</small>	Y	Y
HL-LHC	Y	Y	Y
LBNF + PIP-II	Y, <small>LBNF components delayed relative to Scenario B.</small>	Y	Y, enhanced
ILC	R&D only	R&D, <small>possibly small hardware contributions. See text.</small>	Y
NuSTORM	N	N	N
RADAR	N	N	N

- So... full speed ahead!



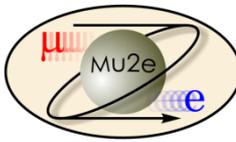
Fermilab Accelerator Complex Today

- Now tha LHC has taken over the Energy Frontier, Fermilab is focusing on intensity-based physics

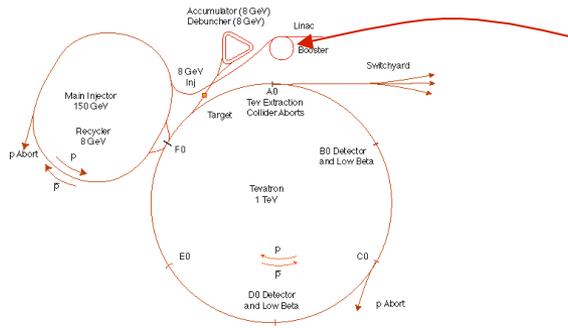


~45 years old!

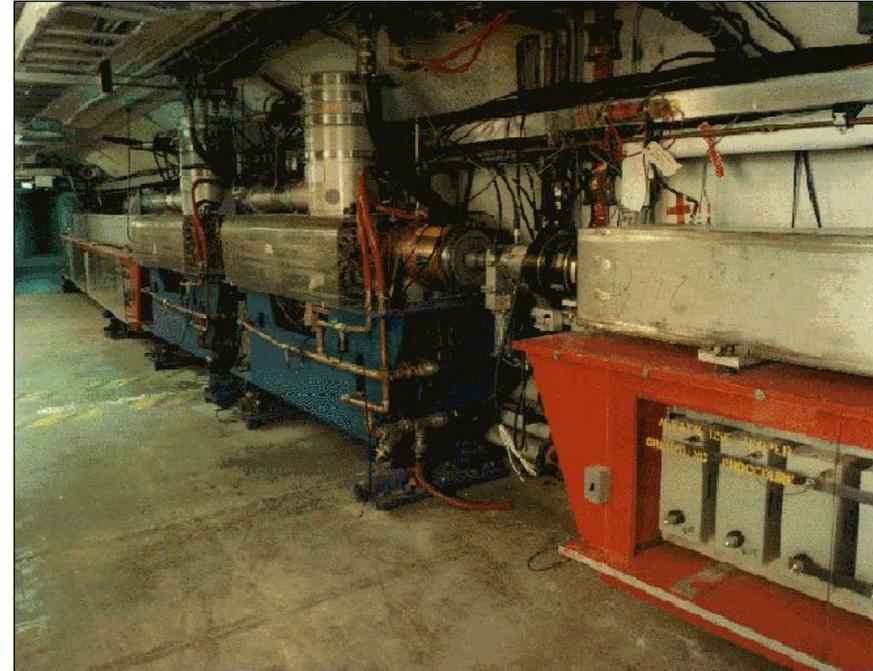
*first permanent magnet storage ring



Fermilab Booster



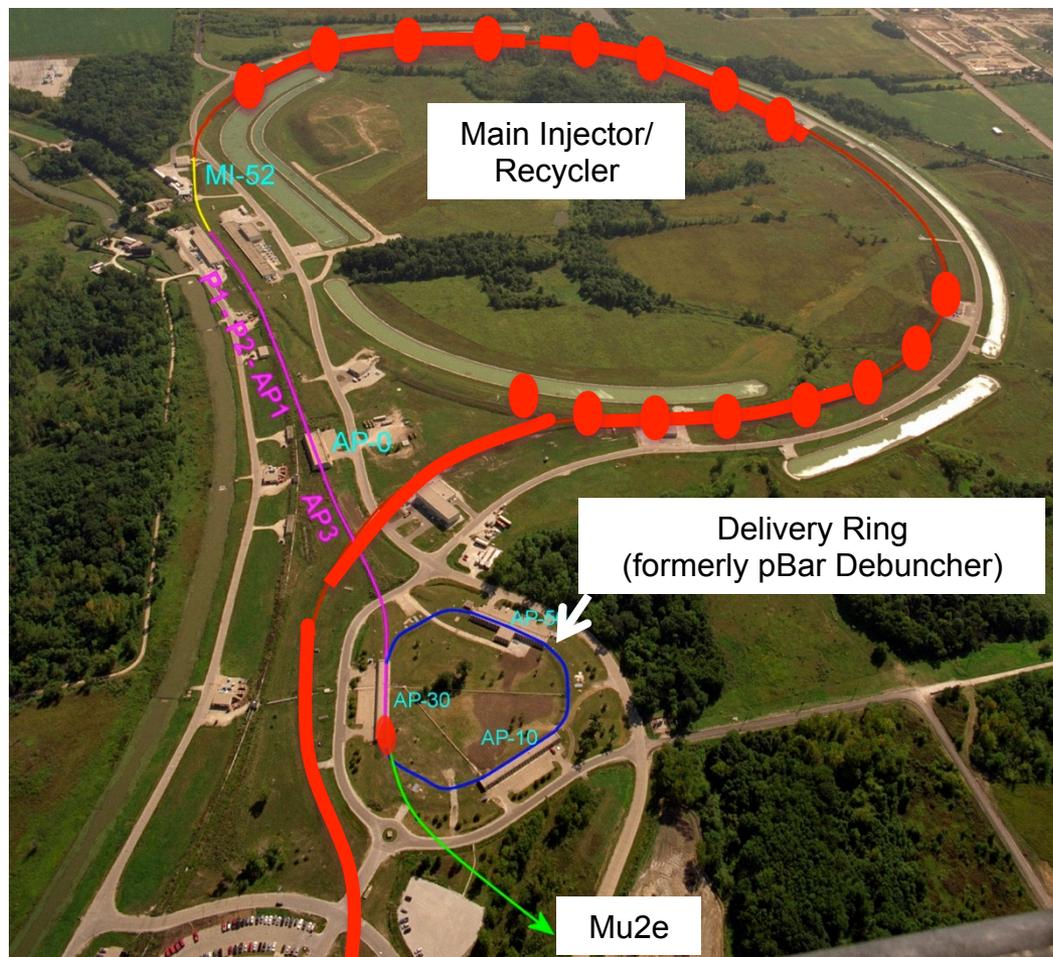
Most “original” part of the complex



- Accelerates protons from 400 MeV to 8 GeV
- Operates in a 15 Hz resonant circuit
 - No time for beam manipulation
 - Can't make required beam structure
- Sets a fundamental clock for the complex
 - 15 Hz “tick”
- Sets a fundamental unit of protons
 - 1 “batch” = up to $\sim 4 \times 10^{12}$ protons
- Since the can't make the beam we need, how do we do it?
 - **By using almost everything else (impossible in Tevatron era)!**



Mu2e Proton Delivery

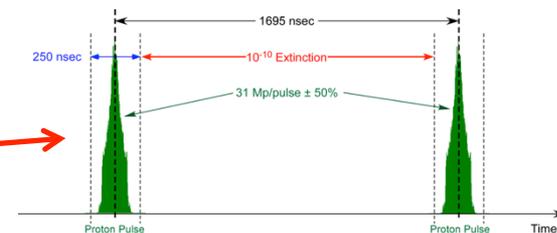


Booster

Mu2e

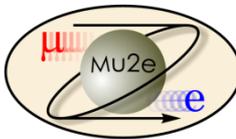
- Two Booster “batches” are injected into the Recycler (8 GeV storage ring). Each is:
 - 4×10^{12} protons
 - 1.7 μ sec long
- These are divided into 8 bunches of 10^{12} each
- The bunches are extracted one at a time to the Delivery Ring
 - Period = 1.7 μ sec
- As the bunch circulates, it is resonantly extracted to produce the desired beam structure.
 - Bunches of $\sim 3 \times 10^7$ protons each
 - Separated by 1.7 μ sec

Exactly what we need →



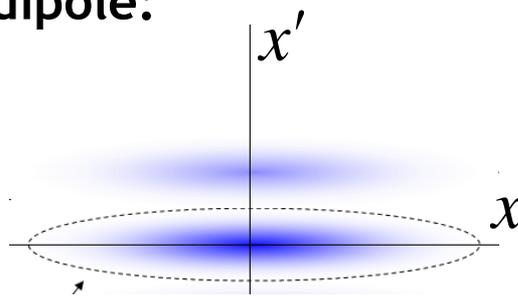


Eliminating out of Time Beam (Extinction)

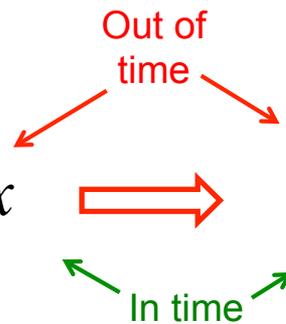


- The bunches from the Delivery Ring will have $\sim 10^{-5}$ extinction
 - We need 10^{-10} to make prompt backgrounds small compared to other backgrounds
- A set of resonant dipoles in the beam line will deflect the beam such that only in-time beam is transmitted through a downstream collimator:

At dipole:



Angular deflection



At collimator:



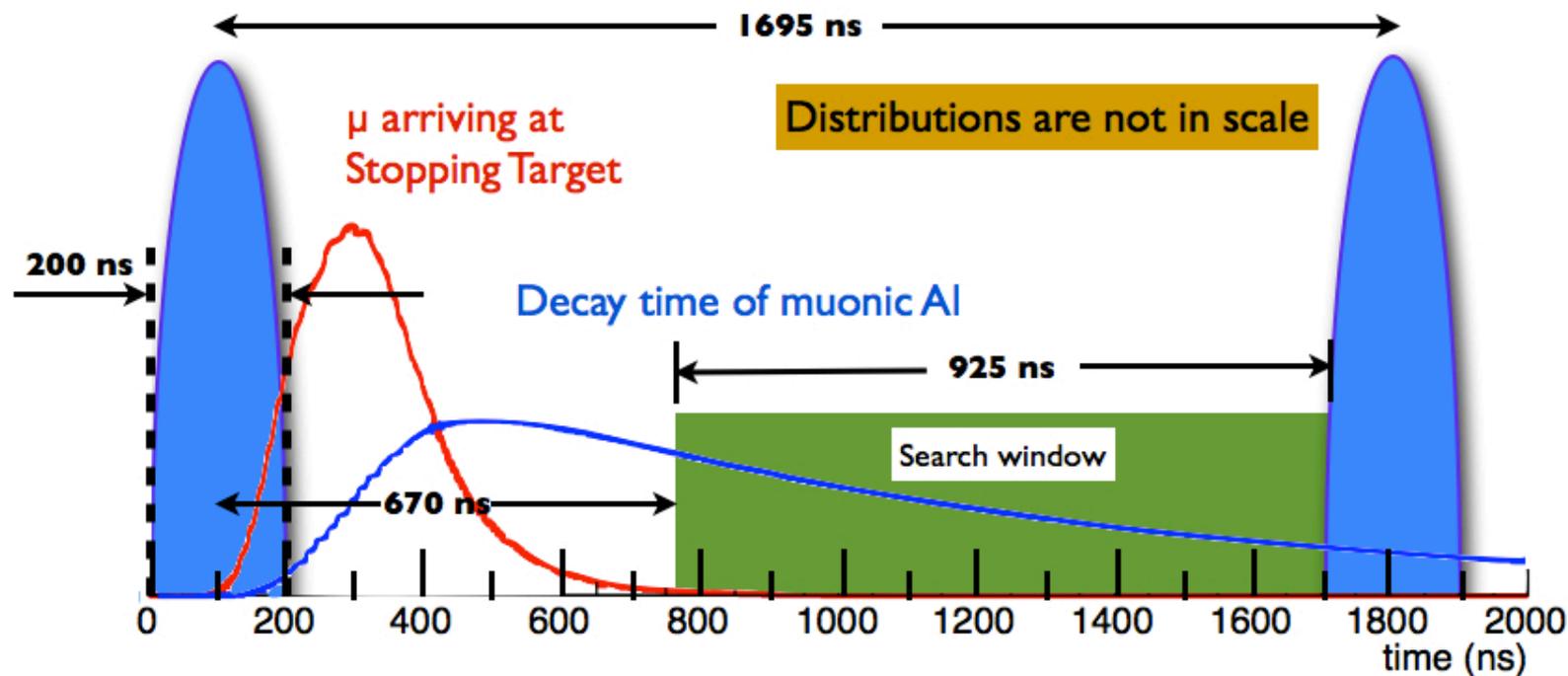
Spatial offset

- Think miniature golf





End Product



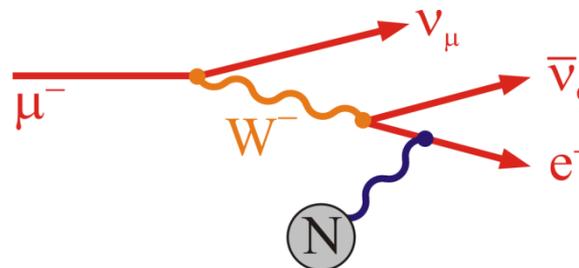
Target data set: $\sim 3.6 \times 10^{20}$ protons in ~ 3 years



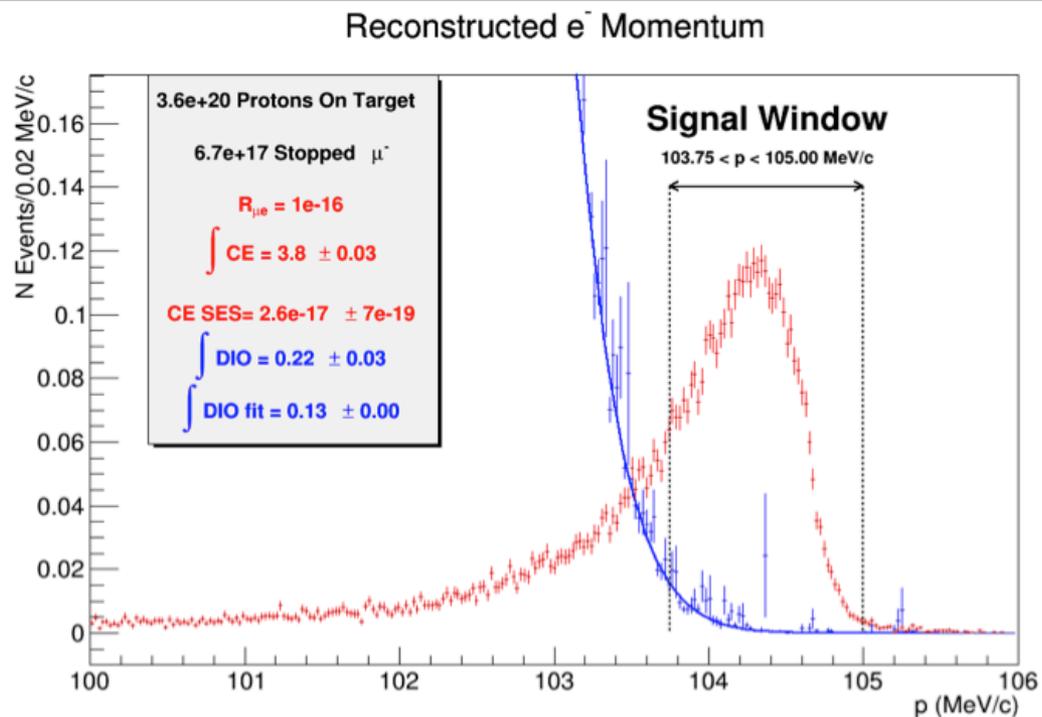
Major Backgrounds Revisited



1. Muon decay in orbit (DIO)

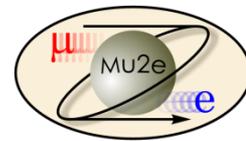


Defeated by good energy resolution





Major Backgrounds (cont'd)

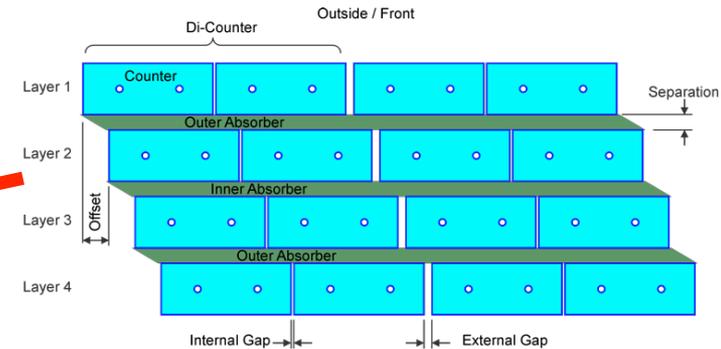
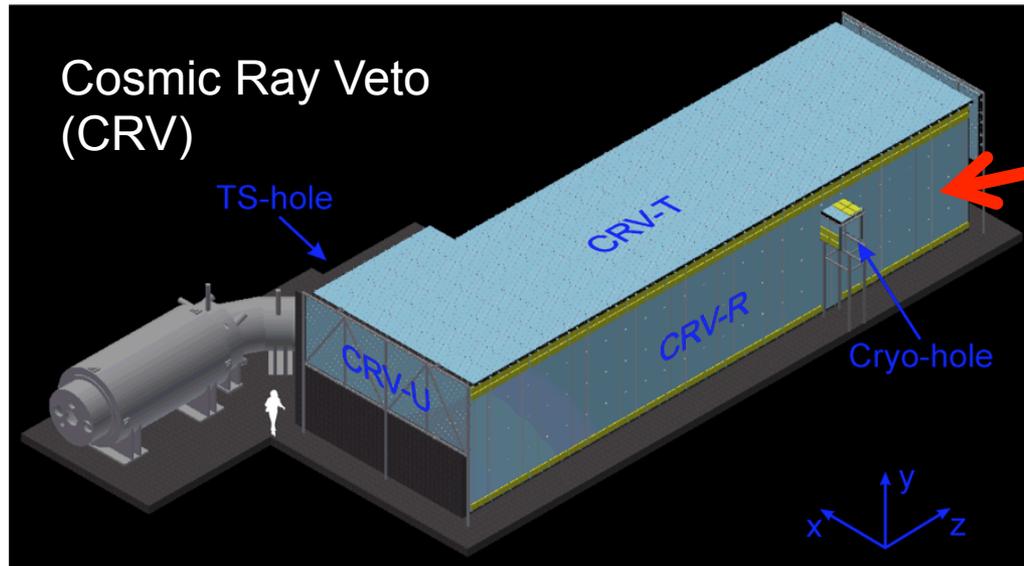


2. Beam Related Backgrounds

Suppressed with 10^{-10} extinction (just talked about this)

3. Asynchronous Backgrounds: Cosmic Rays

Suppressed by active and passive shielding



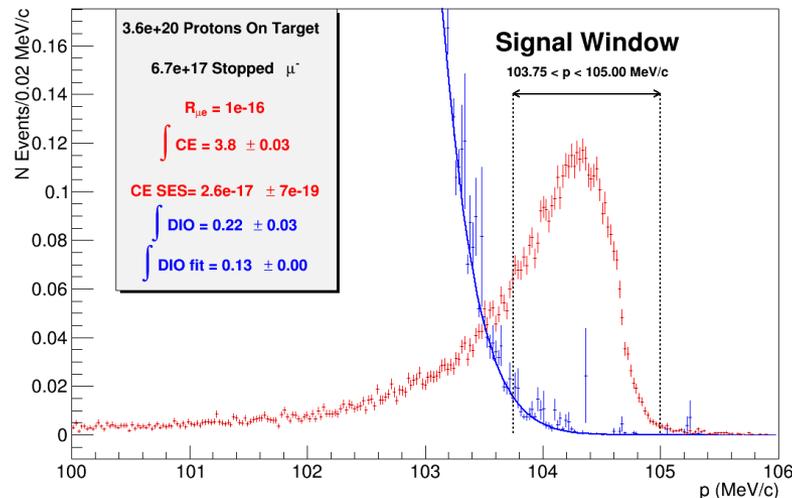
- Four layers of scintillator surround experiment
- Efficiency goal: $>99.99\%$



Sensitivity



- Full GEANT4 Simulation
- 3.6×10^{20} protons on target
 - 3 years nominal running
- Cuts chosen to maximize sensitivity



Parameter	Value
Physics run time @ 2×10^7 s/yr.	3 years
Protons on target per year	1.2×10^{20}
μ^- stops in stopping target per proton on target	0.0019
μ^- capture probability	0.609
Total acceptance x efficiency	$(8.5 \pm_{0.9}^{1.1})\%$
Single-event sensitivity with Current Algorithms	$(2.87 \pm_{0.27}^{0.32}) \times 10^{-17}$

Single Event Sensitivity: $R_{\mu e} = 2.9 \times 10^{-17}$



Significance



➤ Backgrounds

Category	Background process	Estimated yield (events)
Intrinsic	Muon decay-in-orbit (DIO)	0.199 ± 0.092
	Muon capture (RMC)	$0.000^{+0.004}_{-0.000}$
Late Arriving	Pion capture (RPC)	0.023 ± 0.006
	Muon decay-in-flight (μ -DIF)	<0.003
	Pion decay-in-flight (π -DIF)	$0.001 \pm <0.001$
Miscellaneous	Beam electrons	0.003 ± 0.001
	Antiproton induced	0.047 ± 0.024
	Cosmic ray induced	0.092 ± 0.020
Total		0.37 ± 0.10

8 GeV is a stupid energy

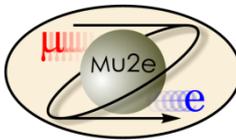


➤ Bottom line:

- Single event sensitivity: $R_{\mu e} = 3 \times 10^{-17}$
- 90% C.L. (if no signal) : $R_{\mu e} < 7 \times 10^{-17}$
- Typical SUSY Signal: ~ 40 events or more

4 order of magnitude improvement!





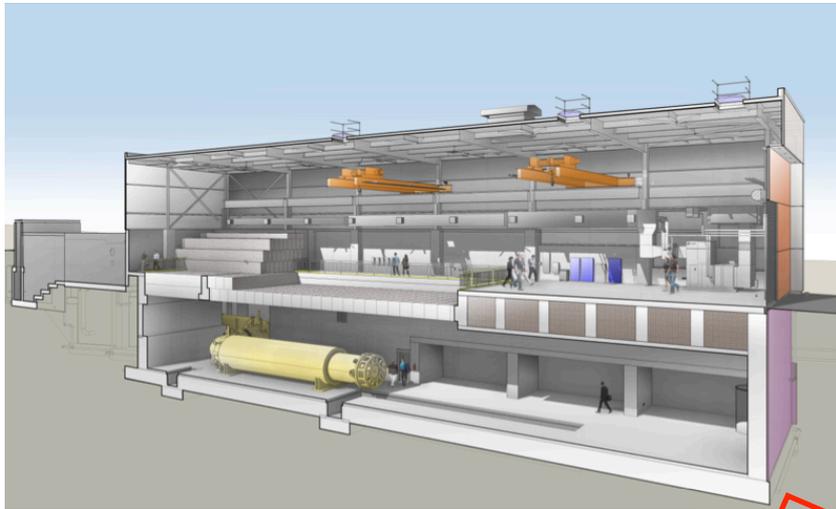
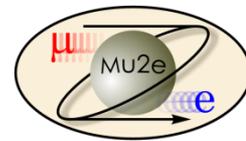
A long time coming

- 1992 Proposed as “MELC” at Moscow Meson Factory
- 1997 Proposed as “MECO” at Brookhaven
(at this time, experiment incompatible with Fermilab)
- 1998-2005 Intensive work on MECO technical design
- July 2005 Entire rare-decay program canceled at Brookhaven
- 2006 MECO subgroup + Fermilab physicists work out means to mount experiment at Fermilab
- Fall 2008 Mu2e Proposal submitted to Fermilab
- November 2008 Stage 1 approval. Formal Project Planning begins
- November 2009 DOE Grants CD-0 ← In DOE project-speak, this is the first “Critical Decision”: Statement of mission need = official existence
- July 2012 CD-1
- March 2015 CD-2/3b ← Approval of baseline and money for long lead elements

Finally, things are really happening!



Civil Construction





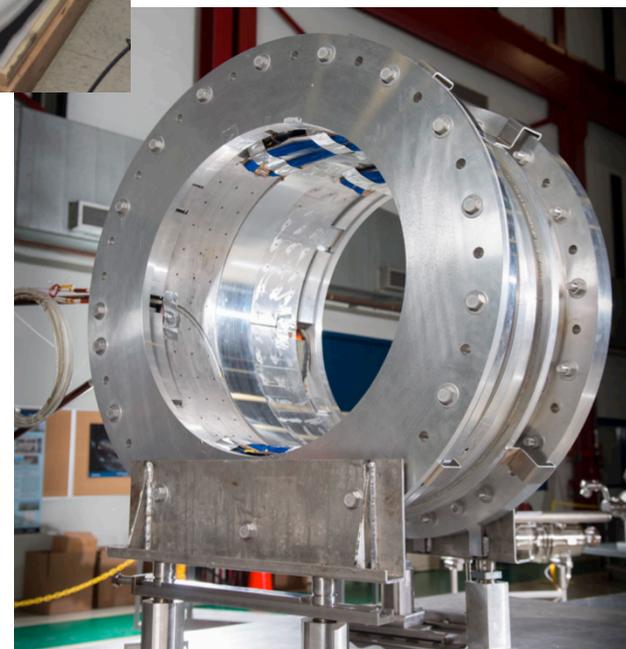
Magnet Procurement and Testing



Cable acceptance

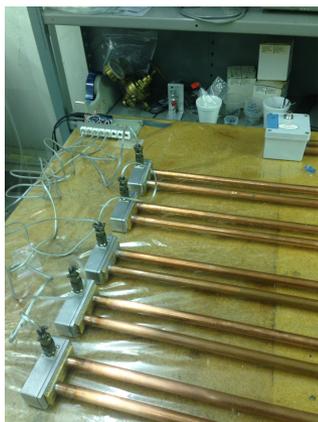


Successful test of Transport Solenoid segment

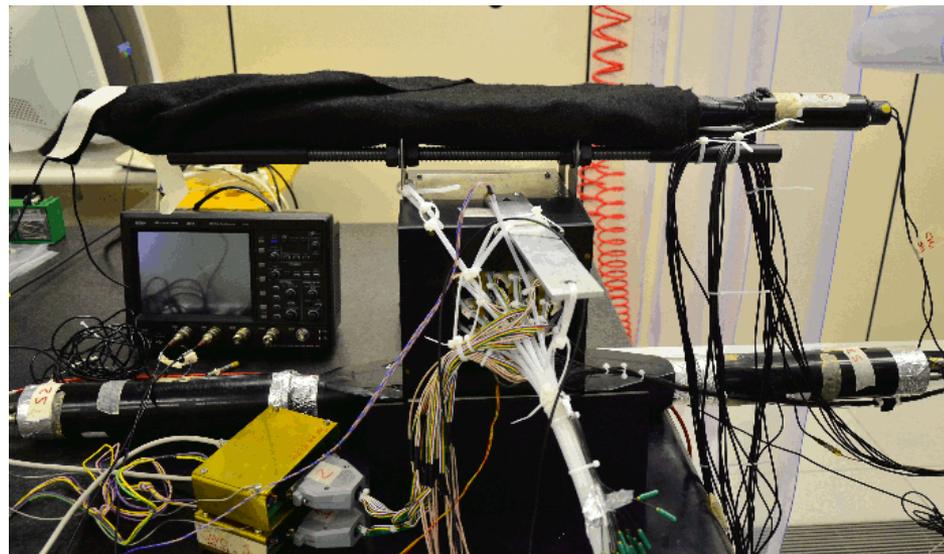




Detectors



Straw Tube Tracker



Calorimeter Crystal Test



Cosmic Ray Veto

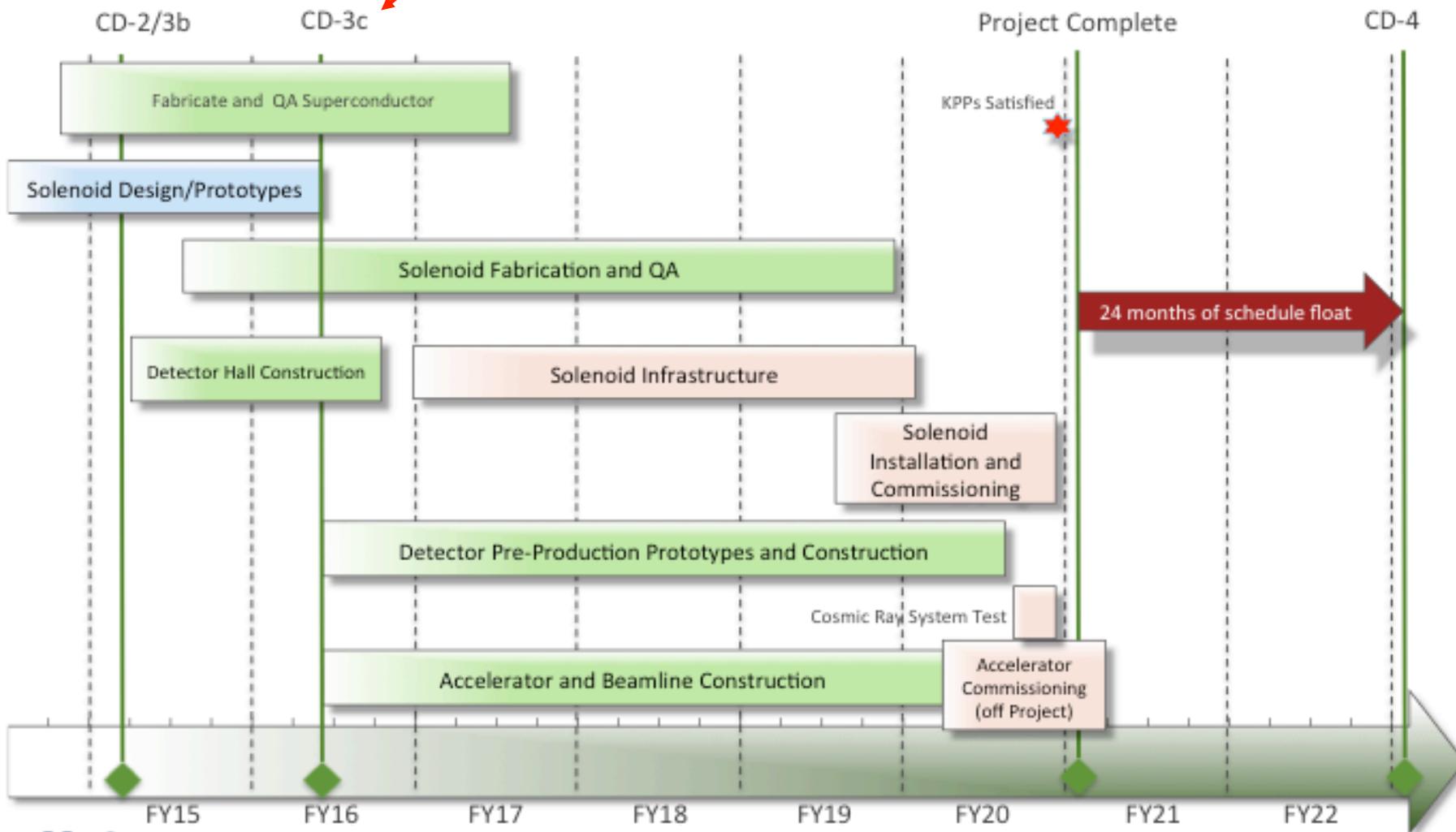




Schedule



Scheduled for June 14



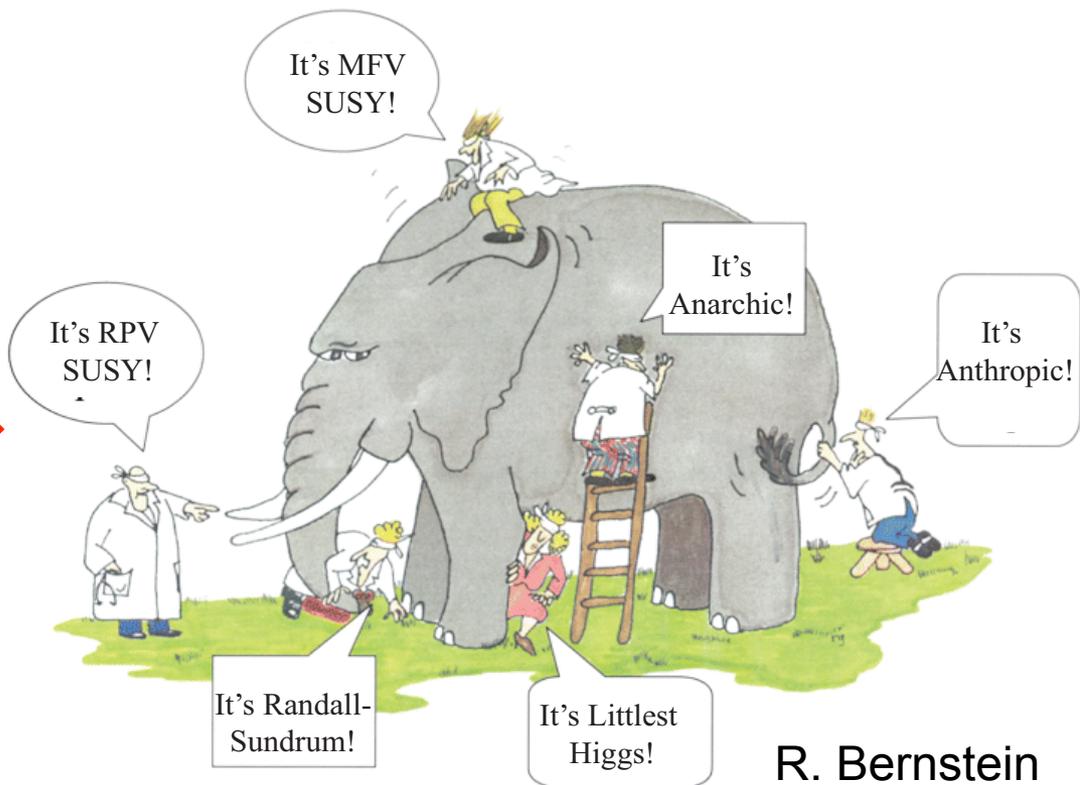
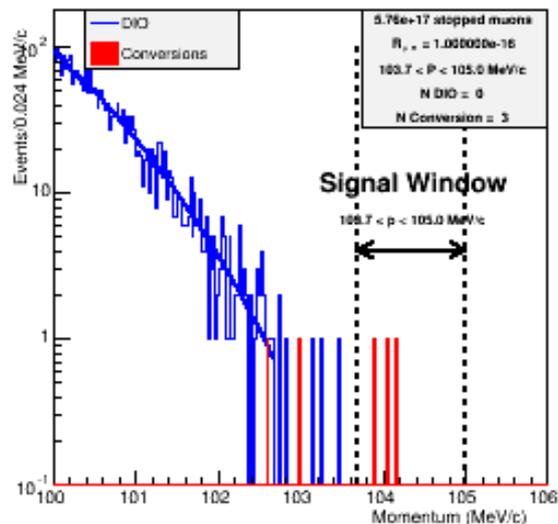


What if we see something?



$$R_{\mu e} = 10^{-16}$$

Toy Mu2e Experiment



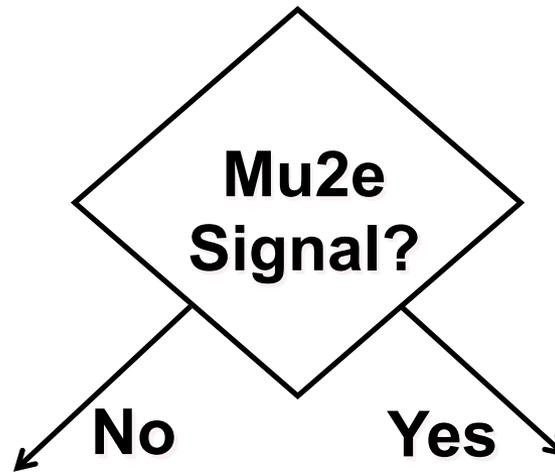
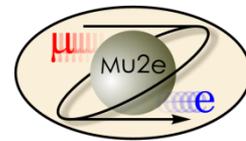
R. Bernstein

Next questions:

- What's the $\mu \rightarrow e \gamma$ signal (if any)?
- What's the target dependence?



Upgrade scenarios

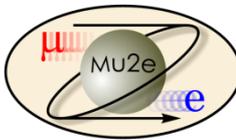


- Both prompt and DIO backgrounds must be lowered to measure
 $R_{\mu e} \sim 10^{-18}$
- Must upgrade all aspects of production, transport and detection.

- Must compare different targets.
- Optimize muon transport and detector for short bound muon lifetimes.
- Backgrounds might not be as important.



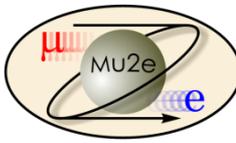
Conclusions



- We have proposed a realistic experiment to measure

$$R_{\mu e} \equiv \frac{\Gamma(\mu^- \text{Al} \rightarrow e^- + \text{Al})}{\Gamma(\mu^- \text{Al} \rightarrow (\text{All Captures}))}$$

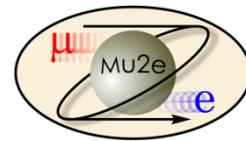
- Single event sensitivity of $R_{\mu e} = 3 \times 10^{-17}$
- This represents an improvement of *four orders of magnitude* compared to the existing limit, or over a *factor of ten* in effective mass reach. For comparison
 - TeV -> LHC = factor of 7 (difference in luminosity makes in comparable)
 - LEP 200 -> ILC = factor of 2.5
- ANY signal would be unambiguous proof of physics beyond the Standard Model
- The absence of a signal would be a very important constraint on proposed new models.



BACKUP SLIDES



Just How Rare is that?



Probability of...	
rolling a 7 with two dice	1.67E-01
rolling a 12 with two dice	2.78E-02
getting 10 heads in a row flipping a coin	9.77E-04
drawing a royal flush (no wild cards)	1.54E-06
getting struck by lightning in one year in the US	2.00E-06
winning Pick-5	5.41E-08
winning MEGA-millions lottery (5 numbers+megaball)	3.86E-09
your house getting hit by a meteorite this year	2.28E-10
drawing two royal flushes in a row (fresh decks)	2.37E-12
your house getting hit by a meteorite today	6.24E-13
getting 53 heads in a row flipping a coin	1.11E-16
your house getting hit by a meteorite AND you being struck by lightning both within the next six months	1.14E-16
your house getting hit by a meteorite AND you being struck by lightning both within the next three months	2.85E-17

← Sindrum limit

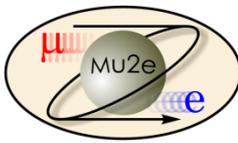
← Single event sensitivity of Mu2e



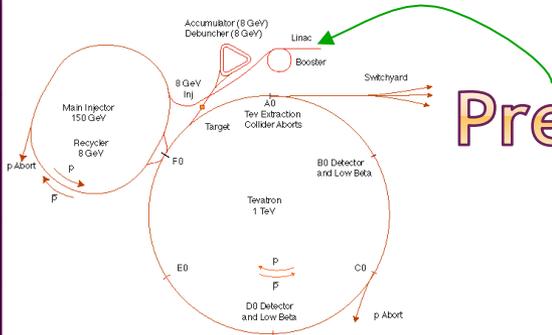
Experimental Challenges for Increased Flux



- At our level of sensitivity, we hit fundamental limits with this technique
 - Simply increasing the proton flux will not improve the limit dramatically
- Improve momentum resolution for the ~ 100 MeV electrons to reject high energy tails from ordinary DIO electrons.
 - Limited by multiple scattering in target and detector plane
 - ➔ go to bunched, mono-energetic muon beam, allowing for thinner target
- Allow longer decay time for pions to decay
- Both of these lead to a decay/compressor ring
- Other issues with increased flux
 - Upgrade target and capture solenoid to handle higher proton rate
 - Target heating
 - Quenching or radiation damage to production solenoid
 - High rate detector
- All of these efforts will benefit immensely from the knowledge and experience gained during the initial phase of the experiment.
- If we see a signal a lower flux, can use increased flux to study in detail
 - Precise measurement of $R_{\mu e}$
 - Target dependence
 - Comparison with $\mu \rightarrow e\gamma$ rate



Preac(cellerator) and Linac

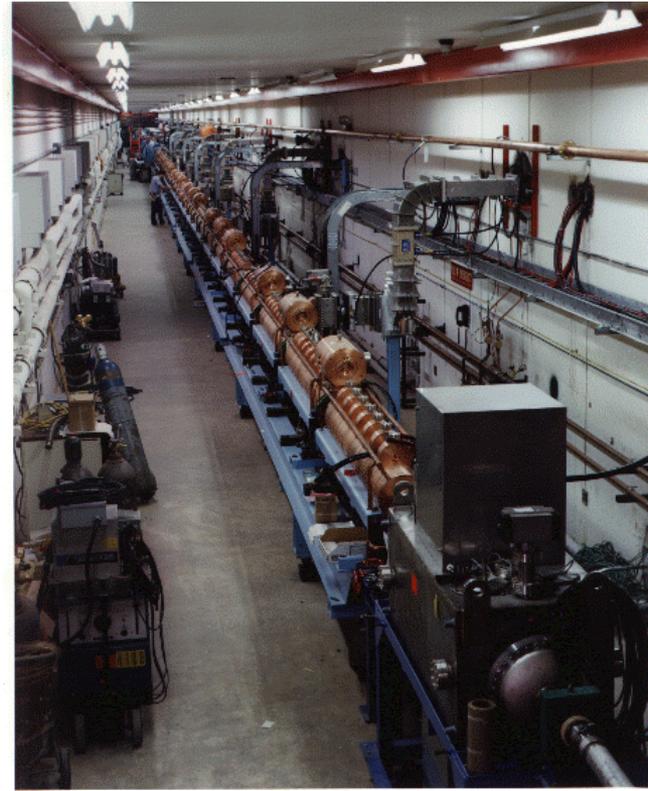


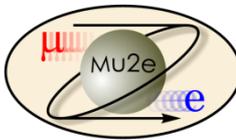
“Preac” - Static Cockcroft-Walton generator accelerates H- ions from 0 to 750 KeV.



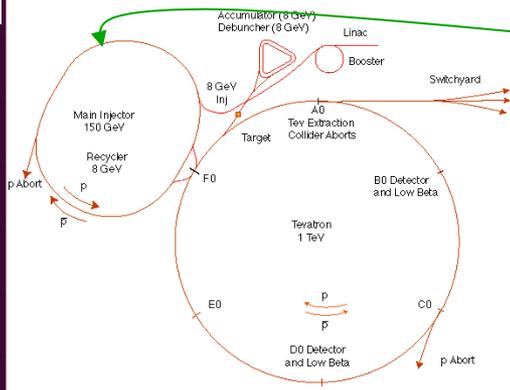
“Old linac”(LEL)- accelerate H- ions from 750 keV to 116 MeV

“New linac” (HEL)- Accelerate H- ions from 116 MeV to 400 MeV

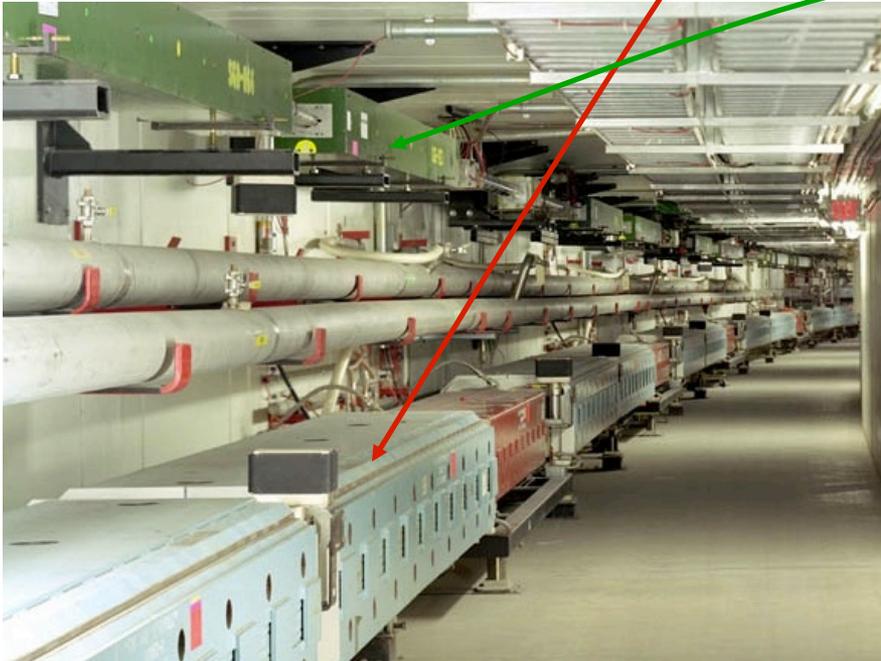




Main Injector/Recycler

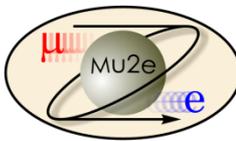


- The **Main Injector** can accept 8 GeV protons OR antiprotons from
 - **Booster**
 - The anti-proton accumulator
 - The **8 GeV Recycler** (which shares the same tunnel and stores antiprotons)
- It can accelerate **protons** to 120 GeV (in a minimum of 1.4 s) and deliver them to
 - The antiproton production target.
 - The fixed target area.
 - The NUMI beamline.
- It can accelerate **protons OR antiprotons** to 150 GeV and inject them into the Tevatron.

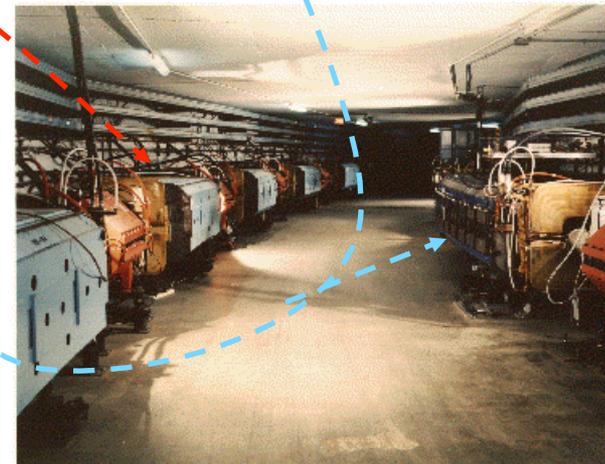
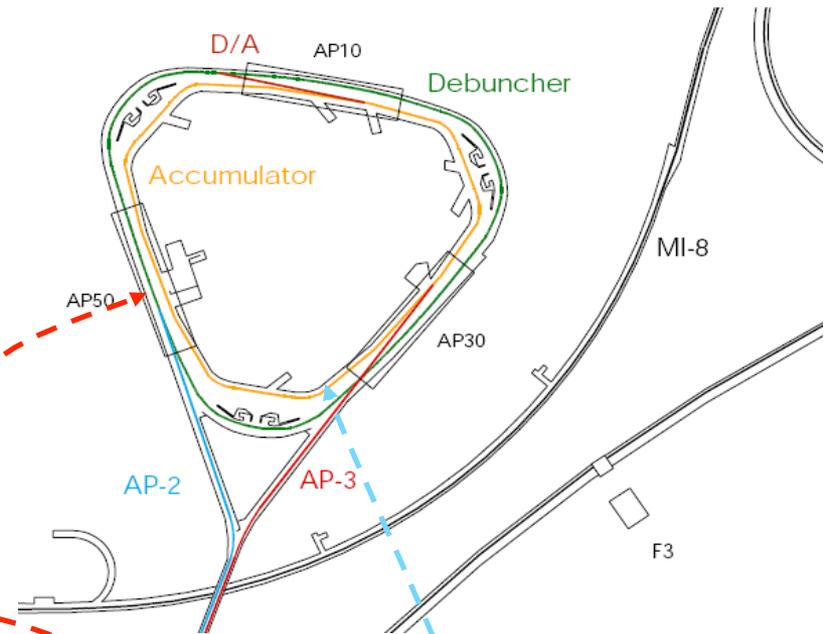




Present Operation of Debuncher/ Accumulator

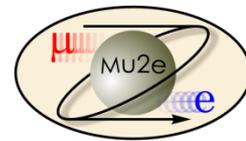


- Protons are accelerated to 120 GeV in Main Injector and extracted to pBar target
- pBars are collected and phase rotated in the “Debuncher”
- Transferred to the “Accumulator”, where they are cooled and stacked
- pBars not used after collider.

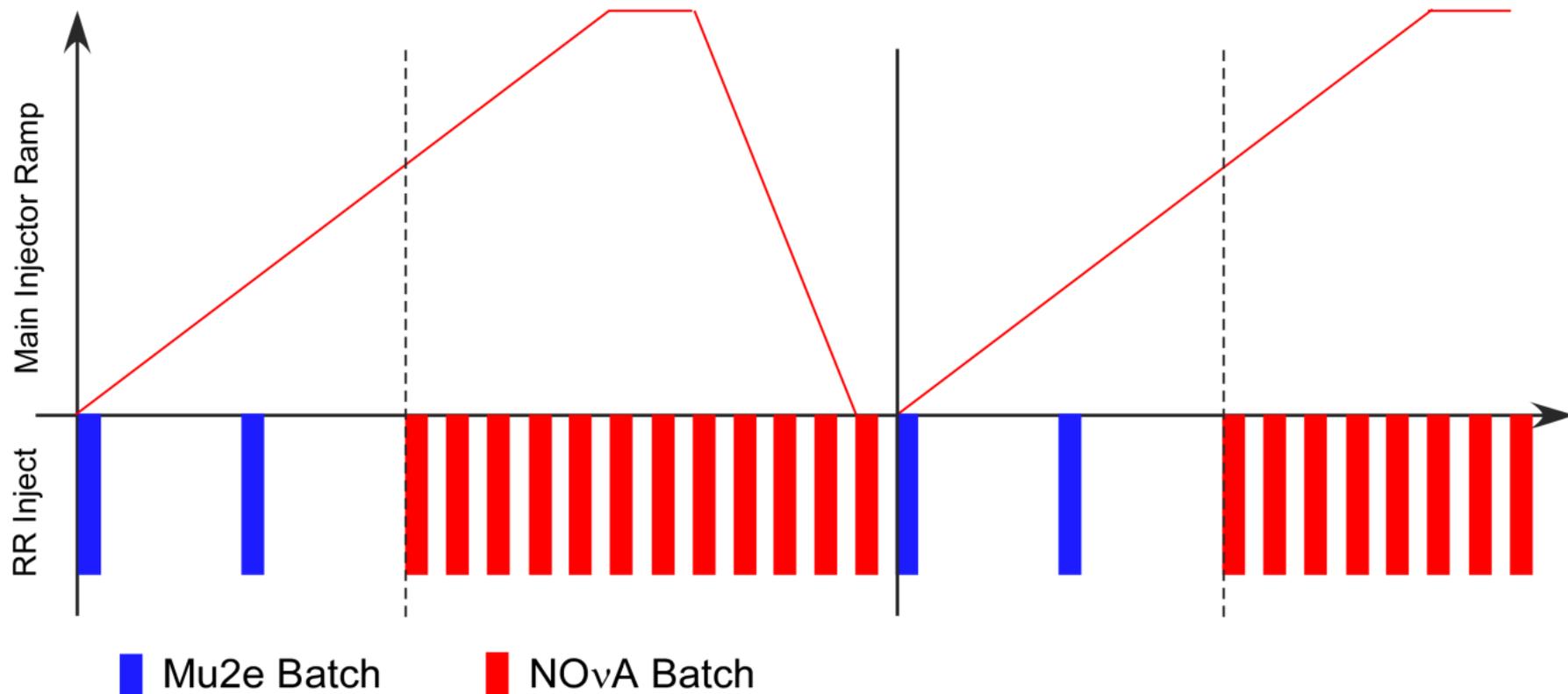


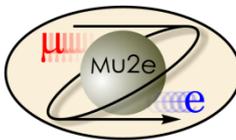


Mu2e in the NOvA era

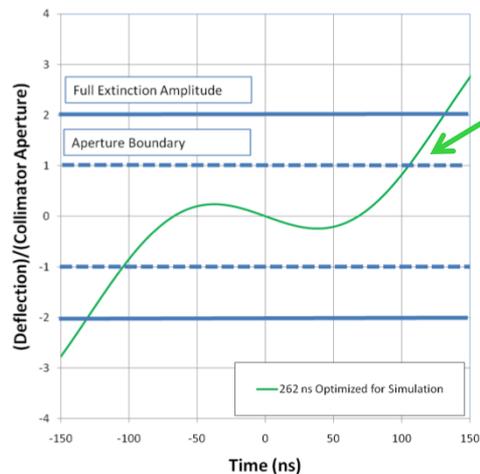


- Beam Delivered in 15 Hz “batches” from the Fermilab Booster



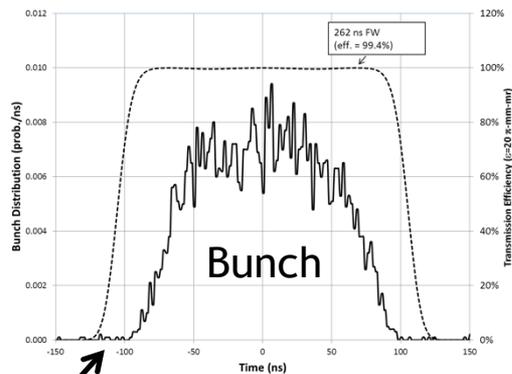


Extinction Performance

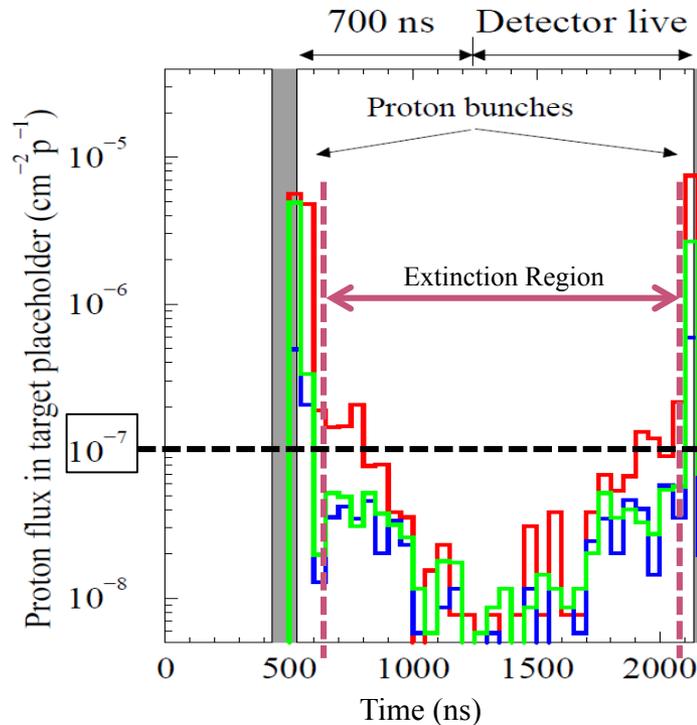


Beam motion in Collimator

Component	Length	Frequency	Peak Field
Low Frequency	3 m	300 kHz	108 Gauss
High Frequency	3 m	3.8 MHz	13 Gauss



Transmission Window



Collimator Material:

- H1–H5: steel
- H1–H5: W
- H1–H3: W, H4–H5: steel

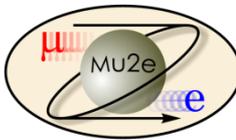
Extinction $< 5 \times 10^{-8}$ over range of interest for optimized collimators

This is multiplied by the Delivery Ring factor to produce a total extinction of $< 5 \times 10^{-12}$

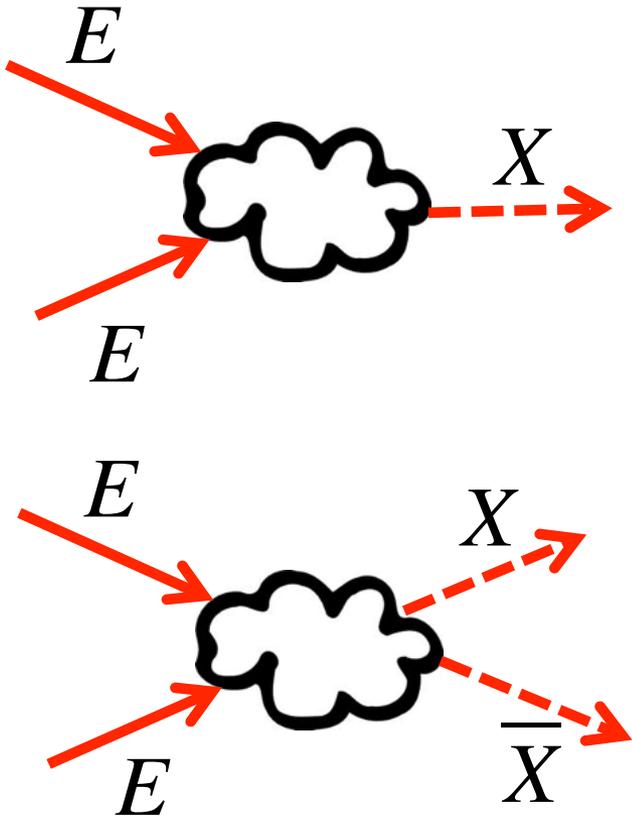
➤ Additional 10^{-5} extinction from beam delivery system



Direct vs. Indirect Observation

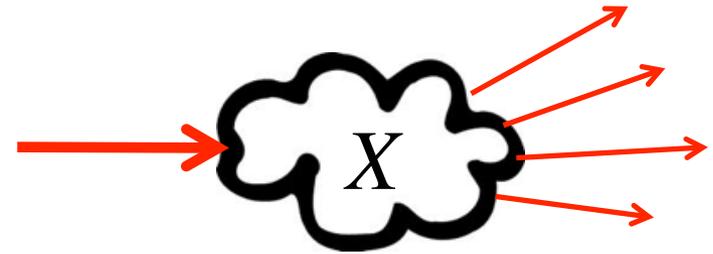


Direct



probe up to $M_X \approx \frac{E}{c^2}$

Indirect

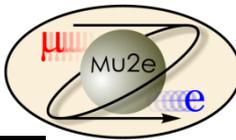


$$M_X c^2 \gg E$$

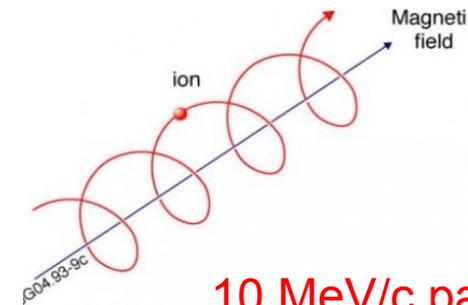
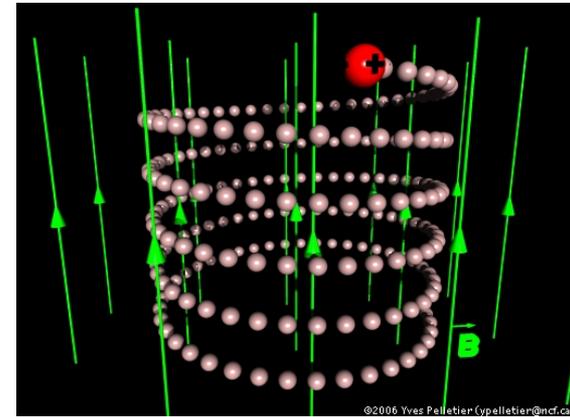
$$\text{Rate} \propto \frac{1}{M_x^4}$$



Review: Particle Motion in a Solenoidal Field



- Generally, particles move in a helical trajectory
- For high momentum particles,
- the curvature is used to measure
- the momentum
- Low momentum particles are effectively “trapped” along the field lines
- A particle trapped along a *curved* solenoidal field will drift out of the plane of curvature with a velocity

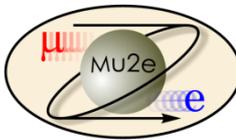


10 MeV/c particle will have a radius of 3 cm in a 1 T field

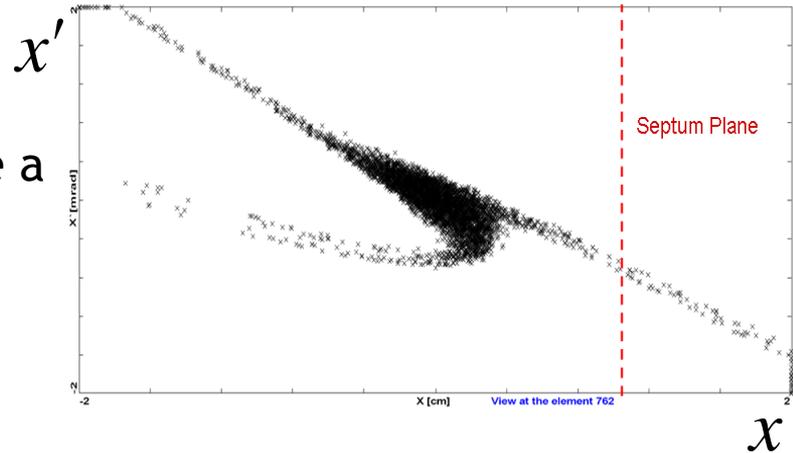
Can be used to resolve charge and momentum! →
$$v_{drift} = \frac{\gamma m}{q} \frac{\hat{R} \times \hat{B}}{RB} (v_{\parallel}^2 + .5v_{\perp}^2)$$



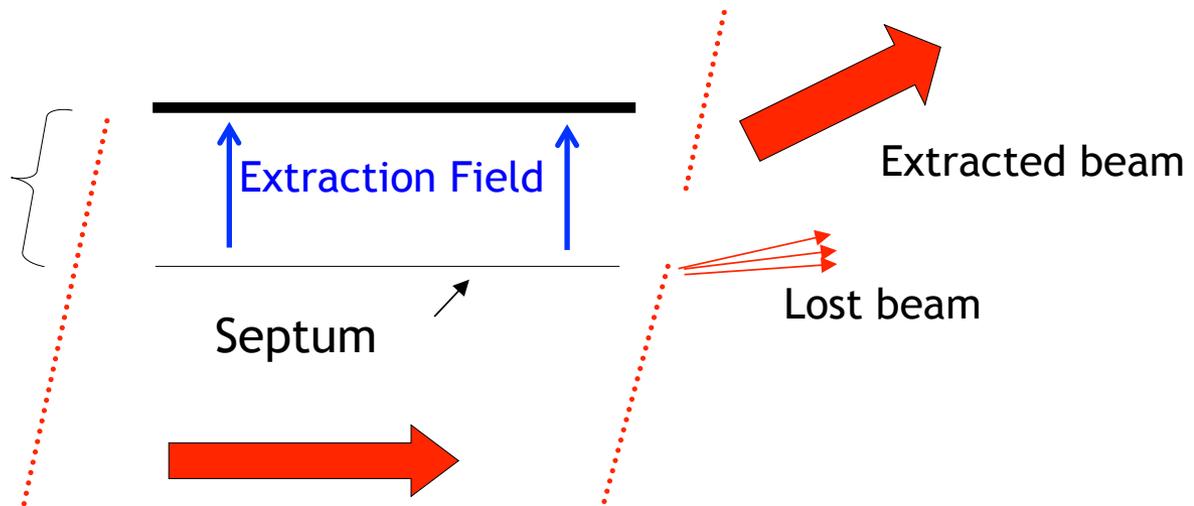
Resonant Extraction



- Extracting all the beam at once is easy, but we want to extract it slowly over ~ 60 ms ($\sim 35,000$ revolutions)
- Use nonlinear (sextupole) magnets to drive a harmonic instability
- Extract unstable beam as it propagates outward
 - Standard technique in accelerator physics



Unstable beam motion
in N (order) turns

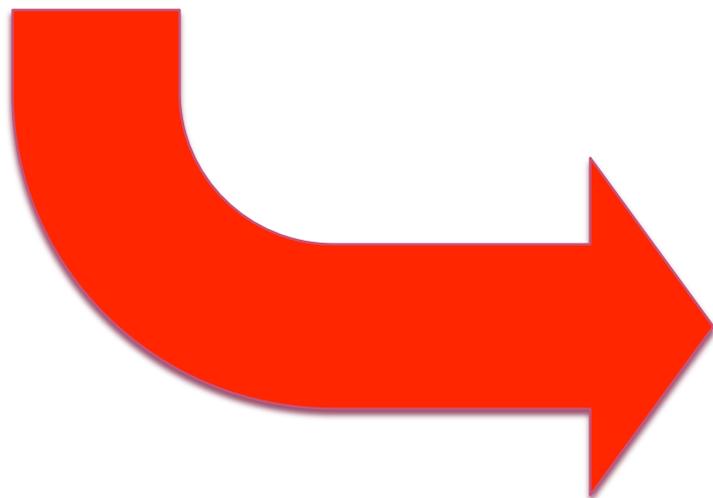
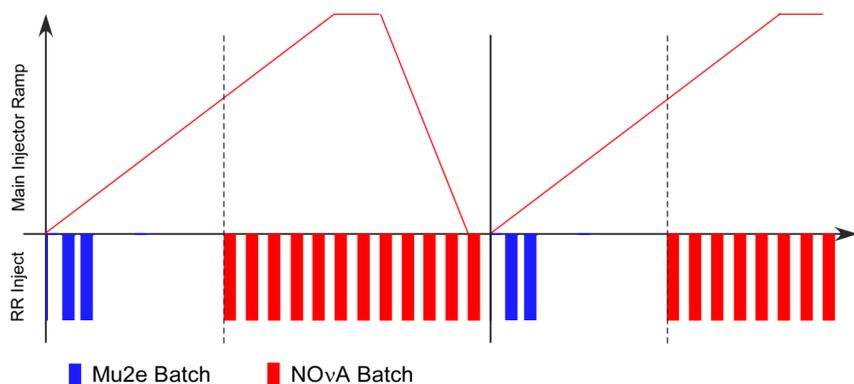




Mu2e Spill Structure

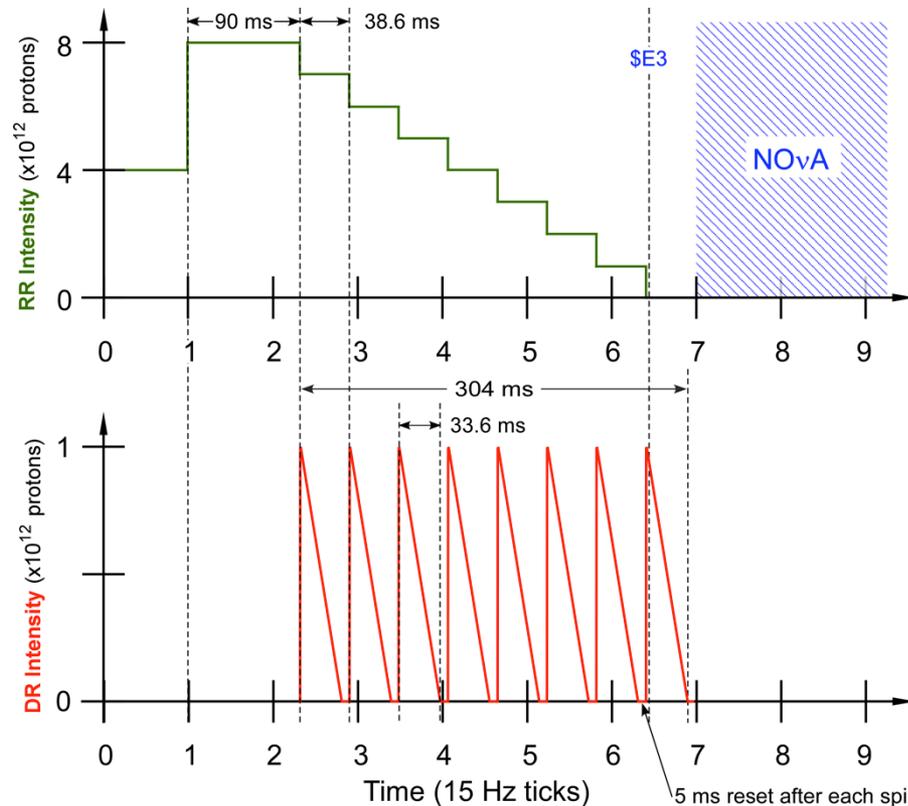


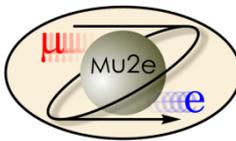
1.33 sec Main Injector cycle



Detail:

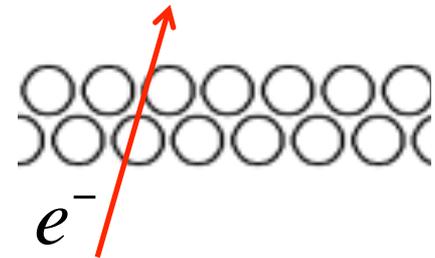
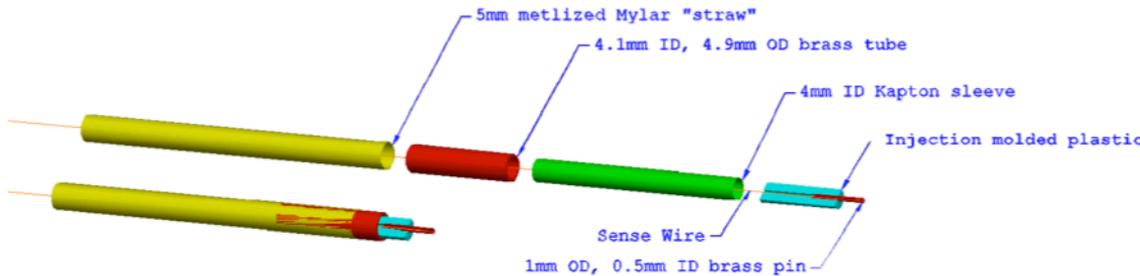
- 3×10^7 p/bunch
- 1.7 μ sec bunch spacing
- ~30% duty factor
- $\sim 1.2 \times 10^{20}$ protons year





Particle Tracking Technology

- To achieve the required resolution, must keep mass as low as possible to minimize scattering
- We've chosen transverse planes of "straw chambers" (~23,000 straws)



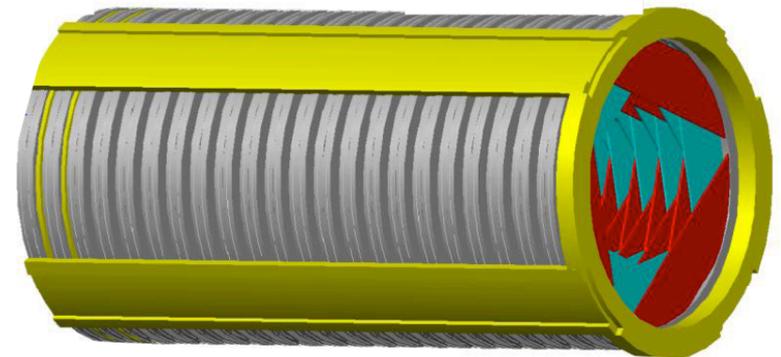
➤ Advantages

- Established technology
- Modular: support, gas, and electronic connections at the ends, outside of tracking volume
- Broken wires isolated

- Track ionizes gas in tube
- Charge drifts to sense wire at center
- Drift time gives precision position

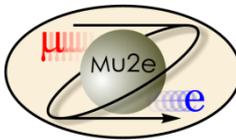
➤ Challenges

- Our specified wall thickness (15 μm) has never been done
- Operating in a vacuum may be problematic



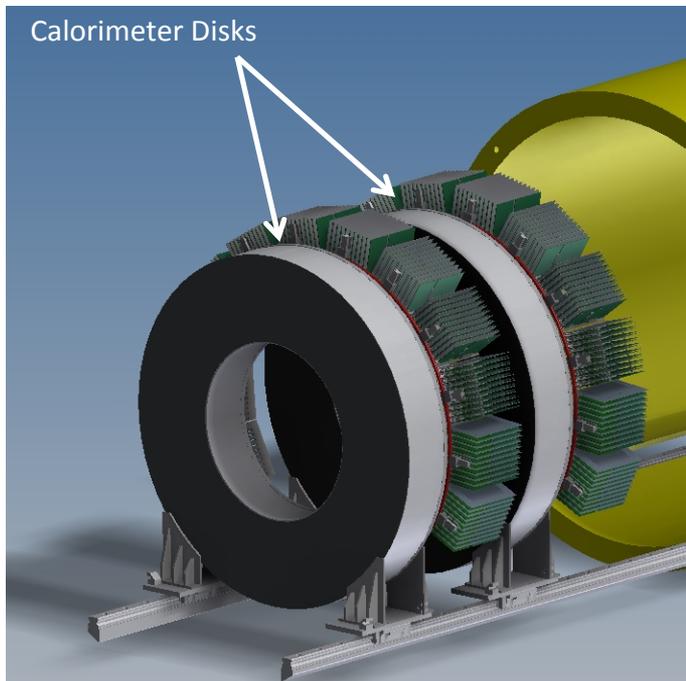


Calorimeter

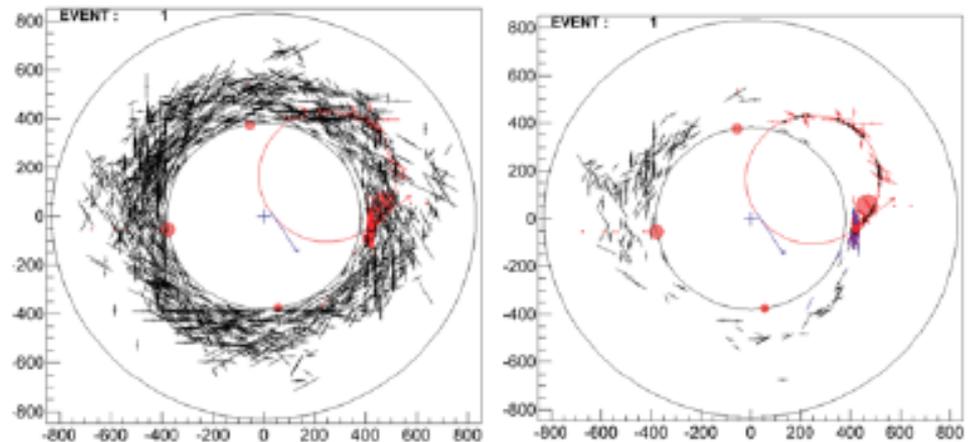


- The Calorimeter will be used to tag electrons
 - Electrons will deposit all of their energy
 - Muons will deposit a small amount of ionization energy
- Two layers of 200 mm long BaF_2 crystals
 - 1860 total

○ Very useful for timing

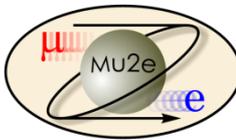


Tracker Hits
Before timing cut After timing cut





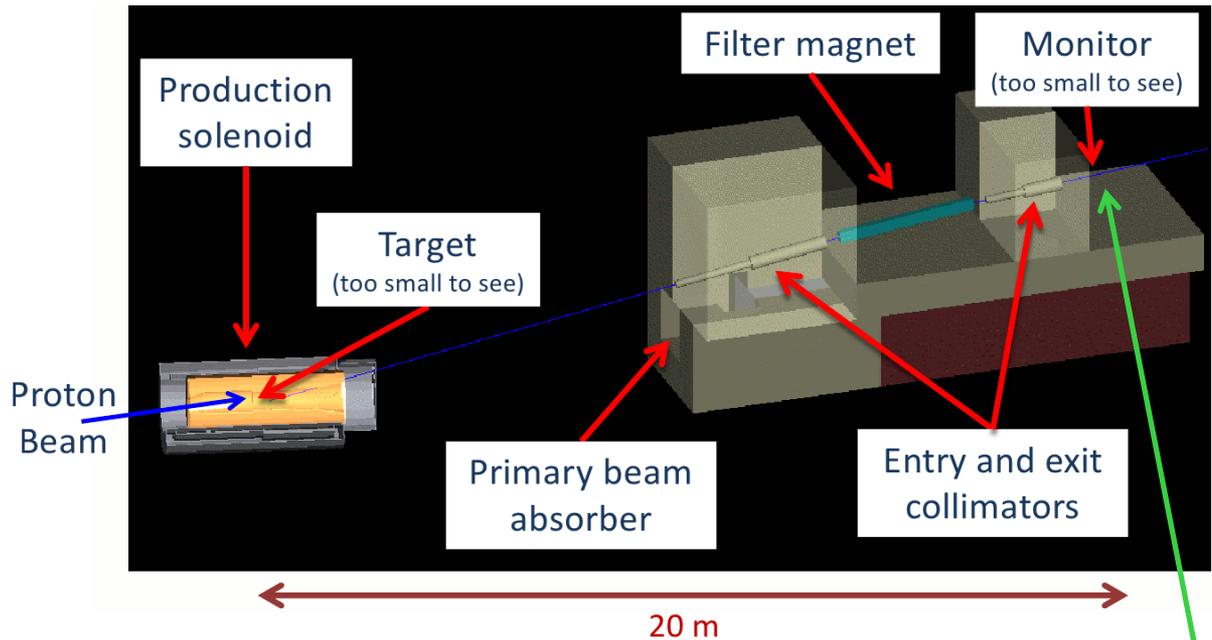
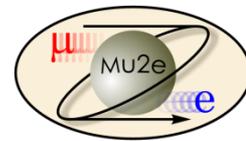
Extinction Monitor



- Achieving 10^{-10} extinction is hard, but it's not useful unless we can verify it.
- Must measure extinction to 10^{-10} precision
 - Roughly 1 proton every 300 bunches!
- Monitor sensitive to single particles not feasible
 - Would have to be blind to the 3×10^7 particles in the bunch.
- Focus on statistical technique
 - Design a monitor to detect a small fraction of scattered particles from target
 - 10-50 per in-time bunch
 - Good timing resolution
 - Statistically build up precision profile for in time and out of time beam.
- Goal
 - Measure extinction to 10^{-10} precision in a few hours

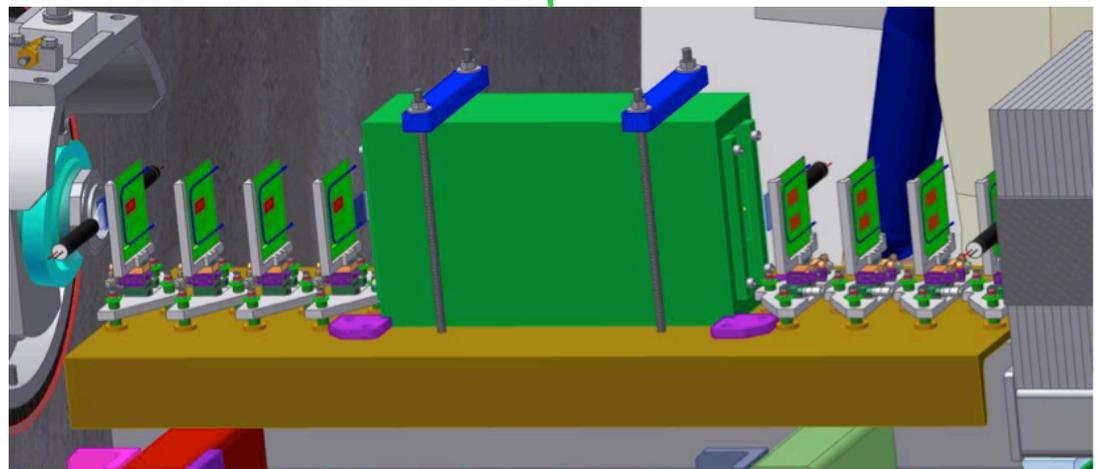


Extinction Monitor Design



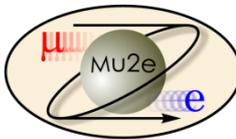
Selection channel built into target dump channel

- Spectrometer based on 8 planes of ATLAS pixels
- Optimized for few GeV/c particles





Target Dependence



- Different models predict different target dependence and different relative rates for $\mu N \rightarrow e N$ and $\mu \rightarrow e \gamma$

V. Cirigliano, R. Kitano, Y. Okada, P. Tuzon., arXiv:0904.0957 [hep-ph]; Phys.Rev. D80 (2009) 013002

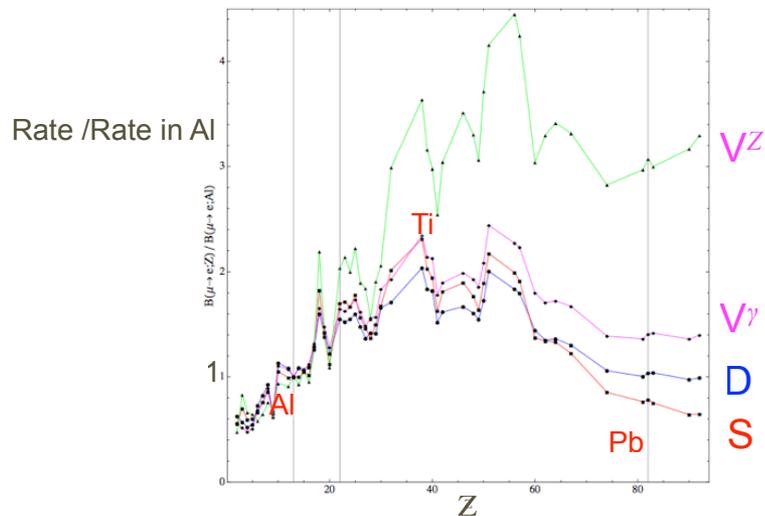
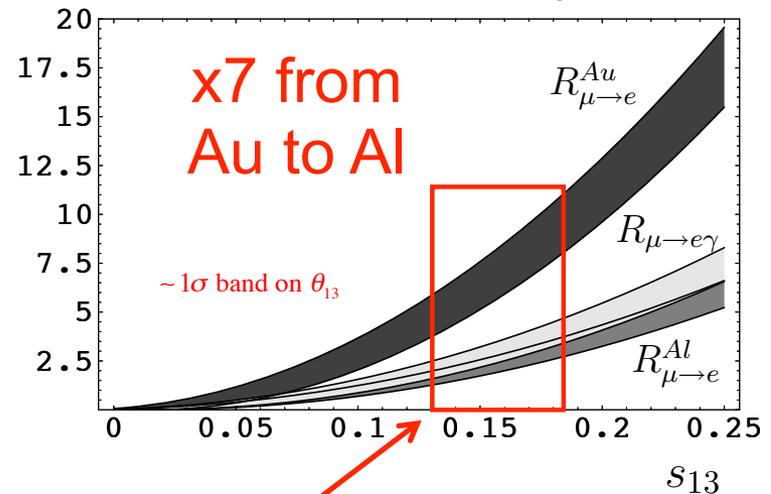


Figure 3: Target dependence of the $\mu \rightarrow e$ conversion rate in different single-operator dominance models. We plot the conversion rates normalized to the rate in Aluminum ($Z = 13$) versus the atomic number Z for the four theoretical models described in the text: D (blue), S (red), $V^{(\gamma)}$ (magenta), $V^{(Z)}$ (green). The vertical lines correspond to $Z = 13$ (Al), $Z = 22$ (Ti), and $Z = 83$ (Pb).

θ_{13} : G. Fogli et al., arXiv:1205.5254



V. Cirigliano, B. Grinstein, G. Isidori, M. Wise
Nucl.Phys.B728:121-134,2005

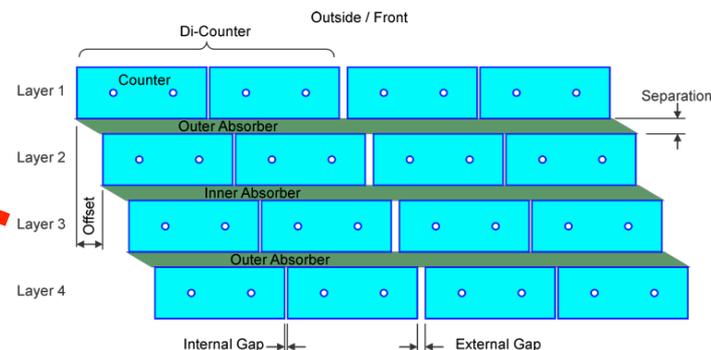
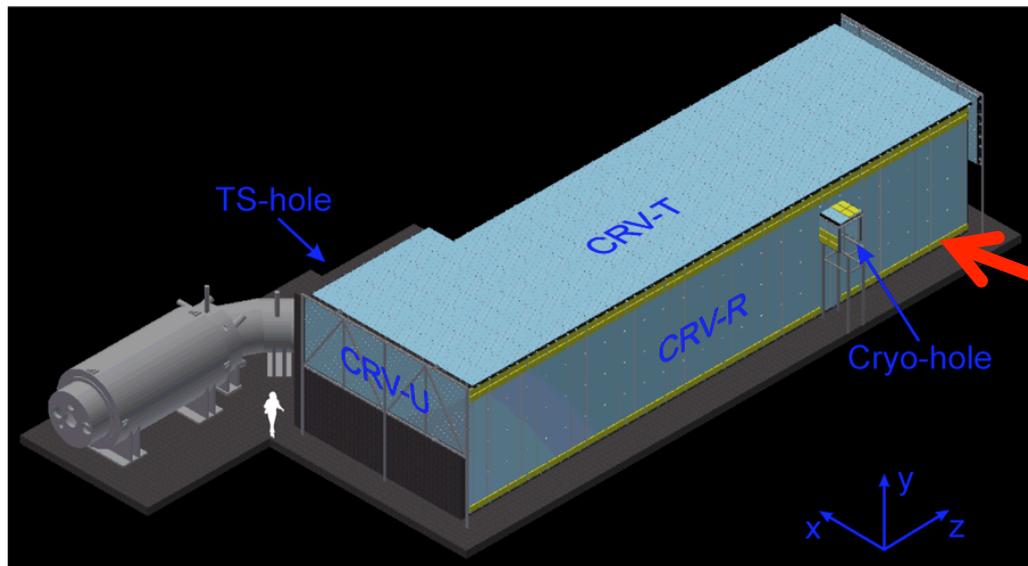
Now we know this!



Cosmic Ray Veto (CRV)



- Multiple layers of scintillator panels surround detector to veto cosmic rays



- Efficiency specification: >99.99%