

# Generating Simplified Models with PYTHIA, Scanning Parameter Space and Presenting the Results

Stephen Mrenna

Fermilab



# Simplified Models

- Process:

$$A_{SM} + B_{SM} \rightarrow C_{SM/NP} + D_{SM/NP}$$

- Particles:

Color – Charge – Mass

- Decays:

$$D \rightarrow E \rightarrow \dots$$

Generate models quickly

Choose masses smartly

Present results usefully



# Particle Properties

```
BLOCK QNUMBERS 9900001 # blobon
    1 0 # 3 times electric charge
    2 1 number of spin states (2S+1)
    3 1 # colour rep (1:singlet,3:triplet,8:octet)
    4 1 # Particle/Antiparticle distinction (0=own anti)
BLOCK QNUMBERS 9900002 # 1blob
    1 6 # 3 times electric charge
    2 1 number of spin states (2S+1)
    3 1 # colour rep (1:singlet,3:triplet,8:octet)
    4 1 # Particle/Antiparticle distinction (0=own anti)
BLOCK QNUMBERS 9900004 # 3blob
    1 3 # 3 times electric charge
    2 1 number of spin states (2S+1)
    3 3 # colour rep (1:singlet,3:triplet,8:octet)
    4 1 # Particle/Antiparticle distinction (0=own anti)

BLOCK MASS
    9900001    700.
    9900002    250.
    9900004    200.
```



# Process and Decays

```
DECAY 9900001 1.00 ! SPECIAL PROCESS PARTICLE
      0.50000      2      21      21
      0.50000      2      9900004      -9900004

DECAY 9900002 1.00
      1.00000      2      -11      -11
DECAY 9900004 1.00
      1.00000      3      1      12      -11
```



# MSUB(482)=1 (MSUB(481)=1 for 2 $\rightarrow$ 1 $\rightarrow$ 2)

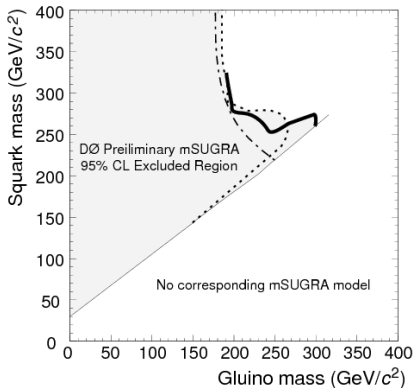
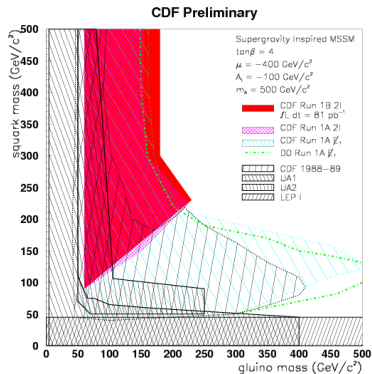
Event listing (summary)

<i>I</i>	<i>particle/jet</i>	<i>KS</i>	<i>KF</i>	<i>orig</i>	<i>p_x</i>	<i>p_y</i>	<i>p_z</i>	<i>E</i>	<i>m</i>
1	!p+	21	2212	0	0.000	0.000	3500.000	3500.000	0.938
2	!p+	21	2212	0	0.000	0.000	-3500.000	3500.000	0.938
=====									
3	!g!	21	21	1	0.000	0.000	346.253	346.253	0.000
4	!g!	21	21	2	0.000	0.000	-251.033	251.033	0.000
5	!g!	21	21	3	0.000	0.000	346.253	346.253	0.000
6	!g!	21	21	4	0.000	0.000	-251.033	251.033	0.000
7	!3blobbar!	21	9900004	0	63.948	56.573	248.709	330.873	200.827
8	!3blob!	21	9900004	0	-63.948	-56.573	-153.489	266.413	200.317
9	!nu_ebar!	21	-12	7	-21.366	90.760	108.690	143.204	0.000
10	!e-	21	11	7	17.823	-48.509	84.768	99.279	0.001
11	!dbar!	21	-1	7	67.491	14.322	55.251	88.391	0.330
12	!nu_e!	21	12	8	17.328	58.108	-89.542	108.142	0.000
13	!e+	21	-11	8	-44.455	-62.735	-53.684	93.776	0.001
14	!d!	21	1	8	-36.821	-51.946	-10.263	64.495	0.330
=====									

(etc)



# mSUGRA transformed to an OSET



$$\tilde{G} \rightarrow \tilde{N}g$$

$g \equiv$  visible, light particle

$\tilde{N} \equiv$  invisible, heavy particle

In the rest frame of  $\tilde{G}$ , the visible momentum from the decay is:

$$p_v = \frac{M_{\tilde{G}}^2 - M_{\tilde{N}}^2}{2M_{\tilde{G}}} = \frac{M_{\tilde{G}}}{2}(1 - r^2), r = M_N/M_G$$

Assume  $\tilde{G}$  is produced in pairs and at threshold

max visible momentum is  $M_G(1 - r^2)$

Motivation behind the discriminant known as  $R$ , or “the razor.”



# Expectations

Models with similar values of  $p_v$  yield similar kinematic distributions?

The event rate will depend upon coupling strength, spin, incoming partons, etc., but not (so much) acceptance

Propose to scan parameter space (and present results limits?) on  $\sigma \times BR^2$  in the  $M_G - p_v$  plane

Model lines defined by orthogonal directions in this space

Alternative to  $\Delta M = M_G(1 - r)$





# Definitions

The physical observable shown is the magnitude of the missing transverse momentum  $mH_T$ , defined using observed jets:

$$mH_T = \left| \sum_i \vec{p}_{Ti} \right|.$$

For one-step decays, one could alternatively show the  $H_T$ , defined as the sum of the magnitude of the individual jet transverse momenta;

$$H_T = \sum_i \left| \vec{p}_{Ti} \right|.$$

Results with PYTHIA and FASTJET anti-kT jets<sup>1</sup>

---

<sup>1</sup>For speed and simplicity, the multiple-interaction model of PYTHIA is not simulated



$p_\nu = M_G(1 - r^2) = 175 \text{ GeV}$  (left) versus  $\Delta M = 100 \text{ GeV}$  (right)

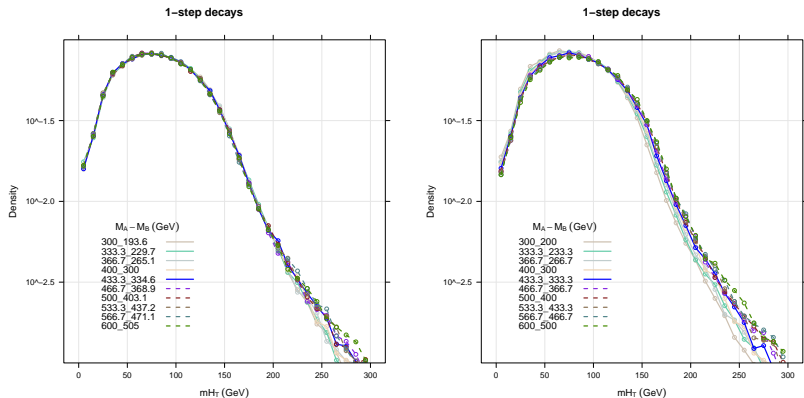
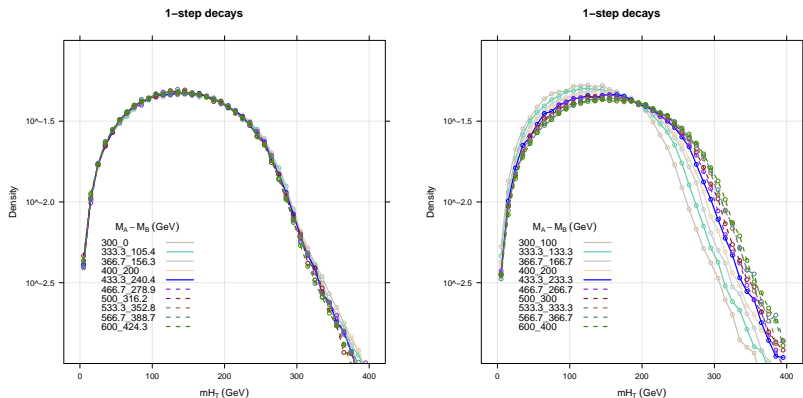


Figure: Distributions of the missing vector transverse momentum  $mH_T$ : (left) for a fixed value of the  $p_\nu$ ; (right) for a fixed mass splitting.



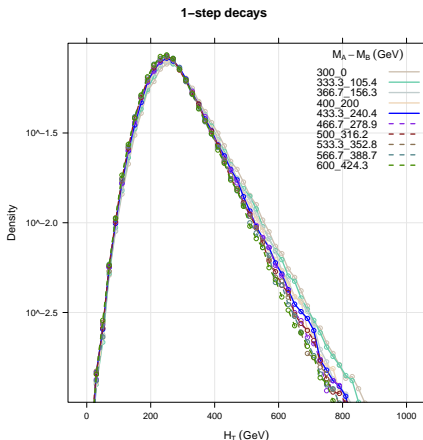
$p_\nu = 300$  GeV (left) versus  $\Delta M = 200$  GeV (right)



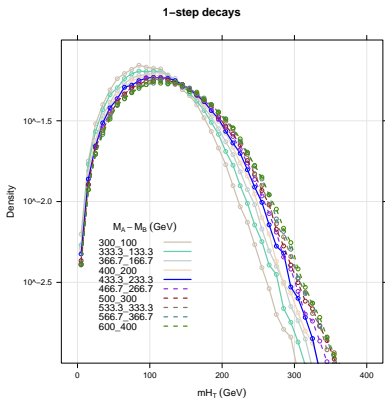
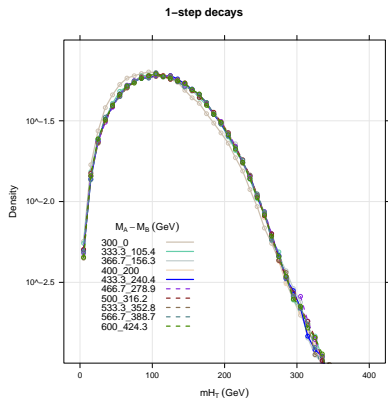
**Figure:** Distributions of the missing vector transverse momentum  $mH_T$ : (left) for a fixed value of the  $p_\nu$ ; (right) for a fixed mass splitting. Similar to Fig. 1, but with a larger mass splitting.



Finally, we note that for such 2-body decays, the  $H_T$  and  $mH_T$  distributions are highly correlated:



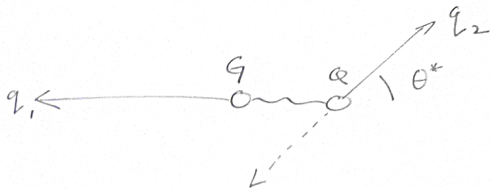
$p_\nu = 300$  GeV (left) versus  $\Delta M = 200$  GeV (right)



The tight correlation between  $mH_T$  and  $H_T$  is reduced.



$$\tilde{G} \rightarrow q_1 \tilde{Q} \rightarrow \tilde{q}_1 q_2 N$$



**Figure:** The kinematic setup for two-step decays.

$$\vec{p}_v = \hat{z}(p_{q_1} + \gamma_Q(p^* \cos \theta^* + \beta_Q E^*)) + \hat{\perp} p^* \sin \theta^*$$

$$-p_{q_1} = p_Q = \frac{M_g^2 - M_Q^2}{2M_g}; p^* = \frac{M_Q^2 - M_N^2}{2M_Q}$$



# “Best” result

Average over a (flat) decay in the rest frame of the cascade particle:

$$\langle p_v \rangle = \frac{M_G}{2(1-a^2)(1-b^2)} \left( 1 - a^4 b^4 - (a^2 + b^2)|a^2 - b^2| + 4a^2 b^2 \log \left( \frac{2a^2 b^2}{a^2 + b^2 - |a^2 - b^2|} \right) \right),$$

$a = M_Q/M_G$  and  $b = M_N/M_Q$

pair production  $\rightarrow 2\langle p_v \rangle$



$(M_G, p_v) = (400, 260)$  GeV (left) and  $(500, 375)$  GeV (right)

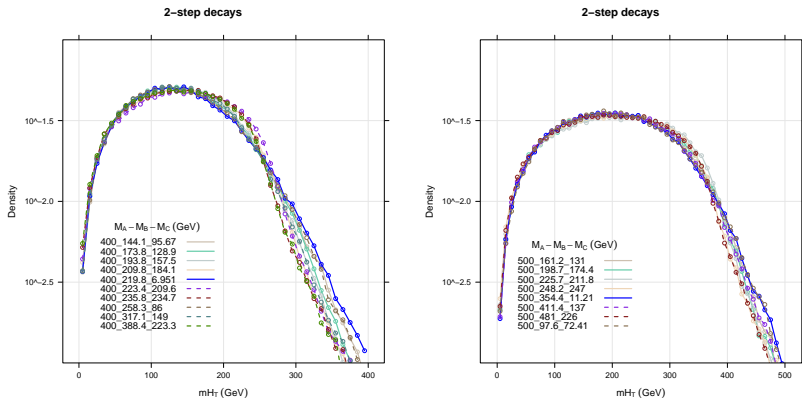


Figure:  $mH_T$  distributions for fixed mother mass and visible momentum  $p_v$





As a first step, consider the decay  $C^\pm \rightarrow \ell^\pm \tilde{\nu}$ .

$$p_\nu = \frac{M_C}{2} \left( 1 - \frac{M_\nu^2}{M_C^2} \right).$$

A related decay chain is  $C^\pm \rightarrow \nu L^\pm \rightarrow \ell^\pm N$ . The above derivation needs a modification to handle the loss of momentum from neutrinos.

$$p_\nu = \frac{M_C}{2} \left( 1 - \frac{M_N^2}{M_L^2} \right) \frac{1}{2} \left( 1 + \frac{M_L^2}{M_C^2} \right).$$

The distributions for fixed  $p_\nu$  are **not** “constant”



- PYTHIA can easily generate simplified models
- The sensitivity of distributions to masses appears in ratios of squares
- Certain decay sequences yield “invariants” under the exchange of these ratios
- We can use this to sample model-space efficiently, or to present results that reflect the sensitive parameters

